

# FYST11 Lecture 12

## BSM II

Thanks to G. Brooijmans, T. Rizzo, L.  
Covi

# This week's topics

- Why go Beyond the SM?
  - What are the problems with the SM?
  - What direct measurements points to physics BSM
- Some attempts at solutions
  - Supersymmetry
  - Extended Higgs sector
  - Extra dimensions
  - A few others
- Searches for DM

# Supersymmetry (SUSY)

# Idea

New symmetry *fermions*  $\leftrightarrow$  *bosons*

This symmetry is the most general extension of Lorentz invariance

SUSY has:  $N_{\text{dof}}(\text{bosons}) = N_{\text{dof}}(\text{fermions})$   
 [cf. SM:  $N_{\text{dof}}(\text{bosons}) \ll N_{\text{dof}}(\text{fermions})$ ]

Spin 0	Spin 1/2	Spin 1	Spin 3/2	Spin 2
sLeptons	Leptons		Gravitino	Graviton
sQuarks	Quarks			
Higgs	Higgsino			
	Photino	Photon		
	Zino	Z		
	Wino	W		
	Gluino	Gluon		

- To create *supermultiplets*, we need to add one *superpartner* to each SM particle
- Superpartners have opposite spin statistics but otherwise equal quantum numbers
- Need to introduce an additional Higgs doublet to the non-SUSY side  $\Rightarrow$  5 Higgs bosons

But where are these partners?!

Supersymmetry must be broken (if realized)

# Particle spectrum (minimal!)

In reality the new states would *mix*

Several ideas of how the supersymmetry is broken – intimately connected with EWK symmetry breaking

		Spin 0	Spin 1/2	Spin 1
<div style="border: 1px solid black; padding: 2px; display: inline-block; margin-bottom: 5px;">SM</div> <div style="border: 1px solid black; padding: 2px; display: inline-block; margin-bottom: 5px;">SUSY</div>	Eigenstates of mass	$\tilde{\ell}_1, \tilde{\ell}_2$	$\ell$	
		$\tilde{q}_1, \tilde{q}_2$	$q$	
		$h^0, H^0, A^0, H^\pm$	$\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$	
			$\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$	$\gamma, Z^0, W^\pm$
		$\tilde{g}_a$	$g_a$	

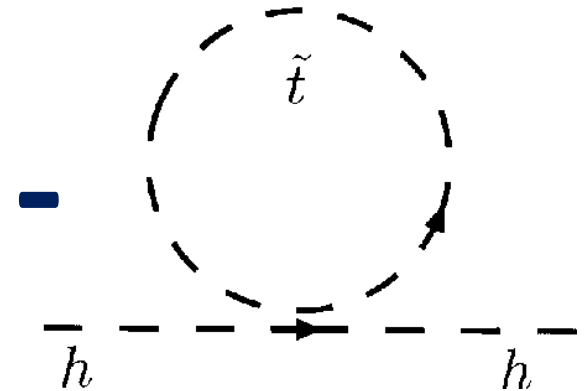
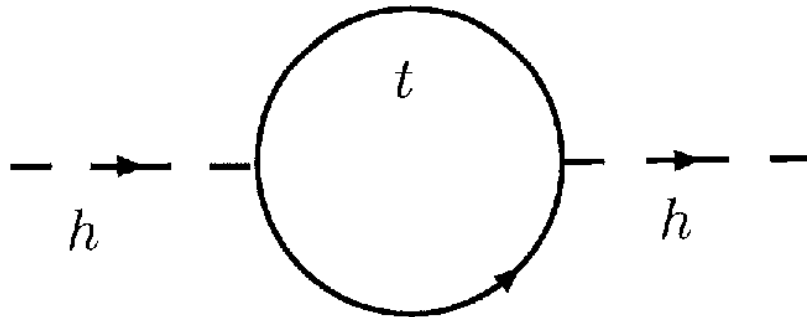
Squark/slepton mixing proportional to SM partner masses  
 → largest for 3<sup>rd</sup> gen.  
 → can become lightest squarks / sleptons

The gauge-mixed physical states that propagate in space and time and that can be observed.  
 Neutralinos: mass eigenstates of photinos, zinos, neutral higgsinos  
 Charginos : mass eigenstates of winos and charged higgsinos

Since we don't know the mechanism, have to introduce  $\mathcal{O}(100)$  new parameters

# SUSY and the hierarchy problem

If Supersymmetry not broken we would have perfect cancellation in the loops!



But as  $m(\tilde{t}) \neq m(t)$  they do not quite cancel, instead just a suppression

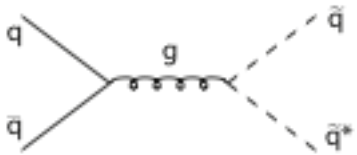
This still gives a decent result if

$$|m(\text{fermion}) - m(\text{boson})| < \mathcal{O}(\text{TeV})$$

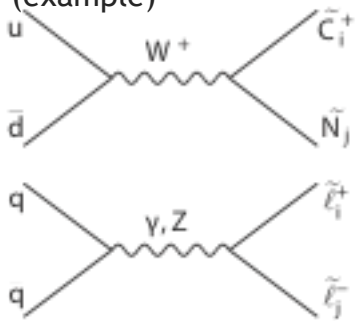
# Once mass spectrum fixed, all cross sections predicted

Spin structure of SUSY spectrum: lower  $\sigma$  than other BSM models, harder to find !

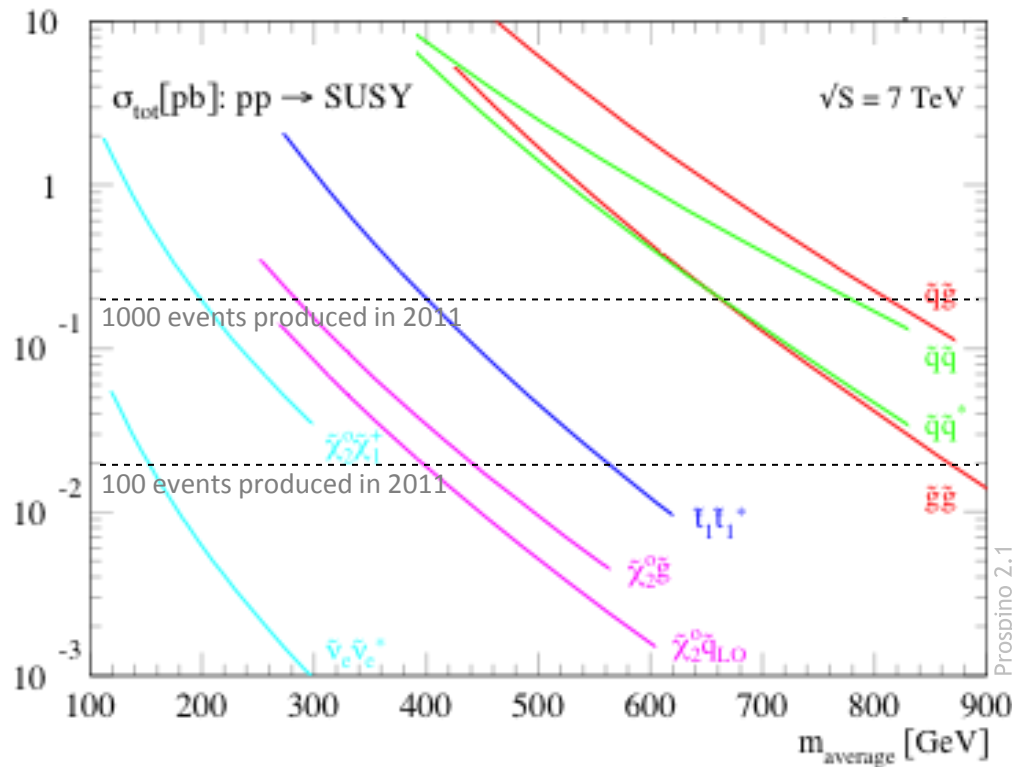
Direct squark pair production (example)



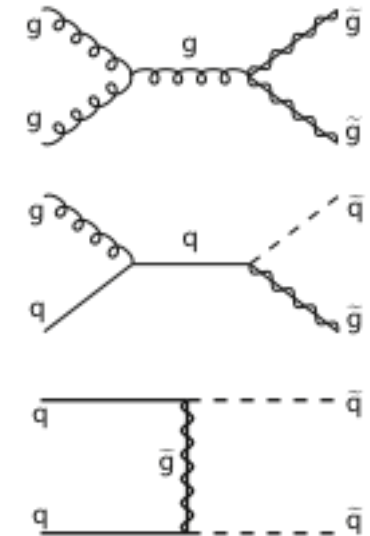
Direct gaugino/slepton pair production (example)



