

FYST17 Lecture 11

BSM I

Thanks to G. Brooijmans, C. Grojean

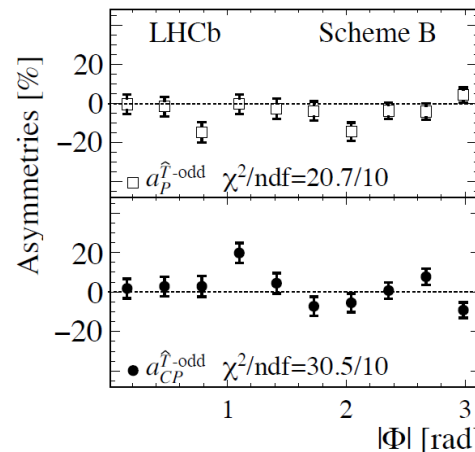
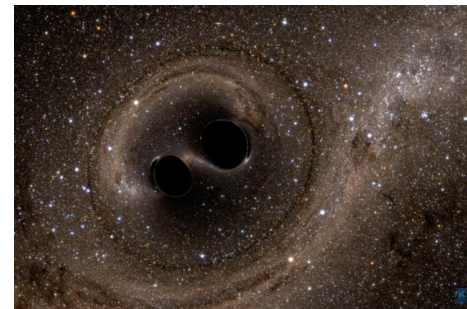
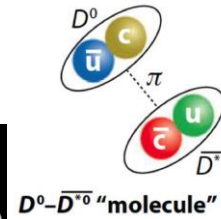
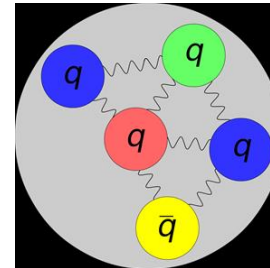
Today & Monday

- Why go Beyond the SM?
 - What are the problems with the SM?
 - What direct measurements points to physics BSM
- Some attempts at solutions
 - Supersymmetry
 - Extended Higgs sector
 - Extra dimensions
 - A few others
- Searches for DM, gravitational waves

Any direct evidence?

Certainly a few measurements that are not incorporated in the current Standard Model:

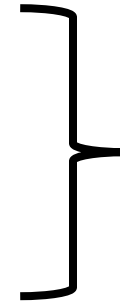
- Exotic baryons (X, pentaquarks etc)
- Neutrino masses!
- (Gravitational waves)
- The new LHCb CP violation measurement (although $<4\sigma$)



Status of the Standard Model

19 parameters (+ ν masses)

Tested to precision level $10^{-3} - 10^{-12}$



Extremely
successful!

But empirically incomplete

Structure quite complicated

Aesthetically unacceptable

Many problems with naturalness

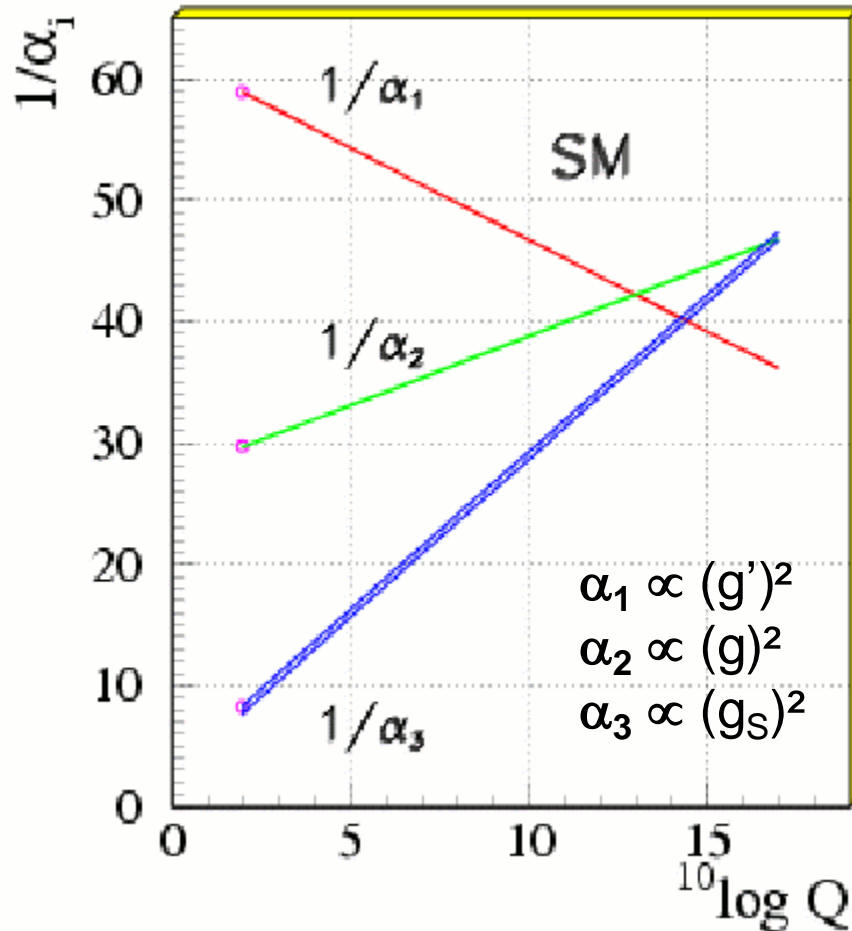
No quantum gravity

Missing answers to "big" questions

Examples of answers we need

- *What is the origin of CP violation?*
- *What is the origin of the matter/anti-matter asymmetry*
- *Why three gauge forces (so far)? And three generations?*
- *Why is the strong interaction strong? Why only left-handed particles participate in weak force?*
- *Gravity? Is there a unified description of all forces?*
- *Why is $\text{mass}(W/Z/H) \ll \text{mass}(\text{Planck})$? (Hierarchy problem)*
- *Why is charge quantized?*
- *What is Dark Matter and Dark Energy? (and why Dark Energy now?)*
- *What was the Big Bang?*

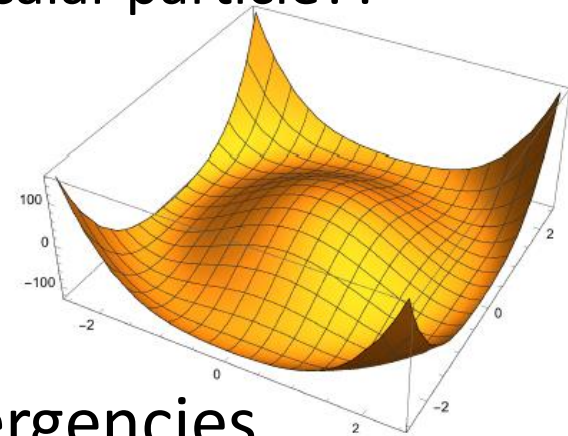
Unification of coupling constants?



Extrapolating the Standard Model coupling constants to higher energies

The Higgs discovery just adds to that list...