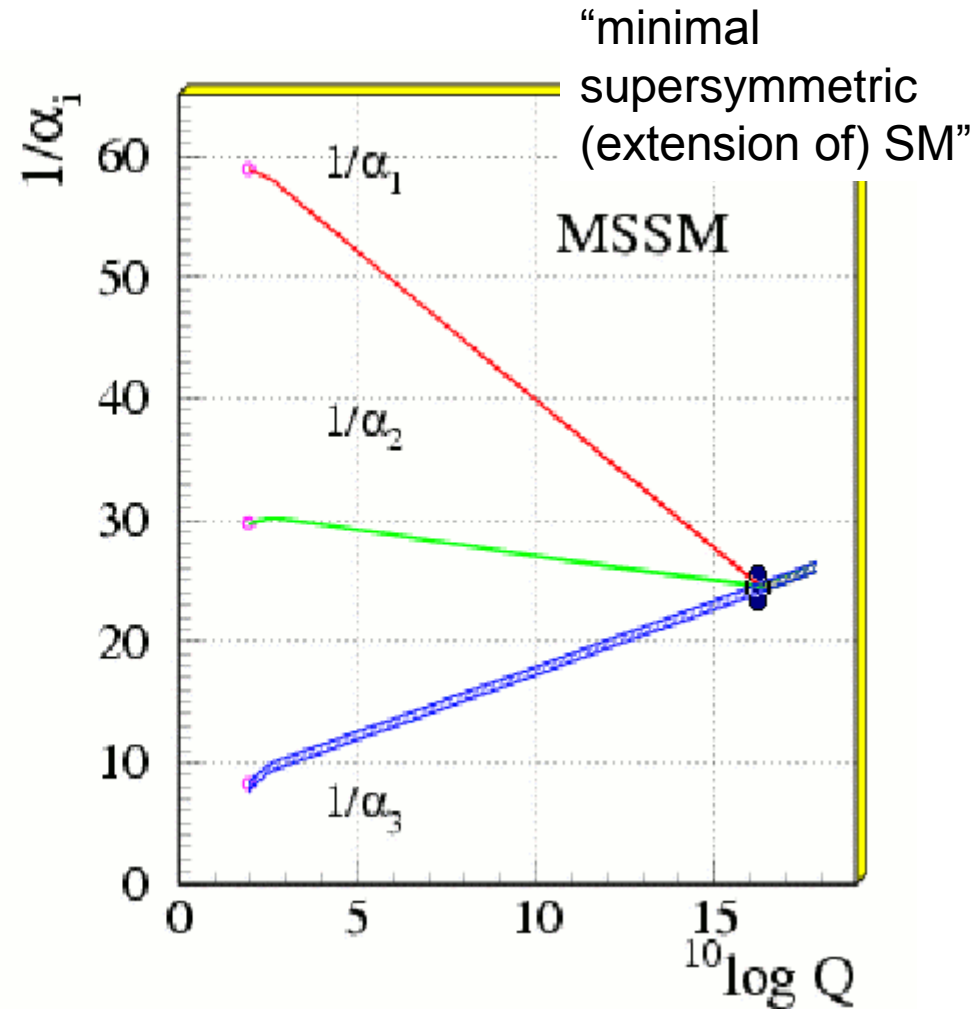
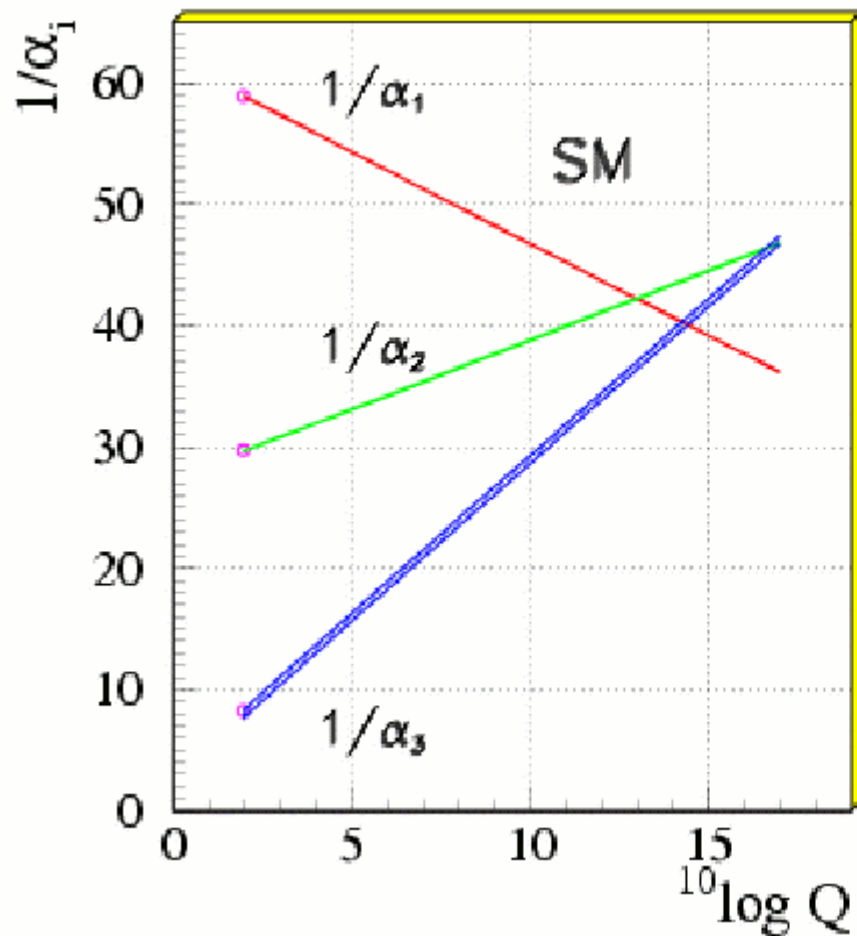
SUSY cross section versus sparticle mass



Gluino & squark production (examples)



# Unification of coupling constants with supersymmetry



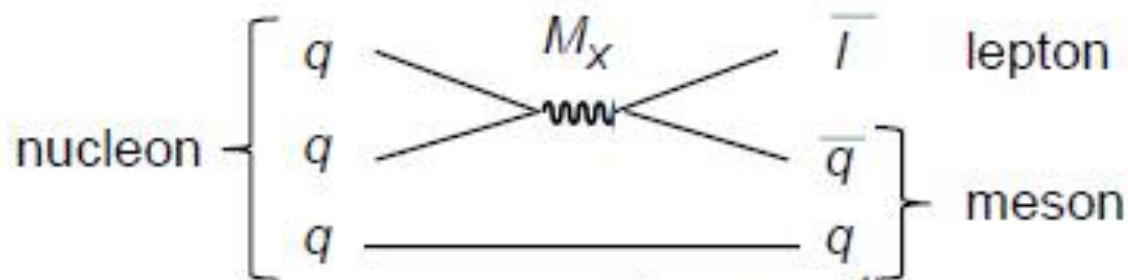


# Proton Decay

(G. Giudice SSLP'15)

in GUT, matter is unstable

decay of proton mediated by new SU(5)/SO(10) gauge bosons



$$\text{GUT: } \tau_p(p \rightarrow e^+ \pi^0) = \left( \frac{M_X}{10^{15} \text{ GeV}} \right)^4 10^{31-32} \text{ yr}$$



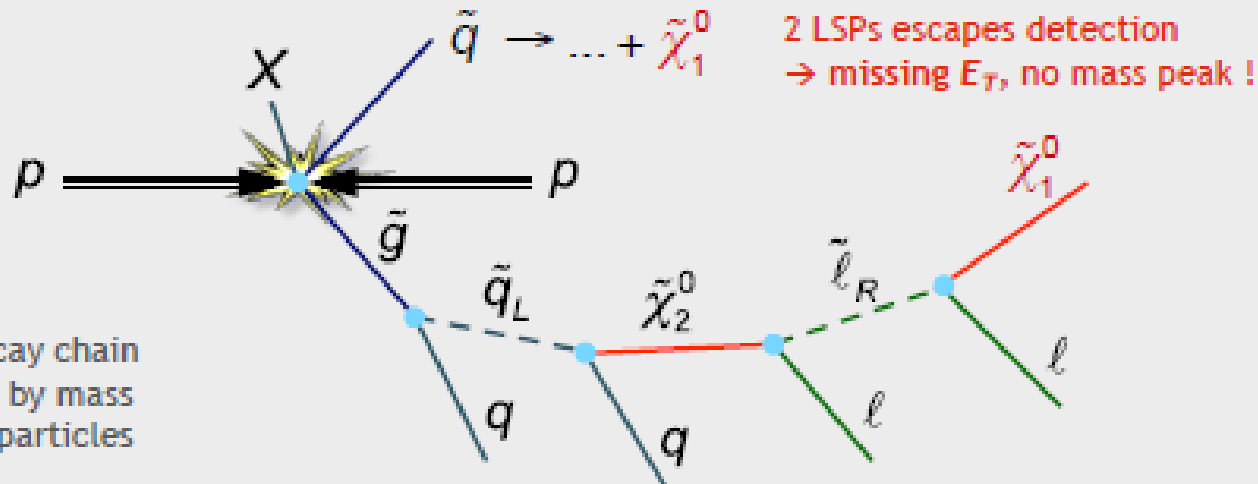
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$$\text{Exp: } \tau_p(p \rightarrow e^+ \pi^0) > 8.2 \times 10^{33} \text{ yr}$$

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# Characteristic SUSY Decay Cascades

- To avoid proton decay, a new conserved quantum number ( $R$ ) is introduced, which forces a SUSY particle to decay in at least one other SUSY particle
- The lightest SUSY particle is thus stable (LSP), and must be neutral and colourless  $\rightarrow$  WIMP (dark matter candidate)
- Typical **LSP is spin- $\frac{1}{2}$  neutralino**. It could also be a gravitino
- With  $R$  parity: SUSY production in pairs only  $\rightarrow$  requires energy  $2 \times$  SUSY mass !

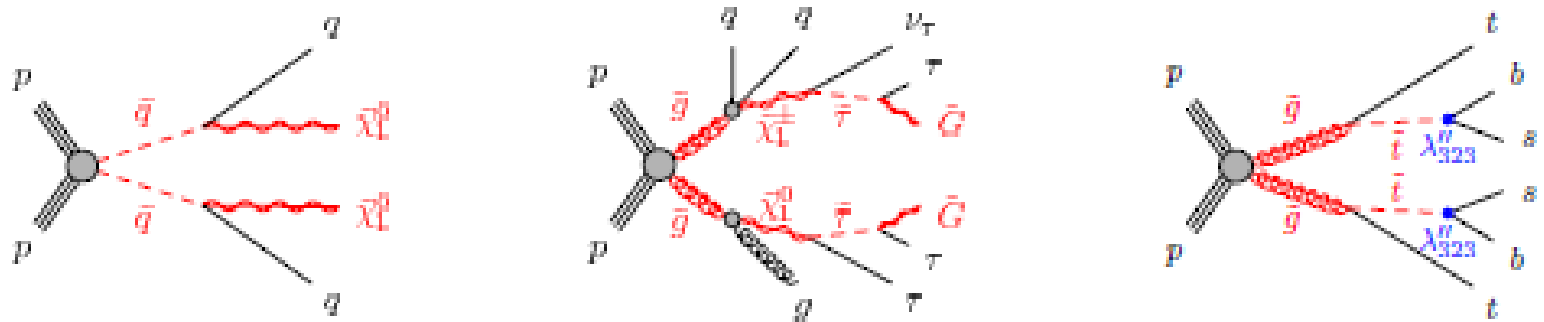


"Typical" SUSY decay chain at the LHC, driven by mass hierarchy of SUSY particles

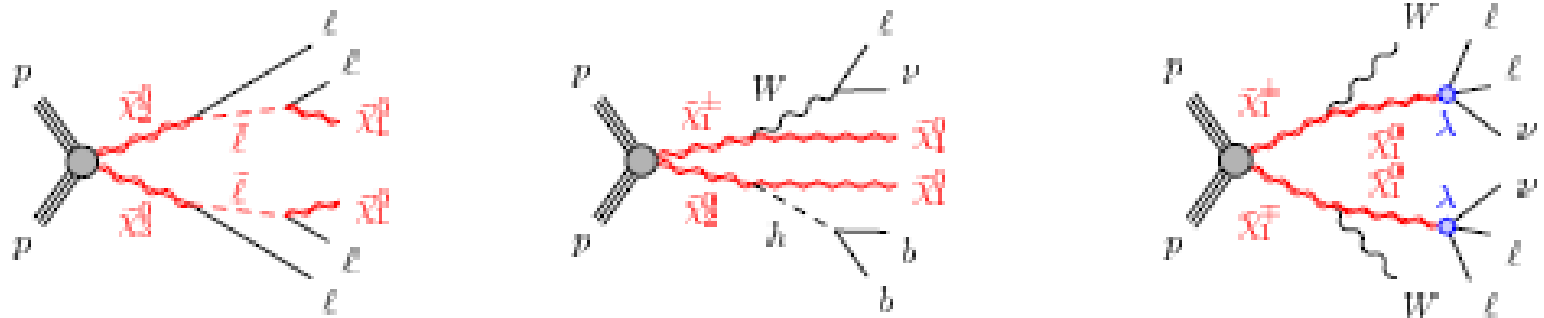
# Canonical SUSY

## ❖ Wide range of signatures

- Strong production... (large cross-section)



- ... or weak production



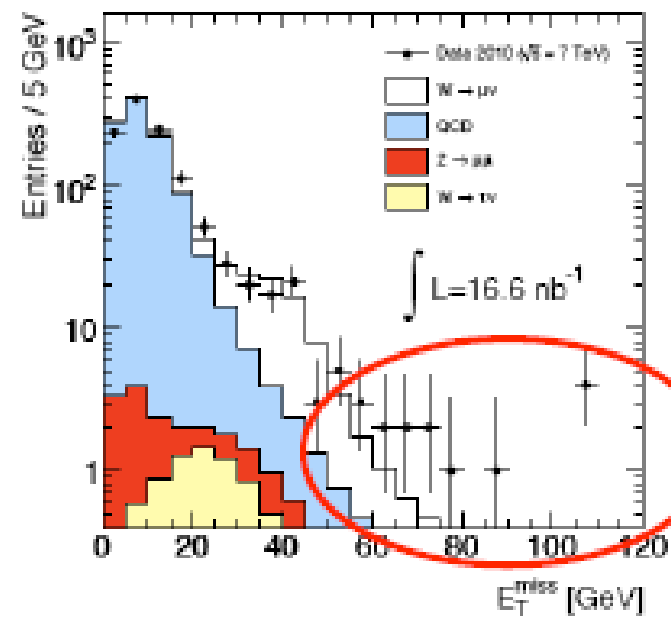
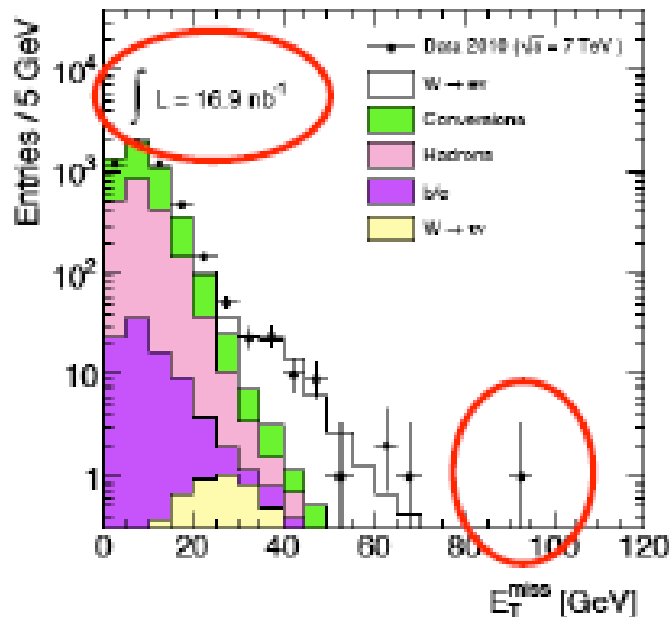
RPV

# Missing ET

❖ “Evil” variable: -  $\Sigma$  (everything else)

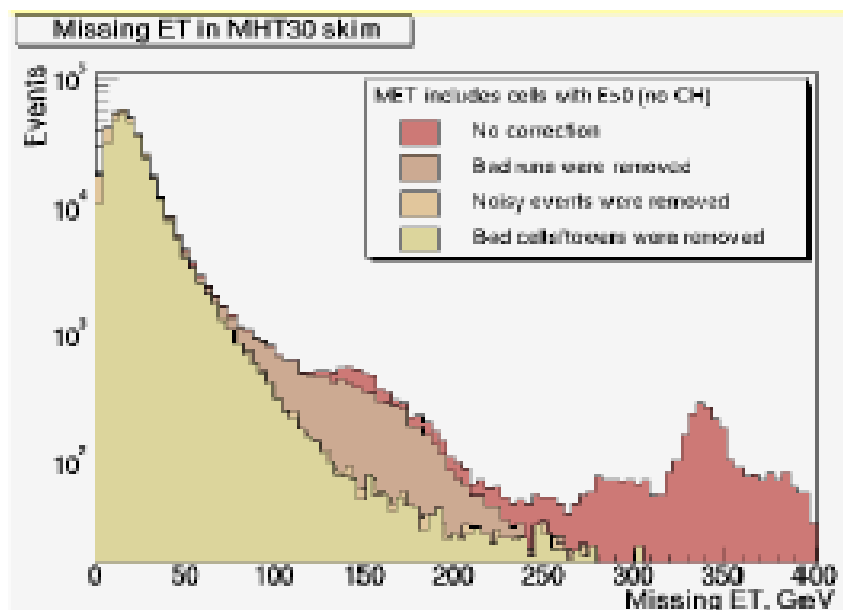
- Need to understand “everything else”
- Good benchmark: leptonic W boson decays

Early 2010

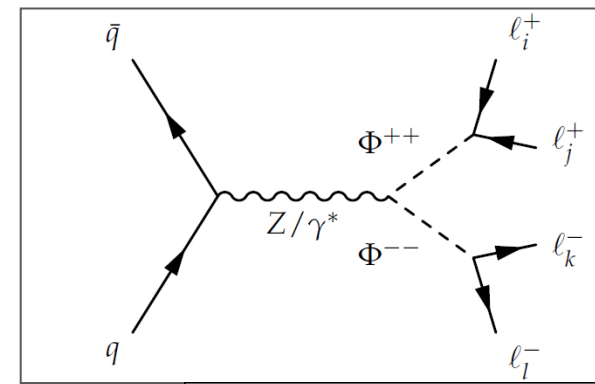


## ❖ Analyses using MET are particularly sensitive

- Requires the full calorimeter to behave, and calorimeter is generally the most sensitive subdetector (analog,  $\sim 16$  bits)
- Easy: basic DQ (high voltage trip, etc.)
- Hard: low frequency
- Can't spot a  $10^{-5}$  Hz (once a day) effect online or in first pass DQ
- But can be biggest part of dataset after cuts!

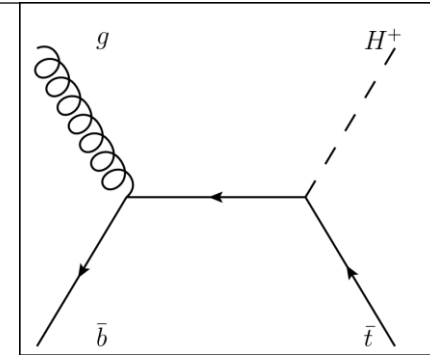


# Extended Higgs sector



In the Standard Model single Higgs doublet, often

written as  $\begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix}$  or  $\begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix}$



**Extended:** Many choices but a few constraints,

for instance suppression of FCNC and  $\frac{M_W}{M_Z} = \cos \theta_W$

- Most successful: 2 Higgs doublet models (2HDMs)
  - Supersymmetry uses this
- See-saw models predict Higgs triplet with  $\varphi^0, \varphi^{+/-}, \varphi^{++/--}$

# General 2HDM Potential

$$\begin{aligned}
 V(\phi_1, \phi_2) = & \lambda_1 \left( |\phi_1|^2 - v_1^2 \right)^2 + \lambda_2 \left( |\phi_2|^2 - v_2^2 \right)^2 \\
 & + \lambda_3 \left[ \left( |\phi_1|^2 - v_1^2 \right) + \left( |\phi_2|^2 - v_2^2 \right) \right]^2 \\
 & + \lambda_4 \left[ |\phi_1|^2 |\phi_2|^2 - \left( \phi_1^{*T} \phi_2 \right) \left( \phi_2^{*T} \phi_1 \right) \right] \\
 & + \lambda_5 \left[ \text{Re} \left( \phi_1^{*T} \phi_2 \right) - v_1 v_2 \cos \xi \right]^2 \\
 & + \lambda_6 \left[ \text{Im} \left( \phi_1^{*T} \phi_2 \right) - v_1 v_2 \sin \xi \right]^2
 \end{aligned}$$

*All  $\lambda$  are real.*

*From "Higgs Hunter's guide".*

$$\begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}_1 = \begin{pmatrix} 0 \\ v_1 \end{pmatrix} \frac{1}{\sqrt{2}} \quad ; \quad \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}_2 = \begin{pmatrix} 0 \\ v_2 e^{i\xi} \end{pmatrix} \frac{1}{\sqrt{2}} \quad ; \quad \tan \beta = \frac{v_2}{v_1}$$

# Higgs Boson Spectroscopy

- One Charged Higgs with mass:

$$m_{H^\pm} = \sqrt{\lambda_4 (v_1^2 + v_2^2)}$$

- One CP-odd neutral Higgs with mass:

$$m_{A^0} = \sqrt{\lambda_6 (v_1^2 + v_2^2)}$$

- And two CP-even higgs that mix.

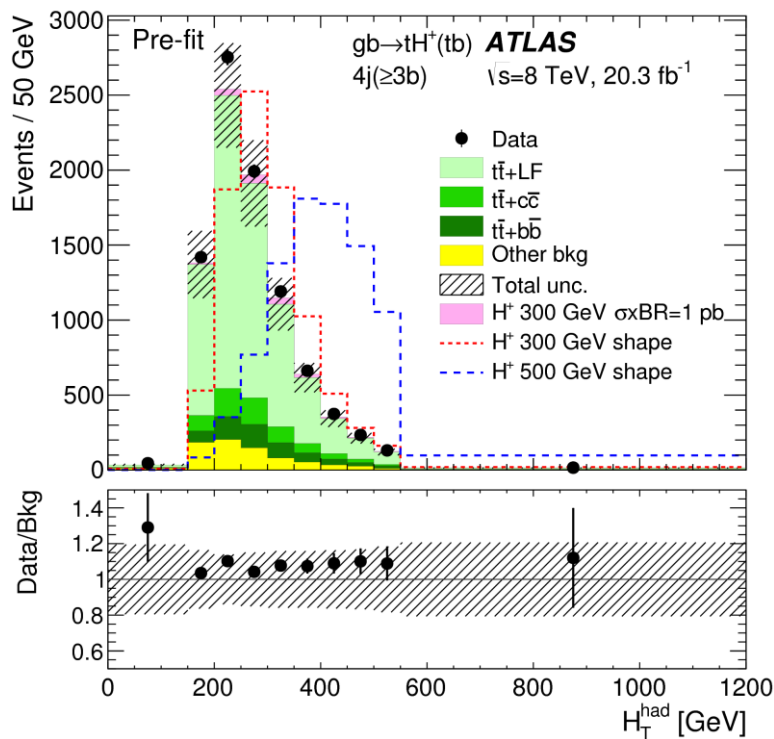
$$M = \begin{pmatrix} 4v_1^2(\lambda_1 + \lambda_3) + v_2^2\lambda_5 & (4\lambda_3 + \lambda_5)v_1v_2 \\ (4\lambda_3 + \lambda_5)v_1v_2 & 4v_2^2(\lambda_2 + \lambda_3) + v_1^2\lambda_5 \end{pmatrix}$$