- What is it, really, a condensate in our Universe?
- Is it elementary?
 - If yes, why is there only 1 fundamental scalar particle??
- Why does it have mass² $\mu^2 < 0$?!
- Higgs mechanism gives quadratic divergencies
 - (see later)



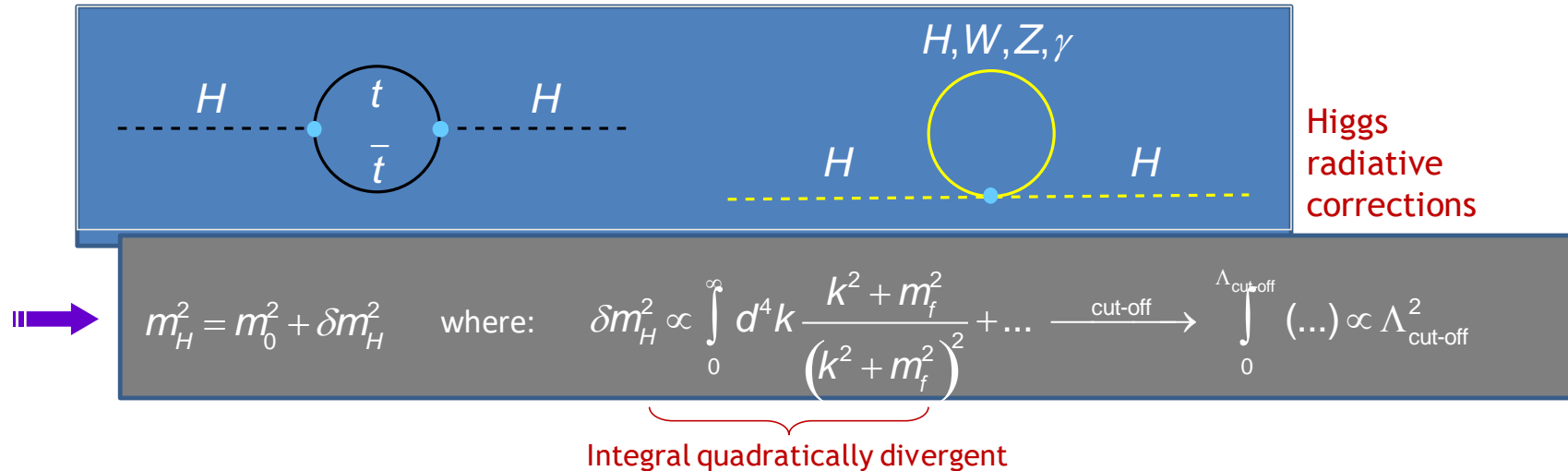
Is the Standard Model really fundamental?

- Does not appear so ($\gtrsim 25$ parameters?!)
- Evidence of selective processes:
 - For instance, no neutral colored fermions
 - $q_d = q_e / N(\text{colors}) \Rightarrow$ grand unification?
- **Fragile:** small changes in parameters \Rightarrow very different physics!
 - If $m_d < m_u$: all protons decay \Rightarrow no atoms
 - If $m_e > 4m_p - m_\alpha \Rightarrow$ Sun doesn't burn \Rightarrow no us
 - If $v \gg \text{TeV} \Rightarrow |m_n - m_p|$ large, rapid neutron decay \Rightarrow no chemistry nor life

The “Gauge Hierarchy Problem”

Discover of Higgs boson with mass < 1 TeV means the Standard Model is complete !

However, when computing radiative corrections to the bare Higgs mass a problem occurs:



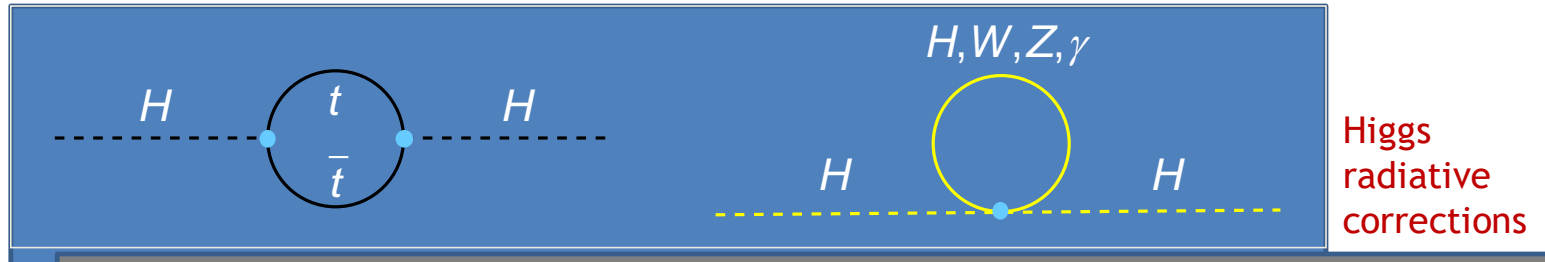
The cut-off sets the scale where new particles and physical laws must come in
 Above the EW scale we only know of two scales: GUT ($\sim 10^{16}$ GeV) and Planck ($\sim 10^{19}$ GeV)
 Such a cut-off would require an incredible amount of finetuning to keep m_H light

$$m_H^2 = (125 \text{ GeV})^2 = m_0^2 + C \cdot \Lambda_{\text{cut-off}}^2$$

The “Gauge Hierarchy Problem”

Discover of Higgs boson with mass $< 1 \text{ TeV}$ means the Standard Model is complete !

However, when computing radiative corrections to the bare Higgs mass a problem occurs:



$m_H^2 = m_0^2 + \delta m_H^2$ where:

$$\delta m_H^2 \propto \int_0^\infty d^4k \frac{k^2 + m_f^2}{(k^2 + m_f^2)^2} + \dots \xrightarrow{\text{cut-off}} \int_0^{\Lambda_{\text{cut-off}}} (\dots) \propto \Lambda_{\text{cut-off}}^2$$

Integral quadratically divergent

Missing protection of scalar Higgs mass is related to **absence of a symmetry principle**. Setting $m_H = 0$ in SM Lagrangian, **does not restore any symmetry in the model**.