5 Higgs bosons!  $h, H, A, H^\pm$

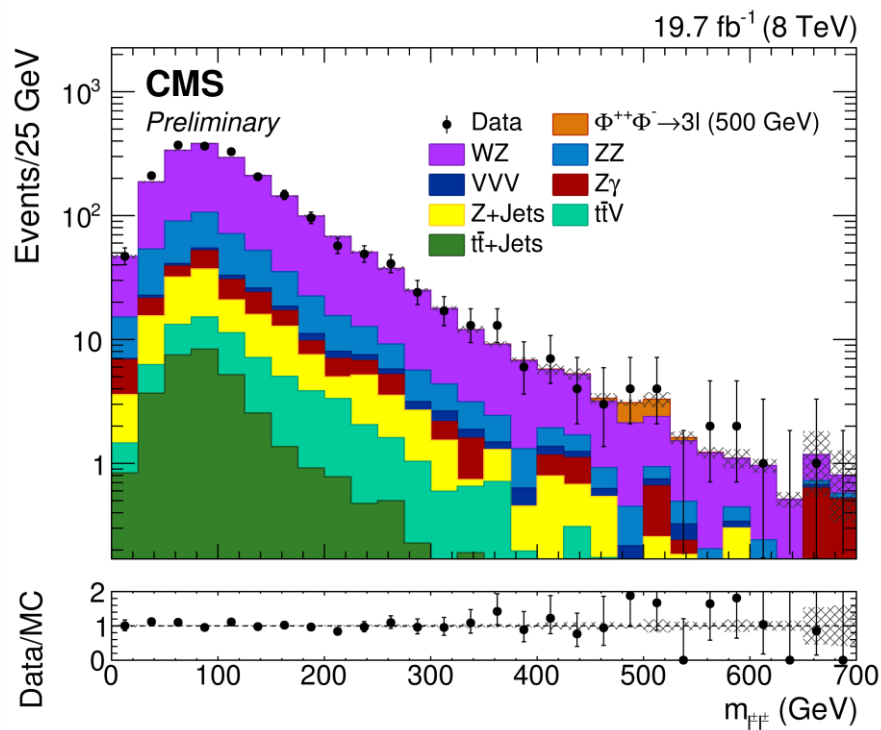


# Examples of searches for extra Higgs bosons

## Singly-charged



## Doubly-charged



Limits around  $\mathcal{O}(200 \text{ GeV})$

600 GeV

# Parity Restoration: Signals

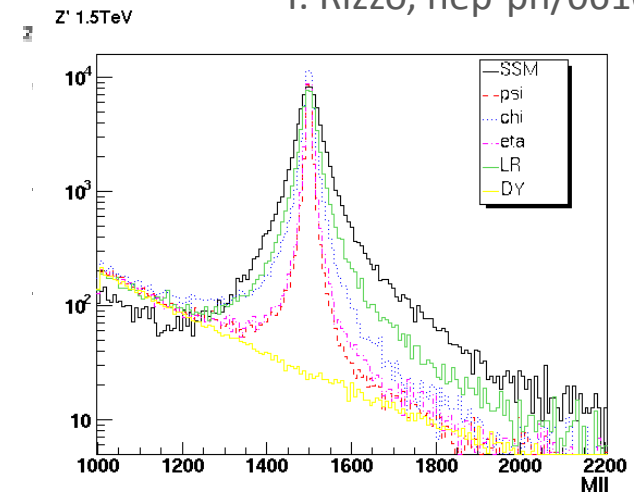
- ❖ Primary signals are (right-handed)  $W'$  (+  $Z'$ )
  - Dilepton resonances ( $Z'$ ) offer clean signals, well-understood backgrounds
    - At LHC, some concern about extrapolation of calibration from  $Z$  to very high energies
    - Electron/muon resolution improves/degrades with  $p_T$
  - $t\bar{t}$  decays visible
  - $\nu_R$  is presumably heavy,  $W'$  may not decay to leptons
    - Only dijet or diboson
    - If  $\nu_R$  lighter than  $W'/Z'$ ,  $\nu_R$  decays become important
- ❖ Note: many kinds of  $Z'$  - review by Langacker
  - $W'/Z'$  would also require new fermions...

arXiv:0801.1345

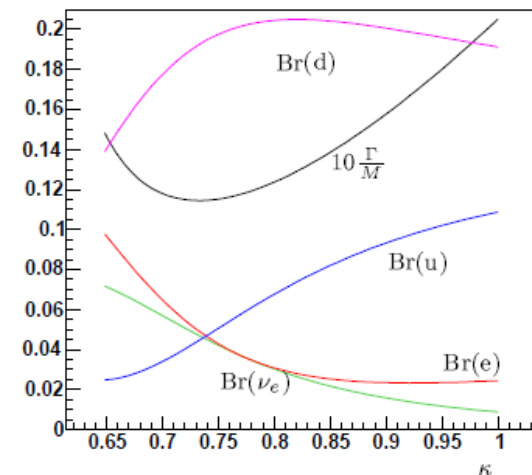
# Z' Production and Decay

T. Rizzo, hep-ph/0610104

- ❖ Production from u, d quarks is dominant at LHC
  - Couplings vary by model
  - E.g. for LR symmetric models,  $\kappa = g_R/g_L$  drives production cross-section (convolute with PDFs) and branching ratios
  
- ❖ Decays somewhat similar to Z (but almost no BR to light neutrinos, decays to top open up), plot assumes  $\nu_R$  heavier



ATL-PHYS-PUB-2005-010

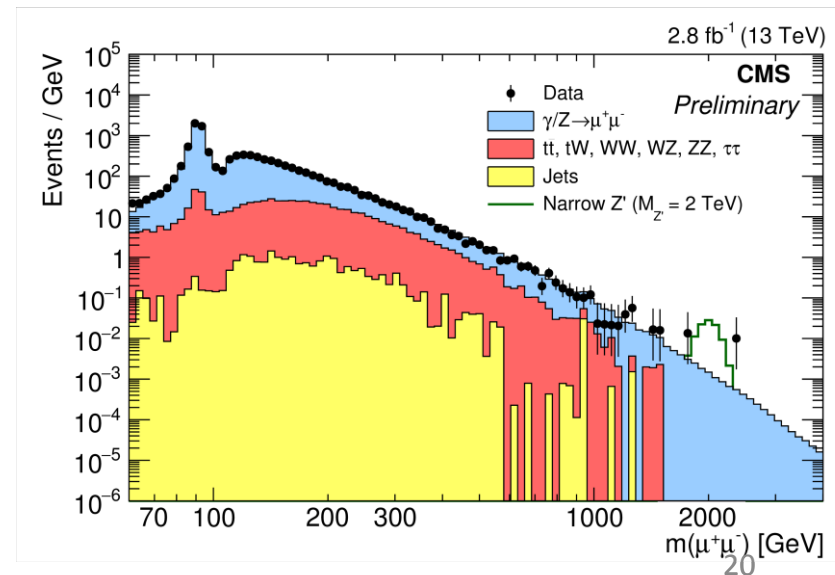
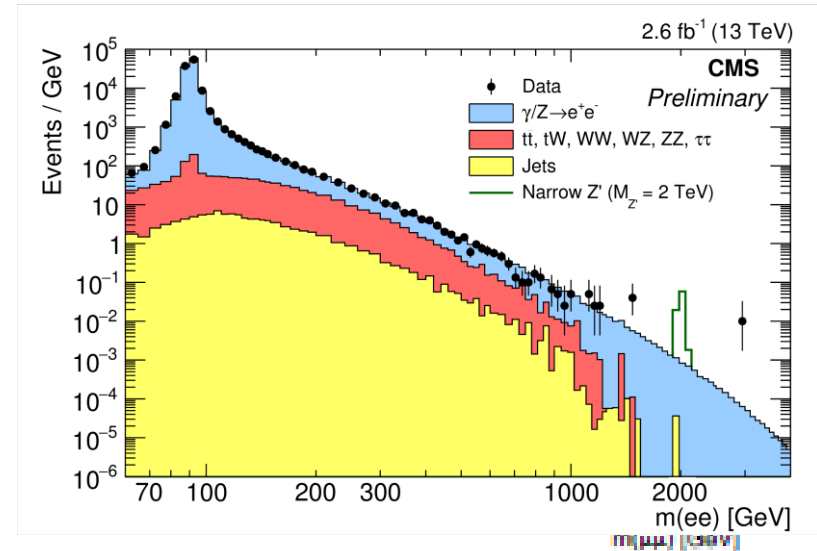


# $Z' \rightarrow ee/\mu\mu$

## ❖ Most promising channels:

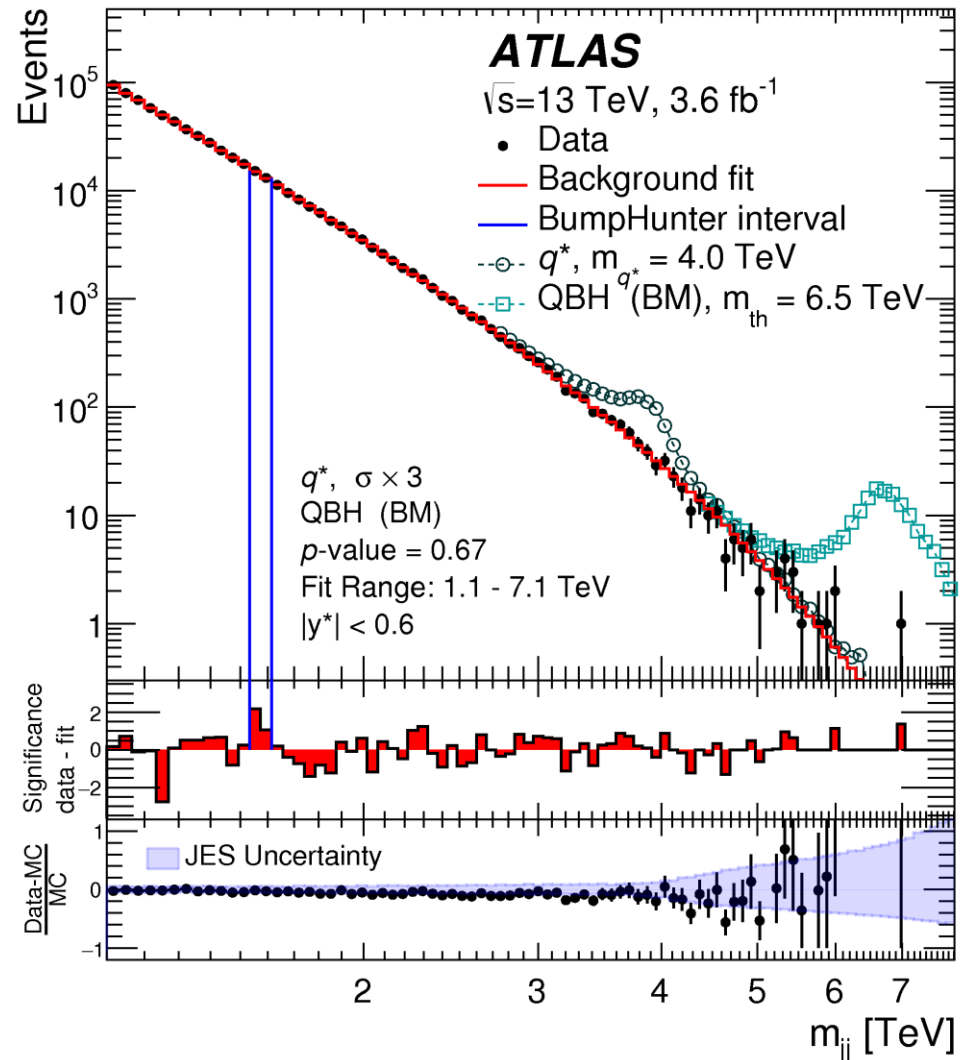
- Backgrounds very low!
- “Self-calibrating”
- In  $ee$ , at high masses, energy resolution dominated by constant term
- 10 GeV for 1.5 TeV electron
- Could measure width!