New physics models should address this. M_H should become a deviation from some exact symmetry, and is thus **intrinsically small** !

$$m_H^2 = (125 \text{ GeV})^2 = m_0^2 + C \cdot \Lambda_{\text{cut-off}}^2$$

Hunting for Answers

- ❖ Get more information
 - Measure particles and their interactions in detail
 - Precision measurements (e.g. LHCb)
 - Observe new particles or interactions
 - Search in new areas in “phase space”
- ❖ Find the underlying pattern(s)
 - Hypothesize, build models
 - Internally consistent? Consistent with data?
 - Suggestions on where to look

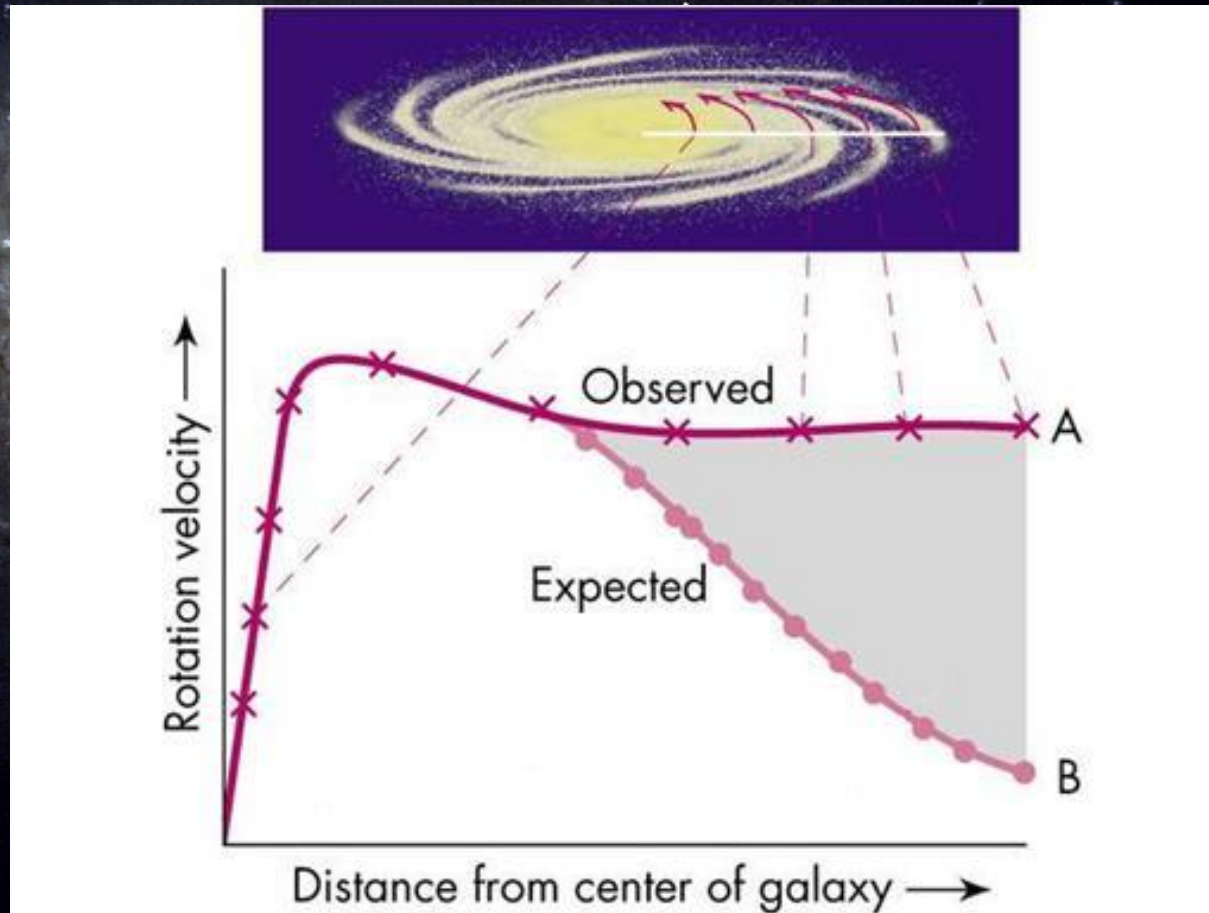
Experiment

Theory

Where to Start?

- ❖ BSM physics **must** couple to SM (weakly?), but is it
 - Resonant?
 - Does it have new massive particles decaying to electrons, muons, quarks, W, Z, ...?
 - “SM-like”?
 - Same but includes some new long-lived particles in the decay chain... (e.g. dark matter candidate)
 - No new “particles” in reach
 - Hidden or too heavy or... don't exist
 - Are there new interactions?

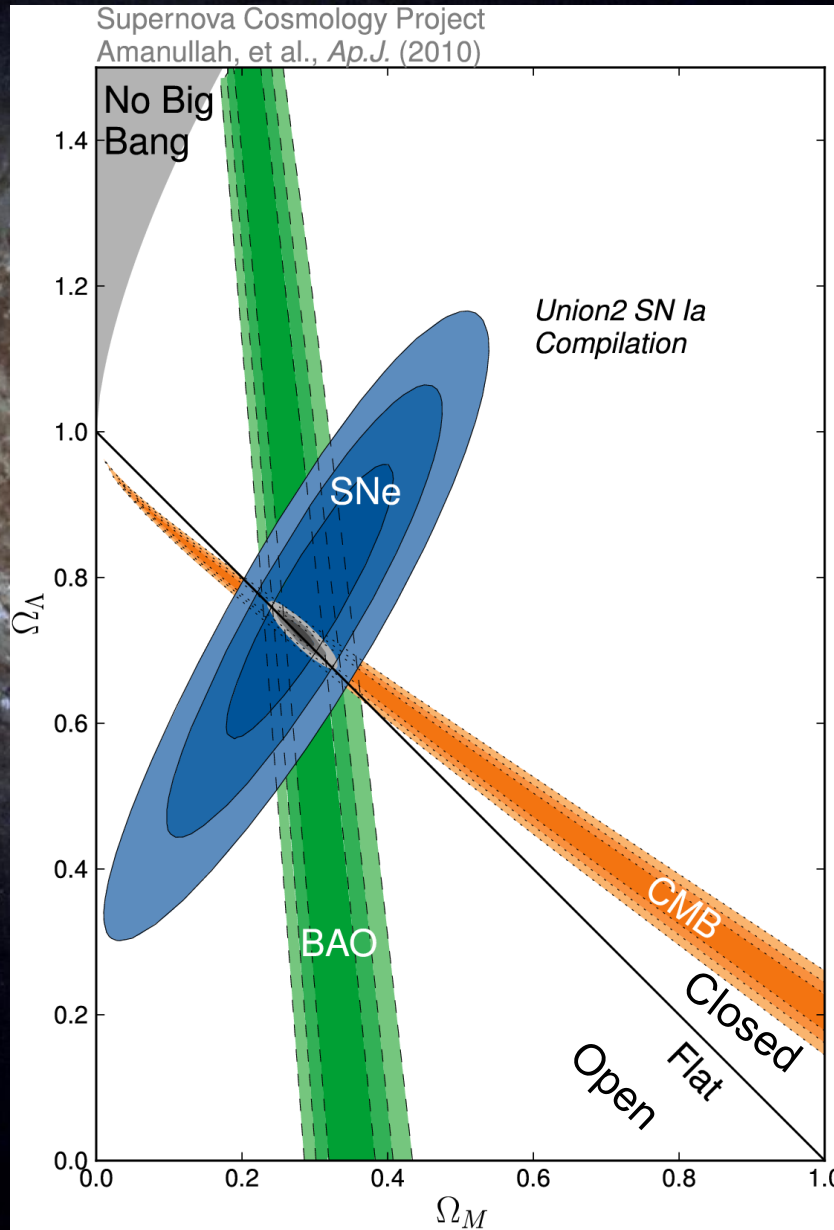
Galaxy rotation curves



**Standard Model only accounts for
~20% of the matter of the Universe!!!**

Supernovae data

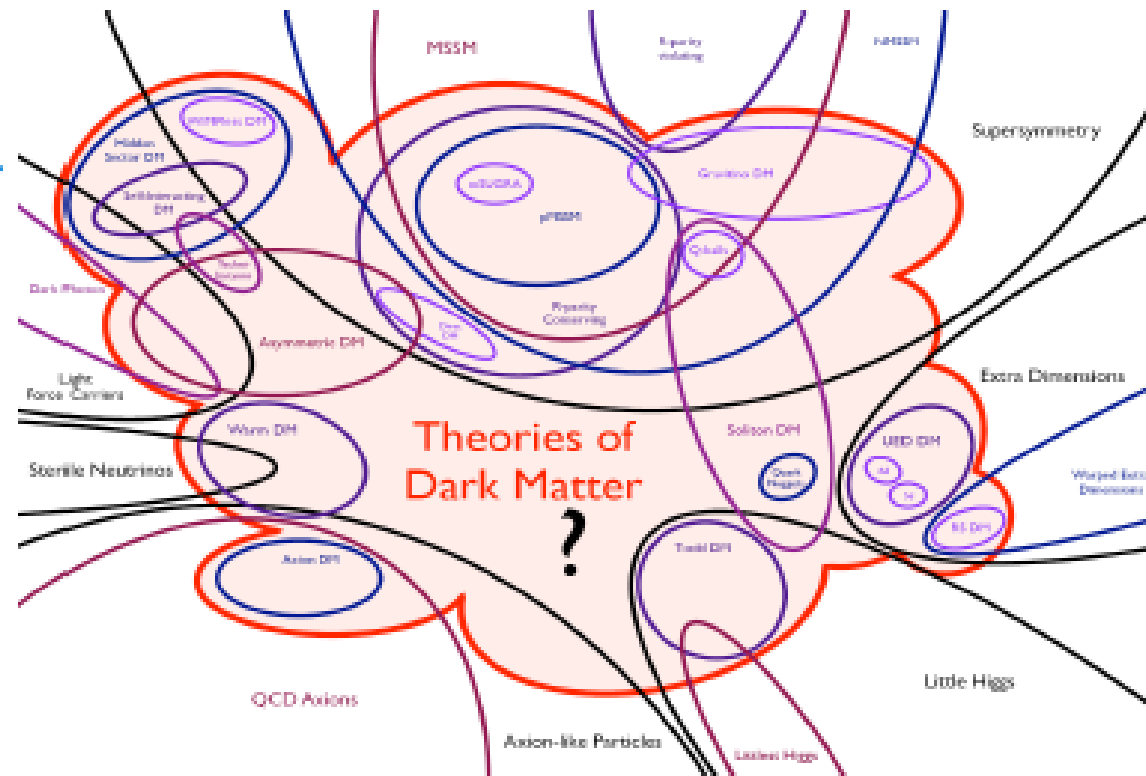
“Cosmological constant” term



Matter only
accounts for
~30% of the
Universe!

“Matter density” term

The energy scale(s) of new physics



T. Tait, DM@LHC '14

The prediction about the mass scale of DM comes with large error bars:

$$10^{-22} \text{ eV} < m_{DM} < 10^{20} \text{ GeV}$$

(ALPs) (Wimpzillas, Q-balls)

Supersymmetry (SUSY)

Idea

New symmetry *fermions* \leftrightarrow *bosons*

This symmetry is the most general extension of Lorentz invariance

SUSY has: $N_{\text{dof}}(\text{bosons}) = N_{\text{dof}}(\text{fermions})$
[cf. SM: $N_{\text{dof}}(\text{bosons}) \ll N_{\text{dof}}(\text{fermions})$]

Spin 0	Spin 1/2	Spin 1	Spin 3/2	Spin 2
sLeptons	Leptons		Gravitino	Graviton
sQuarks	Quarks			
Higgs	Higgsino			
	Photino	Photon		
	Zino	Z		
	Wino	W		
	Gluino	Gluon		

- To create *supermultiplets*, we need to add one *superpartner* to each SM particle
- Superpartners have opposite spin statistics but otherwise equal quantum numbers
- Need to introduce an additional Higgs doublet to the non-SUSY side \rightarrow 5 Higgs bosons

But where are these partners?!

Supersymmetry must be broken (if realized)

Particle spectrum (minimal!)

In reality the new states would mix

Several ideas of how the supersymmetry is broken – intimately connected with EWK symmetry breaking

		Spin 0	Spin 1/2	Spin 1
<div style="border: 1px solid black; padding: 2px; display: inline-block; margin-bottom: 5px;">SM</div> <div style="border: 1px solid black; padding: 2px; display: inline-block; margin-bottom: 5px;">SUSY</div>	Eigenstates of mass	$\tilde{\ell}_1, \tilde{\ell}_2$	ℓ	
		\tilde{q}_1, \tilde{q}_2	q	
		h^0, H^0, A^0, H^\pm	$\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$	
			$\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$	γ, Z^0, W^\pm
		\tilde{g}_a	g_a	

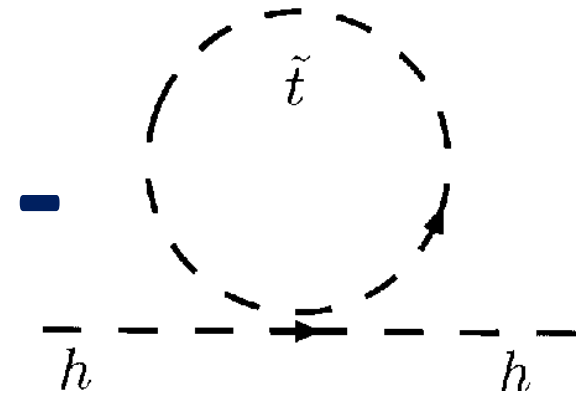
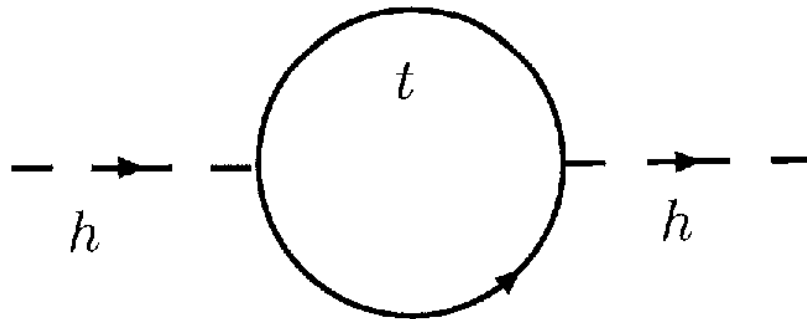
Squark/slepton mixing proportional to SM partner masses
 → largest for 3rd gen.
 → can become lightest squarks / sleptons

The gauge-mixed physical states that propagate in space and time and that can be observed.
 Neutralinos: mass eigenstates of photinos, zinos, neutral higgsinos
 Charginos : mass eigenstates of winos and charged higgsinos

Since we don't know the mechanism, have to introduce $\mathcal{O}(100)$ new parameters

SUSY and the hierarchy problem

If Supersymmetry not broken we would have perfect cancellation in the loops!