## ❖ LHC extended Tevatron reach immediately!



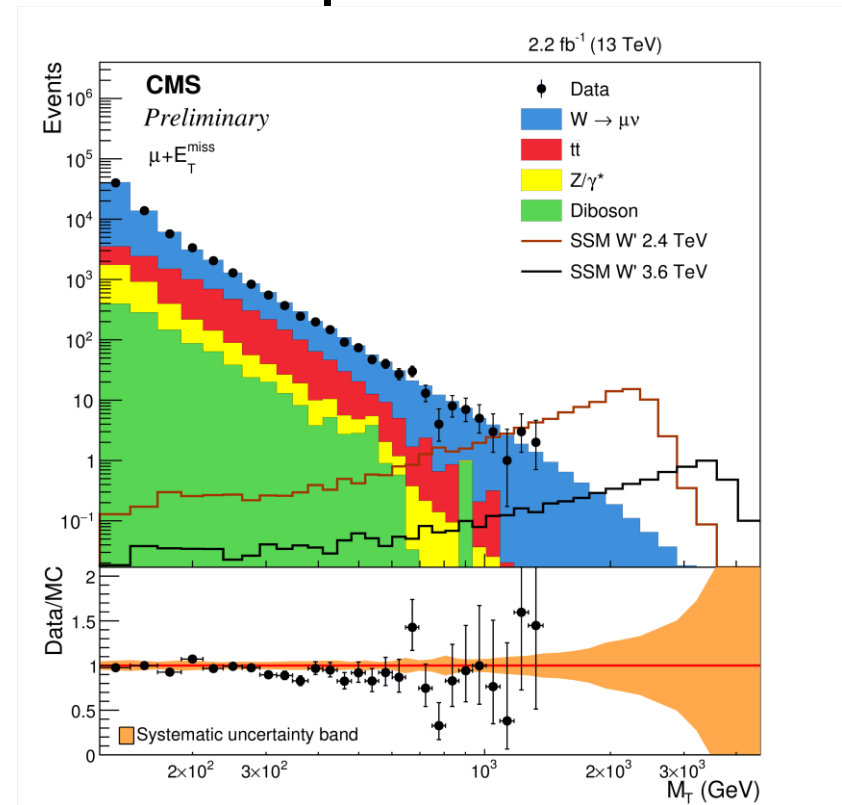
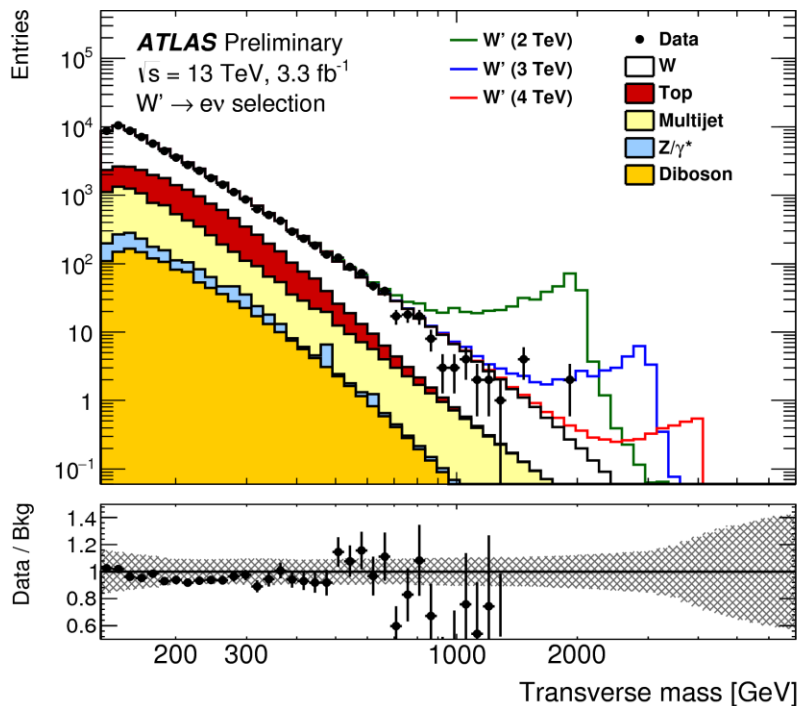
# Dijets

- SM background obviously much larger
  - But single source
  - And opens the door to strongly interacting objects



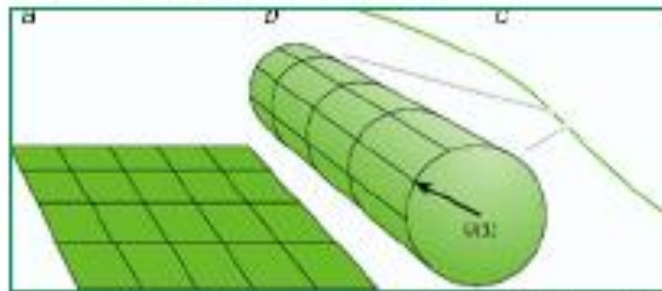
$$W' \rightarrow \mu\nu / e\nu$$

Another very simple selection: lepton + MET



# Extra Dimensions

- ❖ A promising approach to quantum gravity consists in adding extra space dimensions: string theory
  - Additional space dimensions are hidden, presumably because they are compactified



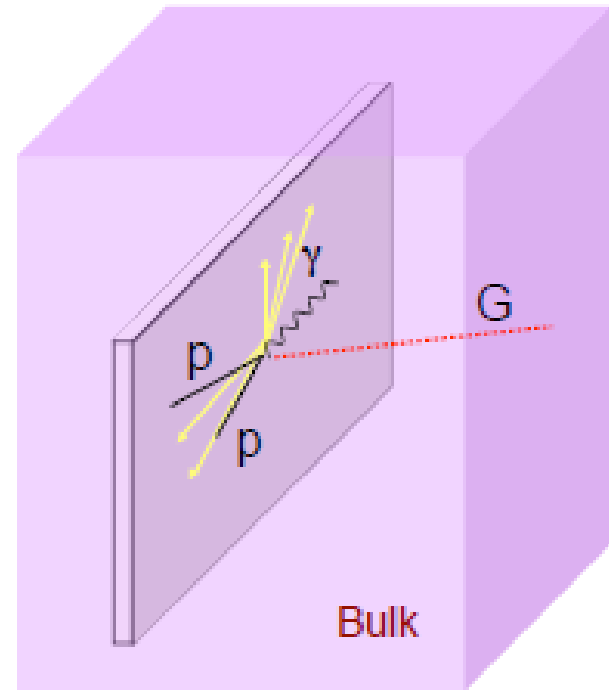
Source: PhysicsWorld

- ❖ Radius of compactification usually assumed to be at the scale of gravity, i.e.  $10^{18}$  GeV
  - In '90 Antoniadis realized they may be much larger...

Phys Lett B246:384-384, 1990

# ADD extra dimensions

- ❖ “Large extra dimension” scenario (developed by Arkani-Hamed, Dimopoulos and Dvali): Phys.Lett. B429 (1998) 263-272
  - Standard model fields are confined to a 3+1 dimensional subspace (“brane”)
  - Gravity propagates in all dimensions
  - Gravity appears weak on the brane because only felt when graviton “goes through”



Drawing by K. Loureiro



# ADD signatures

- ❖ Edges of extra dimensions identified
  - ➔ Boundary conditions
  - ➔ Momentum along extra dimension is quantized
    - Looks like mass to us
    - Very small separations → looks like continuum
    - Called Kaluza-Klein tower
- ❖ Coupling to single graviton very weak, but there are *lots* of them!
  - Large phase space → observable cross-section
    - Impacts all processes (graviton couples to energy-momentum)

❖ Consider processes that involve the bulk (i.e. gravitons)

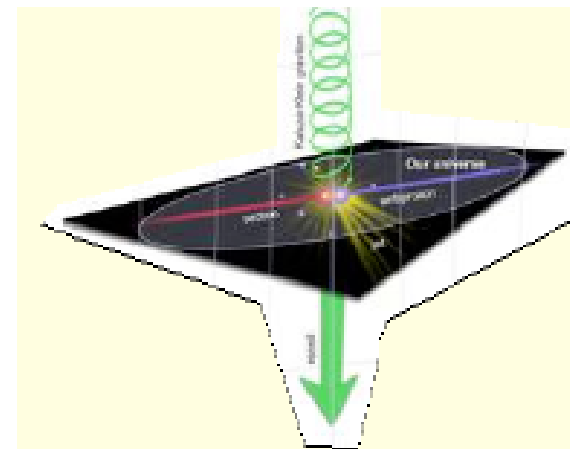
- Translational invariance is broken

➔ Momentum is not conserved ...

- ... because graviton disappears in bulk right away

❖ Look for  $p p \rightarrow \text{jet/photon} + \text{nothing}$  (i.e.  $\cancel{E}_T$ ), or deviations in high mass/angular behavior in standard model processes

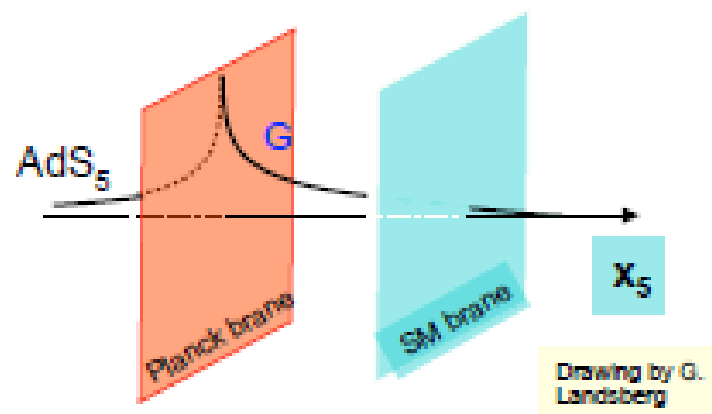
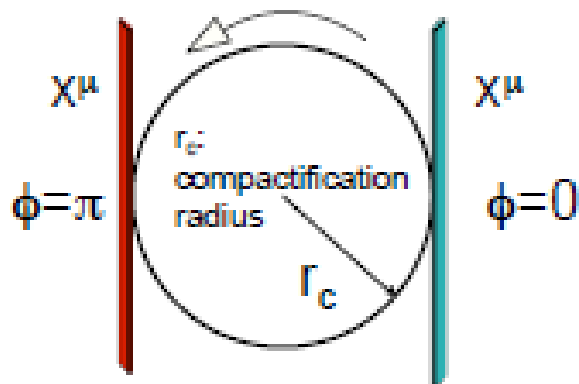
- Graviton has spin 2, couples to energy-momentum!



# Warped extra dimensions

## ❖ “Simple” Randall-Sundrum model:

- SM confined to a brane, and gravity propagating in an extra dimension
- As opposed to the original ADD scenario, the metric in the extra dimension is “warped” by a factor  $\exp(-2kr_c\phi)$
- (Requires 2 branes)



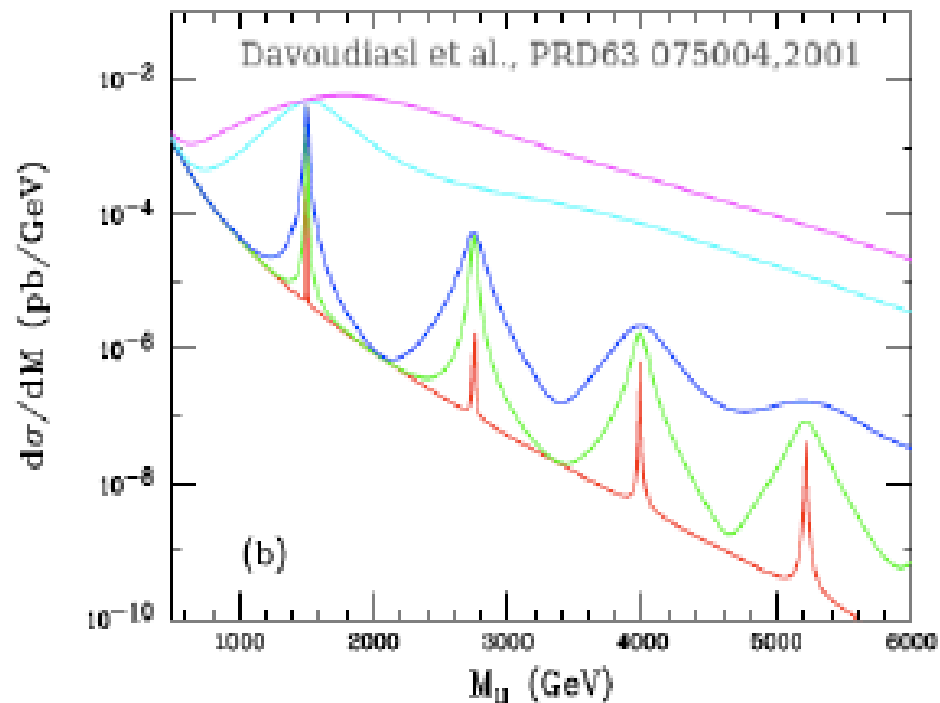
# Graviton excitations

❖ In RS, get a few massive graviton excitations

- Widths depend on warp factor  $k$
- Mass separation = zeros of Bessel function

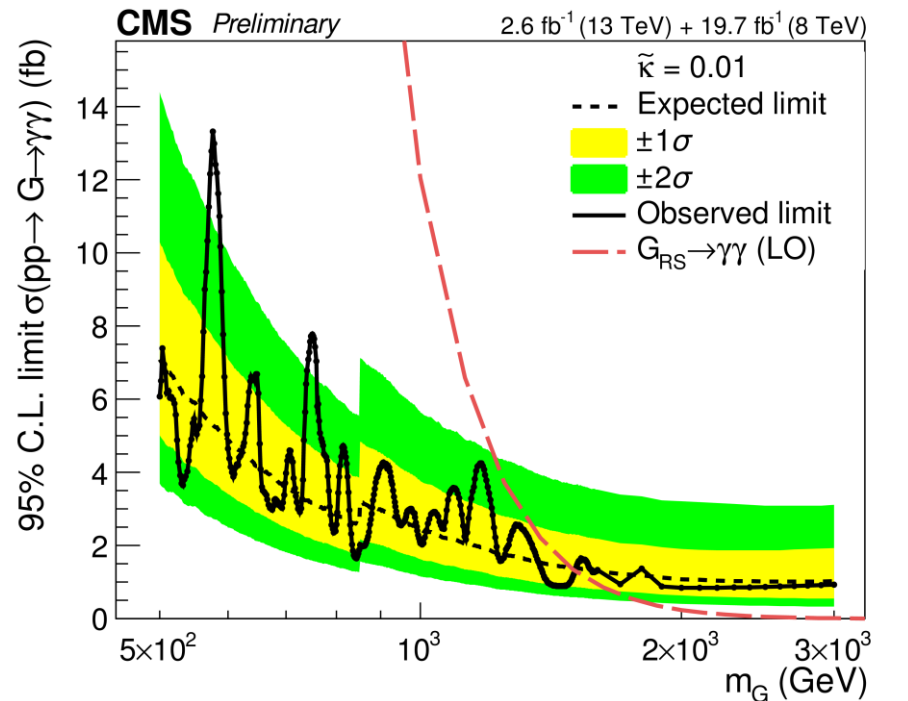
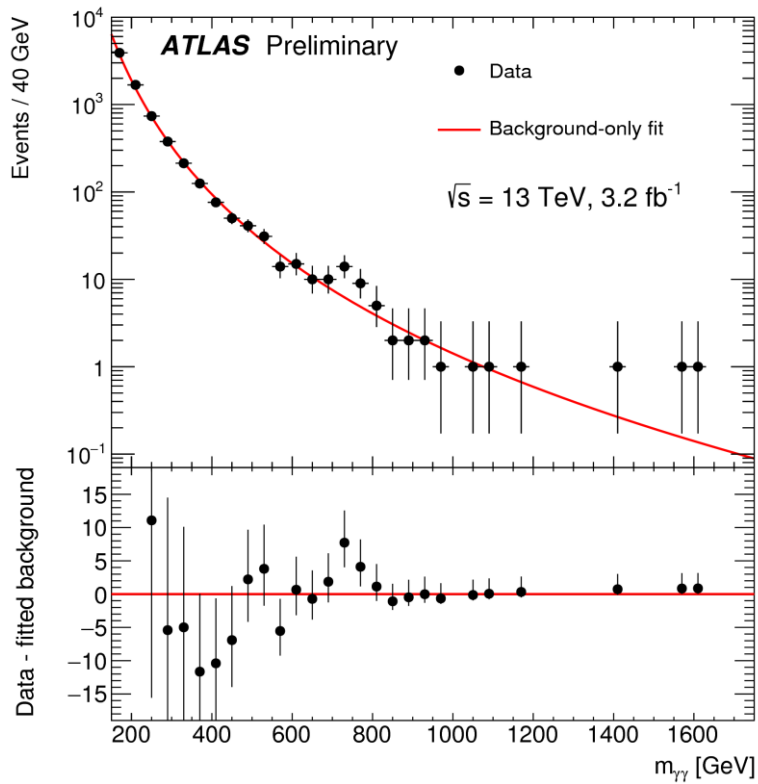
➔ Smoking gun!

(BRs also different than  $Z'$ :  
e.g.  $\gamma\gamma$  allowed)



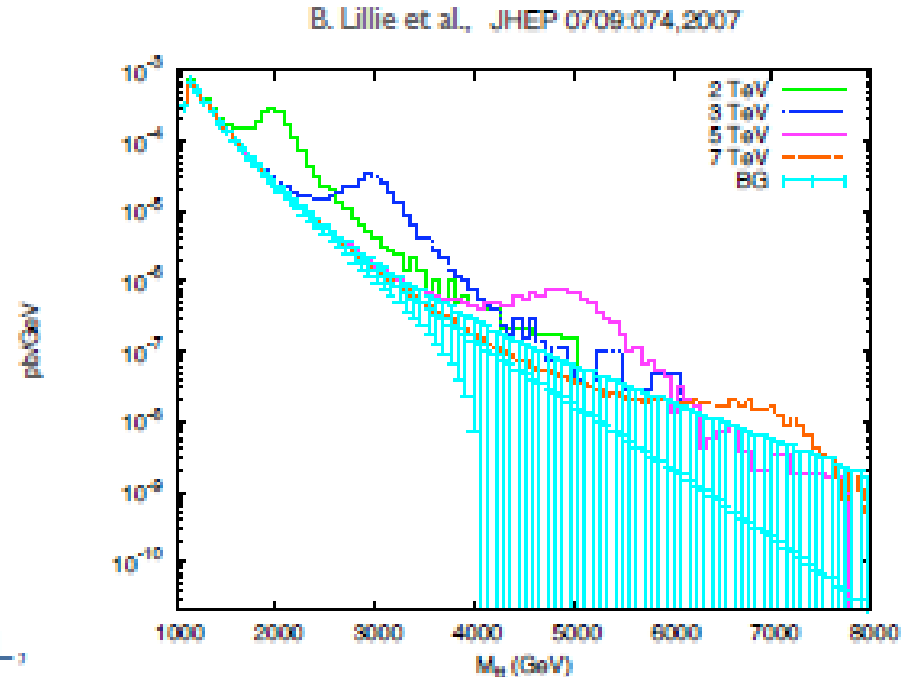
# Example

The infamous  $\gamma\gamma$  bump is an example of a search for RS gravitons:



# Gauge boson excitations

- ❖ Excitations of the gauge bosons are very promising channels for discovery
  - Couplings to light fermions are small
  - Small production cross-sections
  - Large coupling to top,  $W_L$ ,  $Z_L$
  - Look for  $t\bar{t}$ ,  $WW$ ,  $ZZ$  resonances (that can be wide)