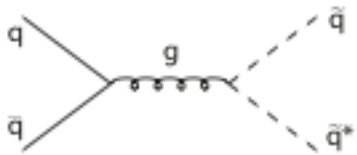
But as $m(\tilde{t}) \neq m(t)$ they do not quite cancel, instead just a suppression

This still gives a decent result if $|m(\text{fermion}) - m(\text{boson})| < \mathcal{O}(\text{TeV})$

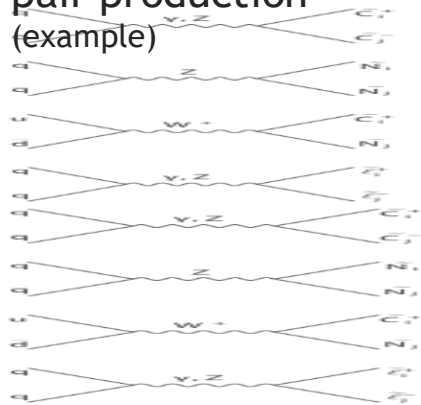
Once mass spectrum fixed, all cross sections predicted

Spin structure of SUSY spectrum: lower σ than other BSM models, harder to find !

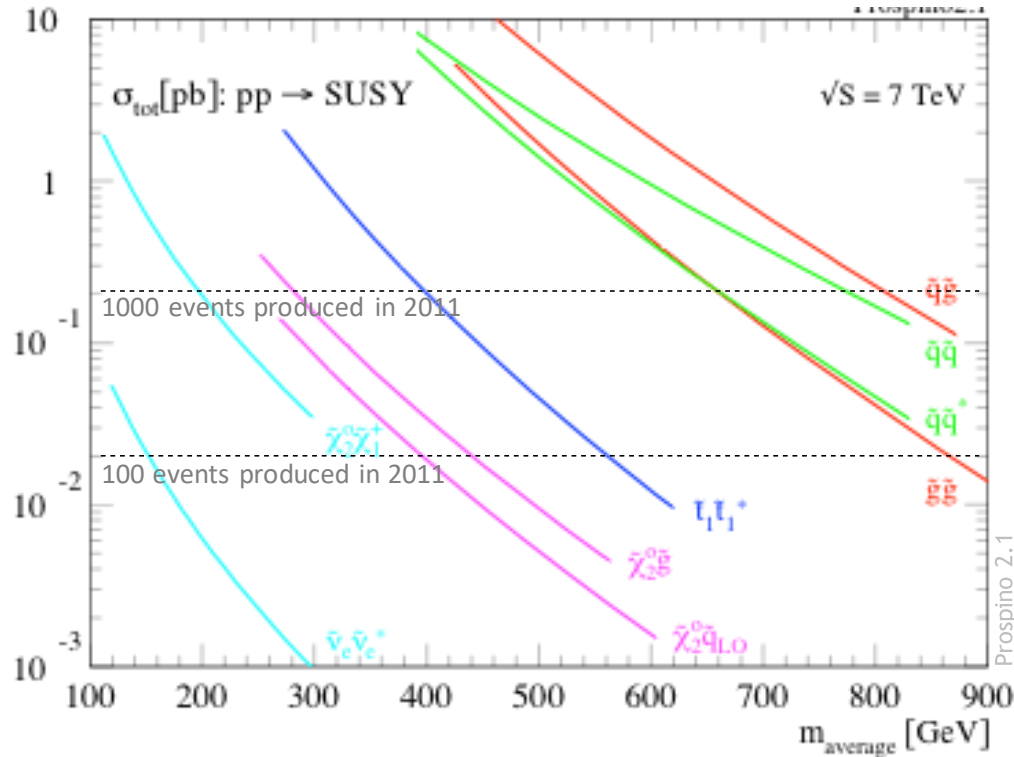
Direct squark pair production (example)



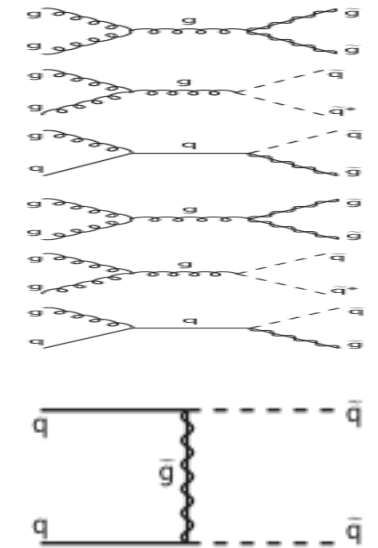
Direct gaugino/slepton pair production (example)