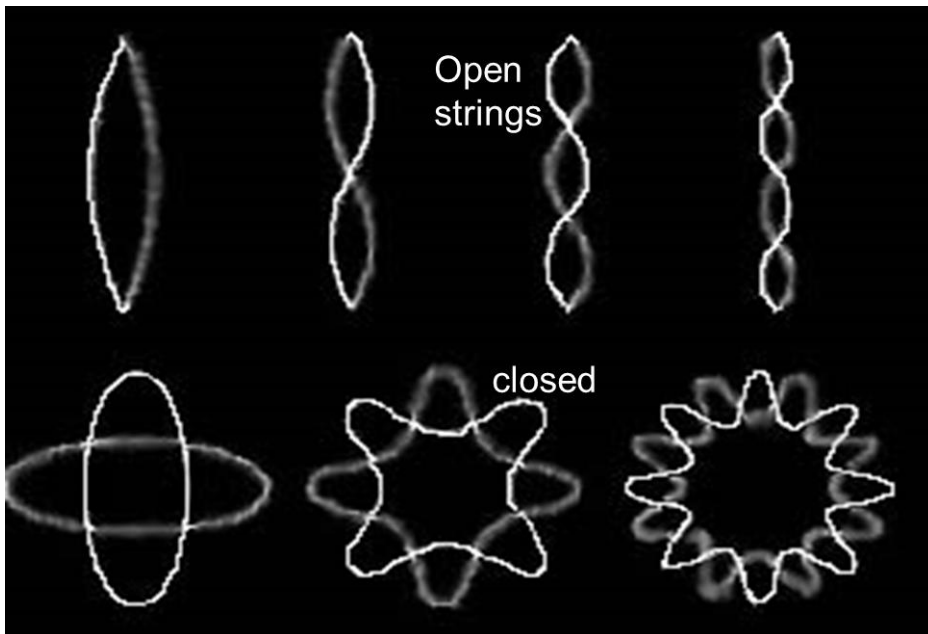
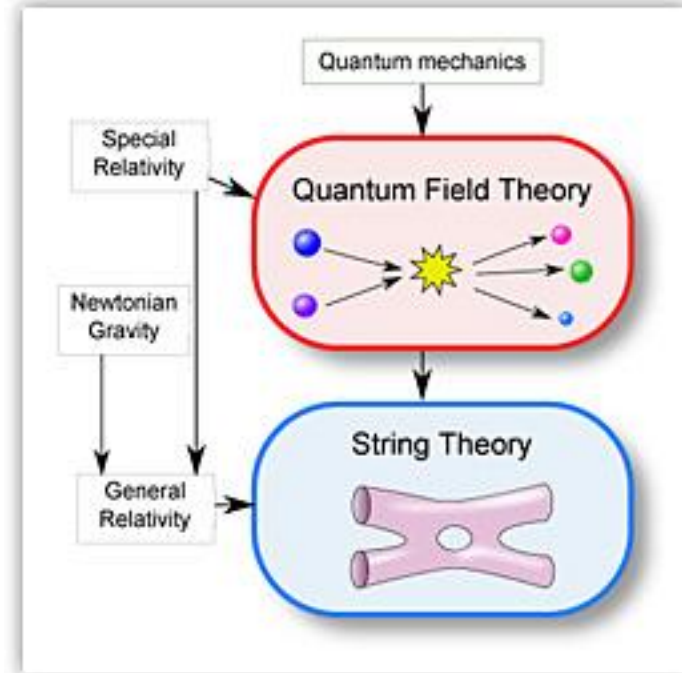


# (super)Strings

Avoid infinities from point-like particles

Different vibration modes = different particles

One fundamental parameter: *string size*



Great idea but we have not yet understood how to test it at current "low" energies

Extra dimensions a must  
Supersymmetry a plus

# Dark Matter Searches

***Lots of models, this is not a unique search!***

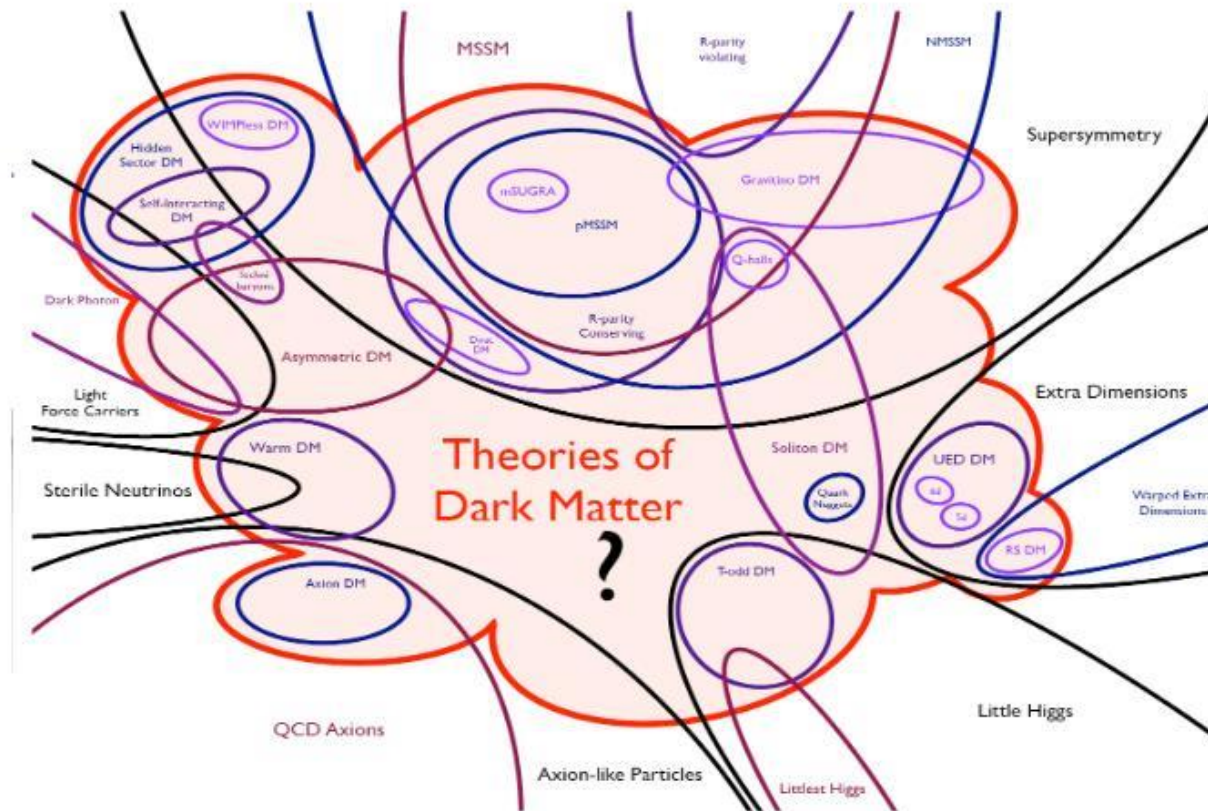
Typically divided in two "classes":

WIMP: weakly interacting massive particle  
( $\approx$  elementary particle)

MACHO: Massive Compact Halo Objects  
(planets, dwarf stars, something large)



# The energy scale(s) of new physics



T. Tait, DM@LHC '14

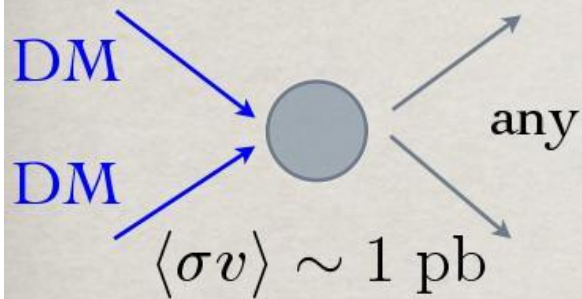
The prediction about the mass scale of DM comes with large error bars:

$$\begin{array}{ccc}
 \text{(WIMPS)} & 10^{-22} \text{ eV} < m_{DM} < 10^{20} \text{ GeV} & \text{(MACHOs)} \\
 & \text{(ALPs)} & \text{(Wimpzillas, Q-balls)}
 \end{array}$$

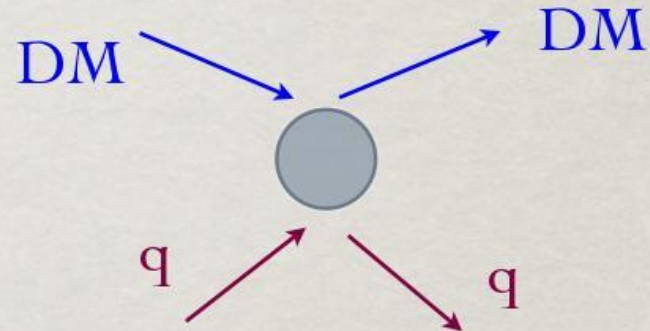
or even black holes  $\sim 10000 \times M_{\text{Sun}}$

# THE WIMP CONNECTION

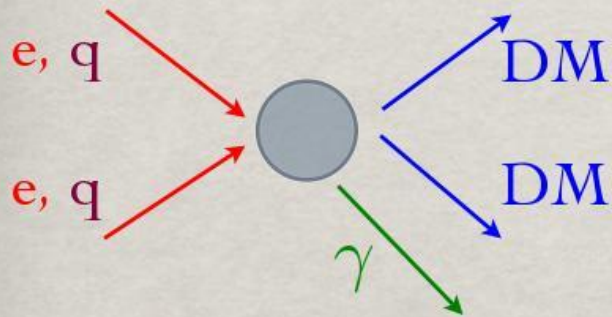
Early Universe:  $\Omega_{CDM} h^2$



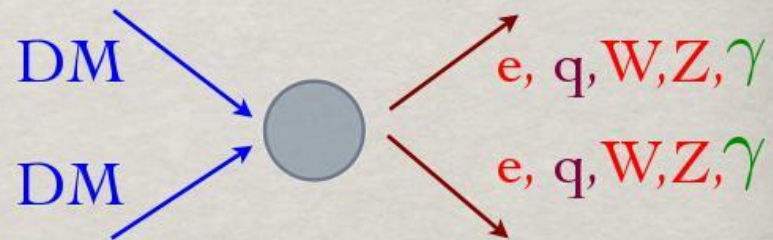
Direct Detection:



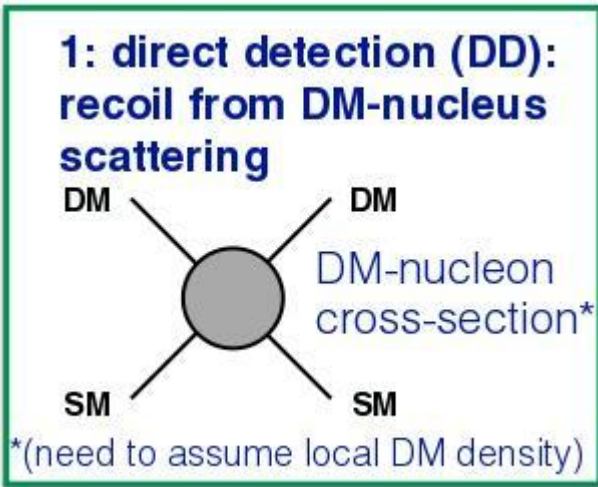
Colliders: LHC/ILC



Indirect Detection:

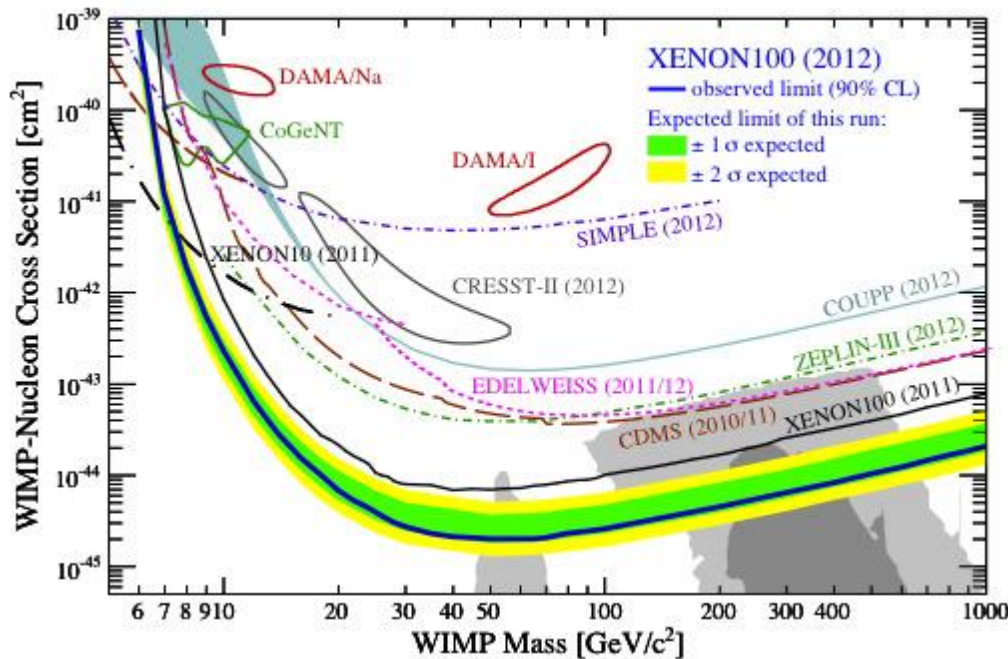
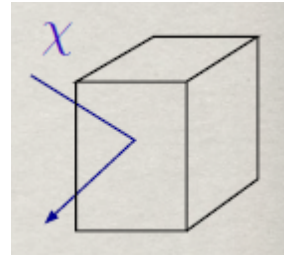


3 different ways to check this hypothesis !!!