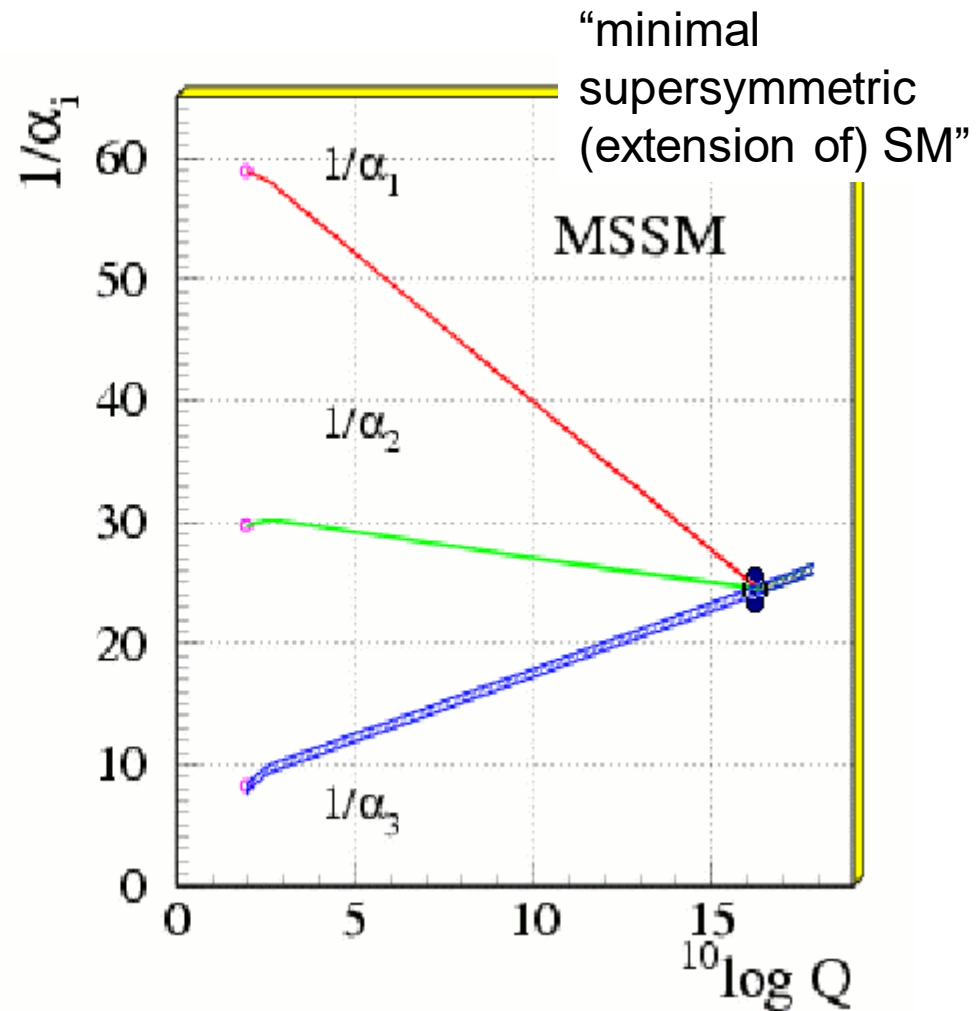
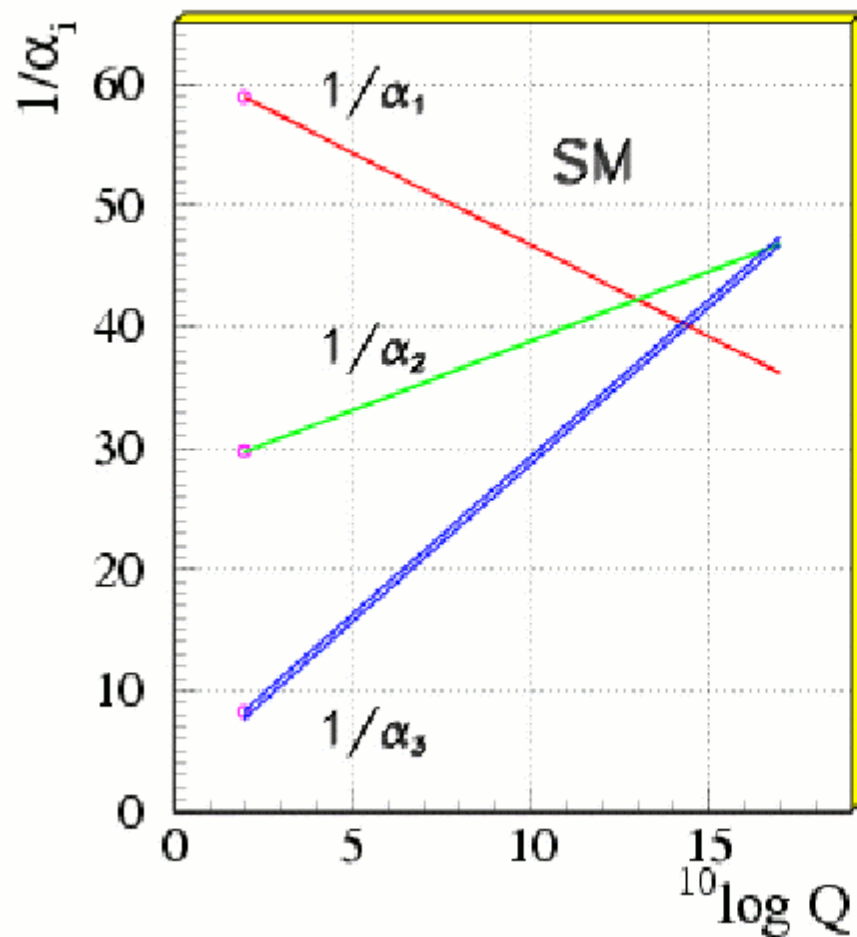
SUSY cross section versus sparticle mass



Gluino & squark production (examples)



Unification of coupling constants with supersymmetry

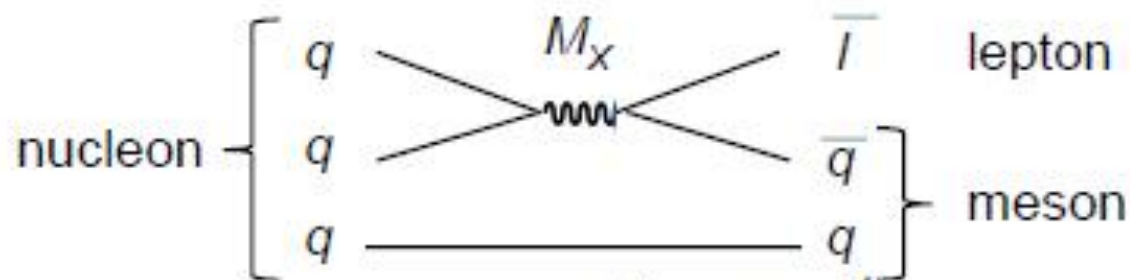


Proton Decay


(G. Giudice SSLP'15)

in GUT, matter is unstable

decay of proton mediated by new SU(5)/SO(10) gauge bosons



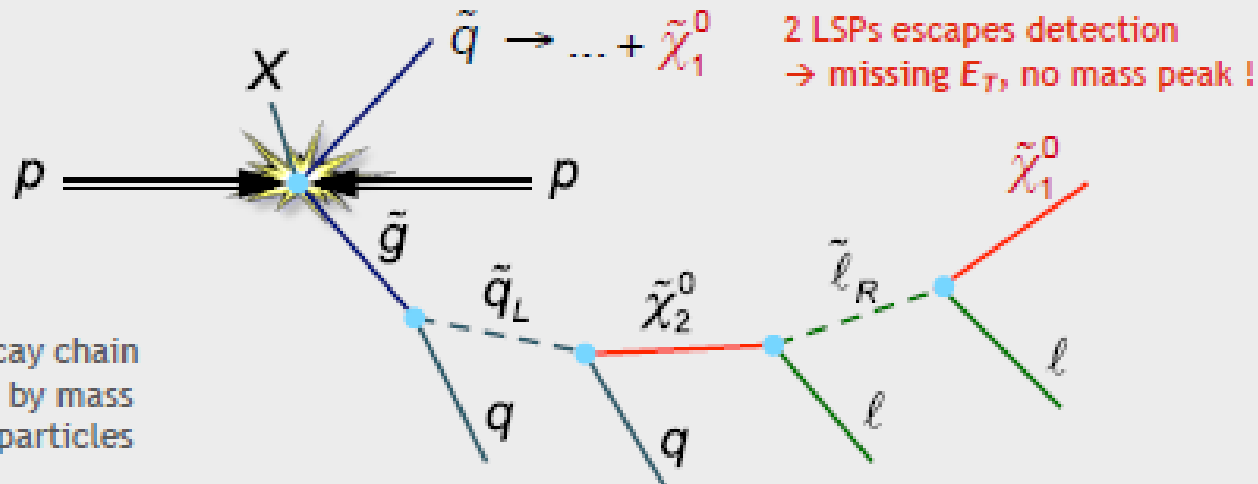
$$\text{GUT: } \tau_p(p \rightarrow e^+ \pi^0) = \left(\frac{M_X}{10^{15} \text{ GeV}} \right)^4 10^{31-32} \text{ yr}$$



$$\text{Exp: } \tau_p(p \rightarrow e^+ \pi^0) > 8.2 \times 10^{33} \text{ yr}$$

Characteristic SUSY Decay Cascades

- To avoid proton decay, a new conserved quantum number (R) is introduced, which forces a SUSY particle to decay in at least one other SUSY particle
- The lightest SUSY particle is thus stable (LSP), and must be neutral and colourless \rightarrow WIMP (dark matter candidate)
- Typical **LSP is spin- $\frac{1}{2}$ neutralino**. It could also be a gravitino
- With R parity: SUSY production in pairs only \rightarrow requires energy $2 \times$ SUSY mass !

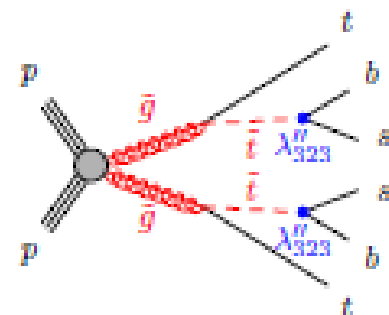
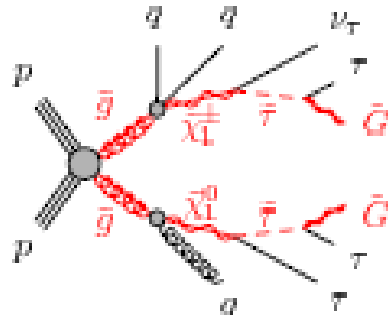
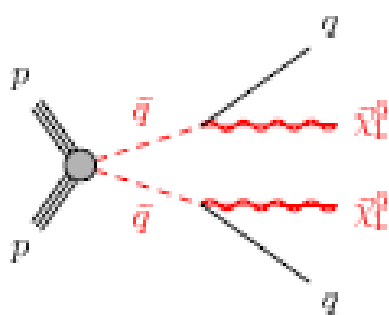


"Typical" SUSY decay chain at the LHC, driven by mass hierarchy of SUSY particles

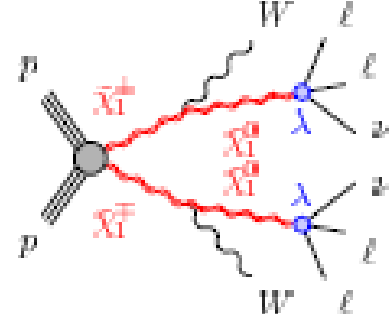
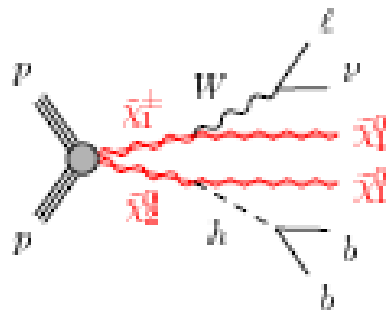
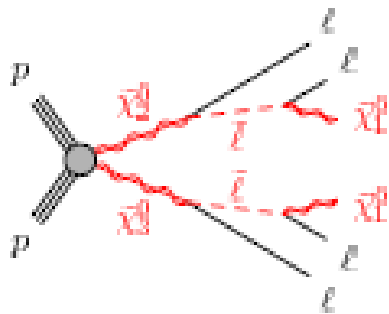
Canonical SUSY

❖ Wide range of signatures

- Strong production... (large cross-section)



- ... or weak production



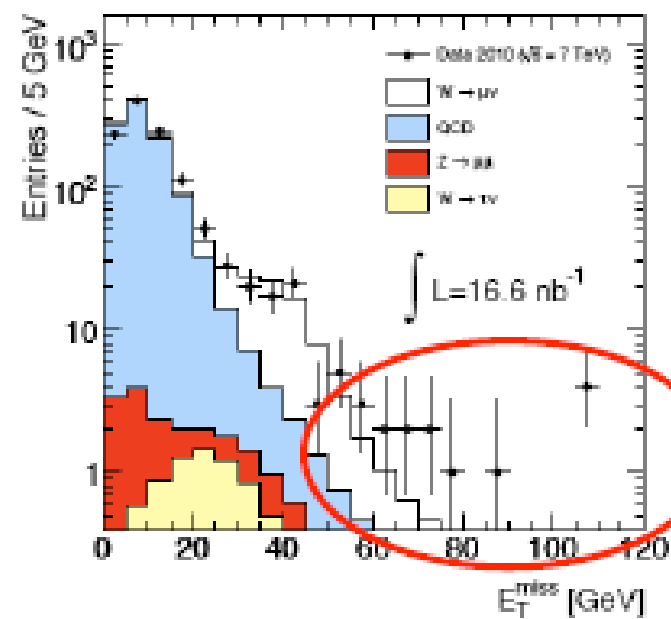
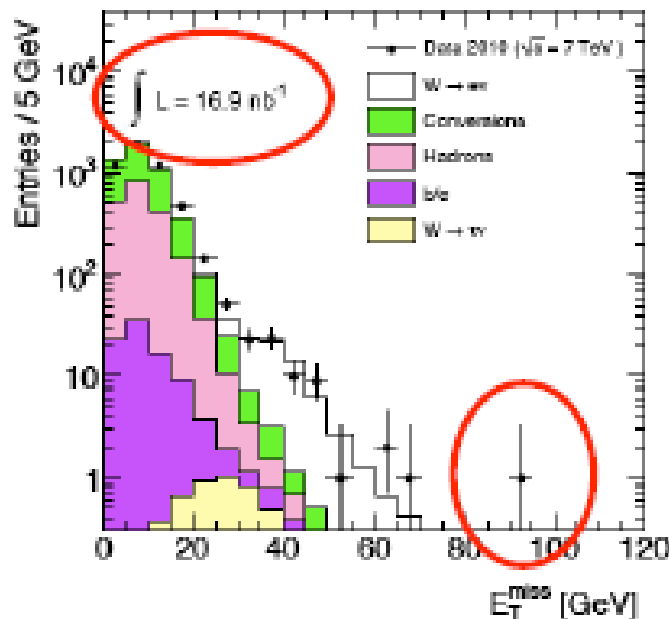
RPV

Missing ET

❖ “Evil” variable: - Σ (everything else)