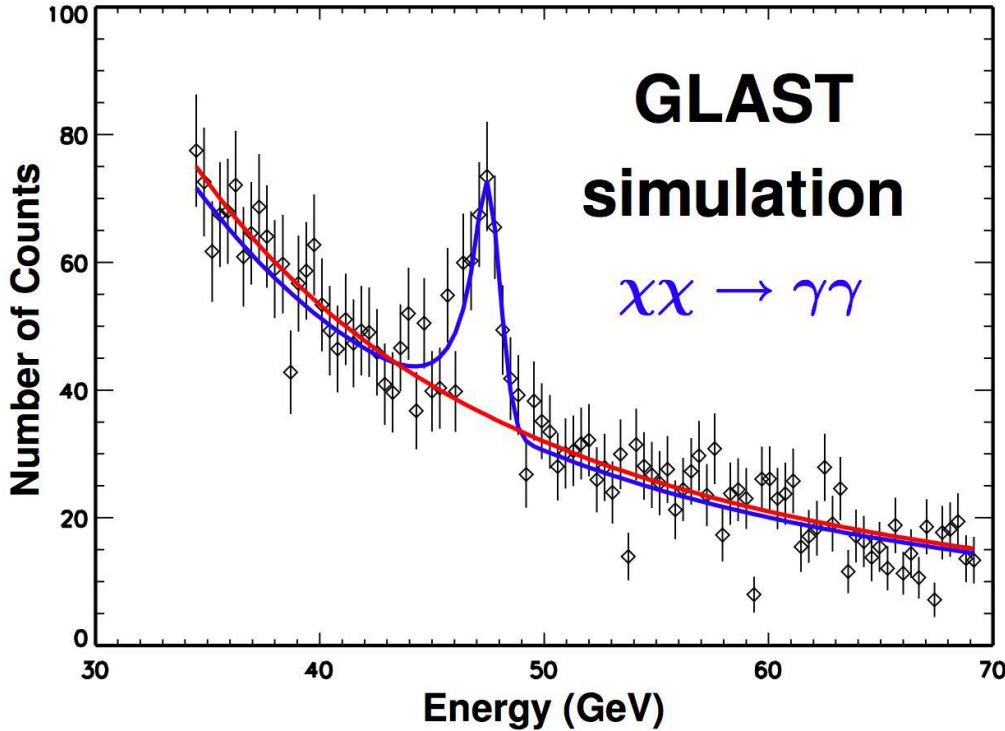
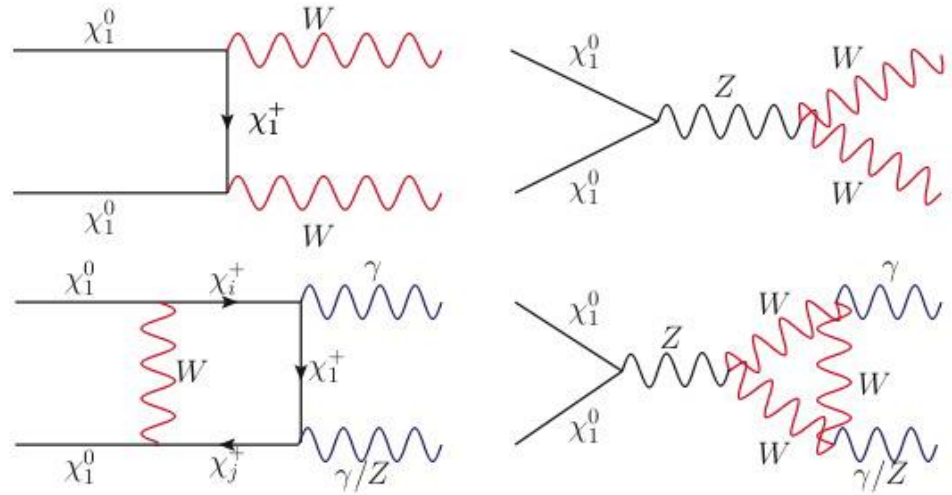
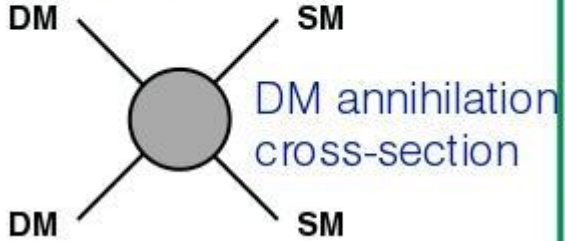
Underground searches  
(experiments: DAMA, Xenon etc

WIMP scatters off nuclei  
Looking for annual  
modulation / DM "wind"



Cross section depends  
on exchange particle:  
Z exchange ruled out  
Now looking for H  
exchange

**2: indirect detection (ID):  
DM-DM annihilation  
products**

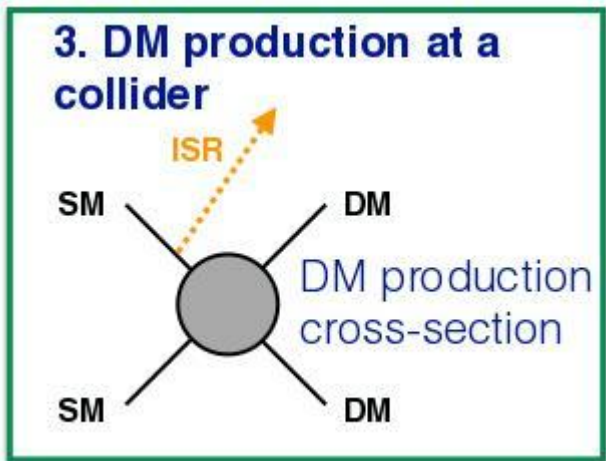


Look for annihilation signals!  
Measure decay products

Experiments: FERMI, PAMELA, AMS etc

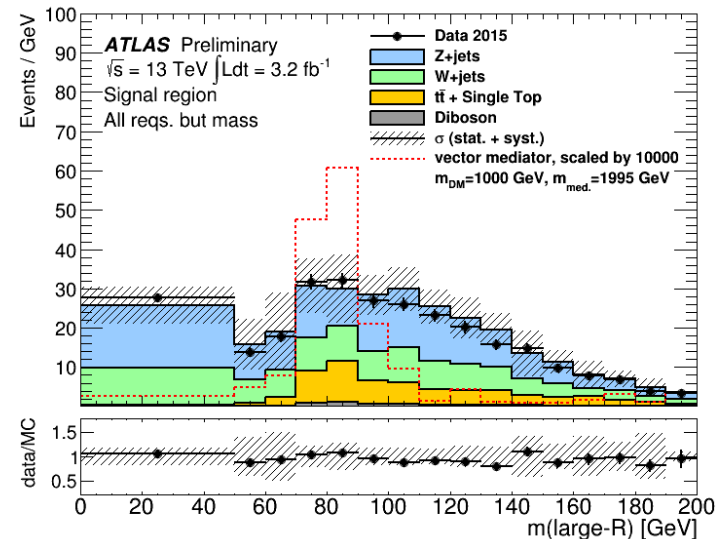
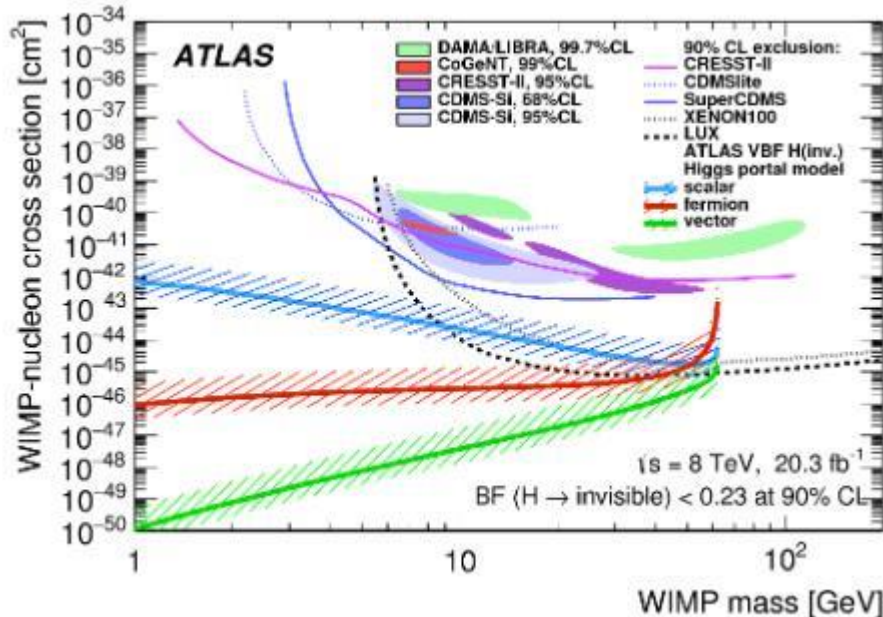
# At the LHC

## 3. DM production at a collider



No Dark Matter interaction with the detector  $\Rightarrow$  signature is missing energy

Use initial state radiation (ISR) to detect it! (e.g jets,  $\gamma$ , W, Z, H)



# Summary/outlook

- Many problems with current Standard model
  - Many new models to take over
- Some important models not mentioned, for instance:
- GUT models
  - Technicolor
  - Hidden valleys
- The LHC energy scale is tuned to be sensitive to many of these, complementary to other current searches
  - Several potential signatures requires new "objects", ie lepton-jets, long-lived heavy particles, "quirks" etc
  - Several good ideas but Nature decides which (if any) are true!