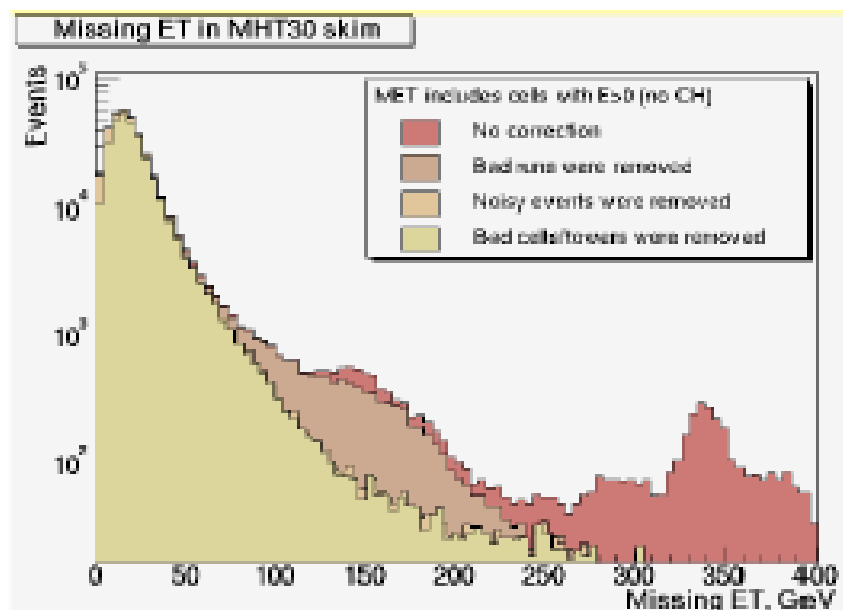
- Need to understand “everything else”
- Good benchmark: leptonic W boson decays

Early 2010

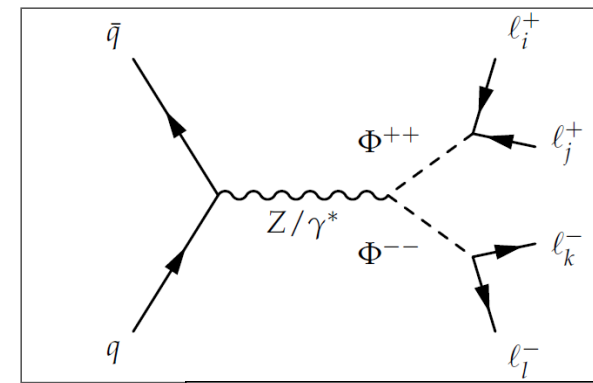


❖ Analyses using MET are particularly sensitive

- Requires the full calorimeter to behave, and calorimeter is generally the most sensitive subdetector (analog, ~ 16 bits)
- Easy: basic DQ (high voltage trip, etc.)
- Hard: low frequency
- Can't spot a 10^{-5} Hz (once a day) effect online or in first pass DQ
- But can be biggest part of dataset after cuts!

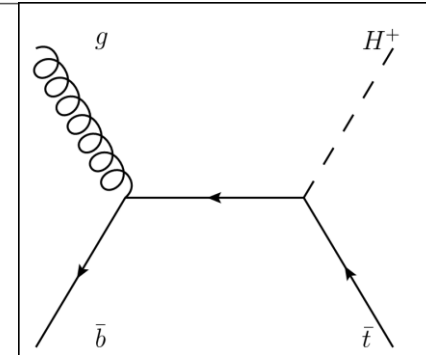


Extended Higgs sector



In the Standard Model single Higgs doublet, often

written as $\begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix}$ or $\begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix}$



Extended: Many choices but a few constraints,

for instance suppression of FCNC and $\frac{M_W}{M_Z} = \cos \theta_W$

- Most successful: 2 Higgs doublet models (2HDMs)
 - Supersymmetry uses this
- See-saw models predict Higgs triplet with φ^0 , $\varphi^{+/-}$, $\varphi^{++/--}$

General 2HDM Potential

$$\begin{aligned}
 V(\phi_1, \phi_2) = & \lambda_1 \left(|\phi_1|^2 - v_1^2 \right)^2 + \lambda_2 \left(|\phi_2|^2 - v_2^2 \right)^2 \\
 & + \lambda_3 \left[\left(|\phi_1|^2 - v_1^2 \right) + \left(|\phi_2|^2 - v_2^2 \right) \right]^2 \\
 & + \lambda_4 \left[|\phi_1|^2 |\phi_2|^2 - \left(\phi_1^{*T} \phi_2 \right) \left(\phi_2^{*T} \phi_1 \right) \right] \\
 & + \lambda_5 \left[\text{Re} \left(\phi_1^{*T} \phi_2 \right) - v_1 v_2 \cos \xi \right]^2 \\
 & + \lambda_6 \left[\text{Im} \left(\phi_1^{*T} \phi_2 \right) - v_1 v_2 \sin \xi \right]^2
 \end{aligned}$$

All λ are real.

From "Higgs Hunter's guide".

$$\begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}_1 = \begin{pmatrix} 0 \\ v_1 \end{pmatrix} \frac{1}{\sqrt{2}} \quad ; \quad \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}_2 = \begin{pmatrix} 0 \\ v_2 e^{i\xi} \end{pmatrix} \frac{1}{\sqrt{2}} \quad ; \quad \tan \beta = \frac{v_2}{v_1}$$

Higgs Boson Spectroscopy

- One Charged Higgs with mass:

$$m_{H^\pm} = \sqrt{\lambda_4 (v_1^2 + v_2^2)}$$

- One CP-odd neutral Higgs with mass:

$$m_{A^0} = \sqrt{\lambda_6 (v_1^2 + v_2^2)}$$

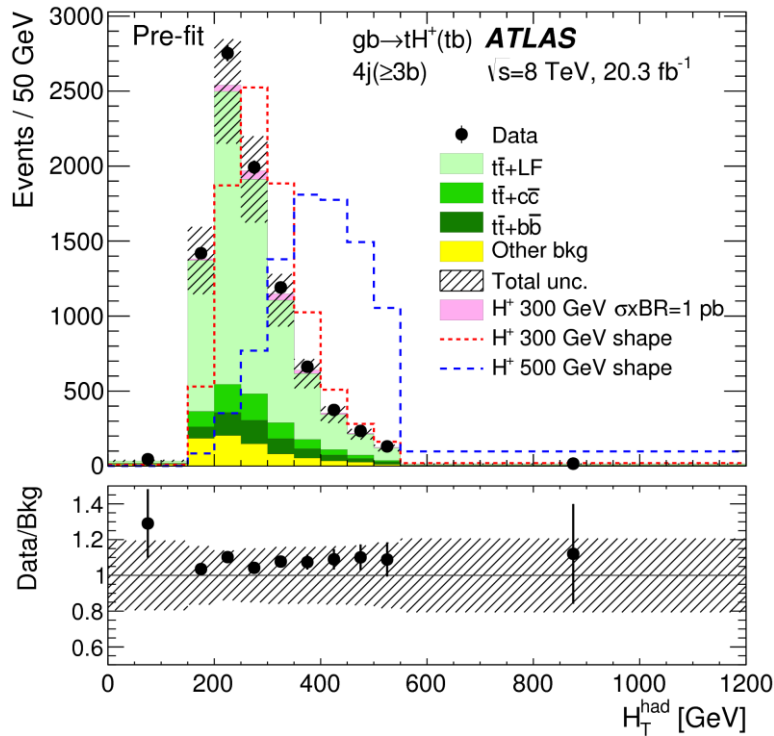
- And two CP-even higgs that mix.

$$M = \begin{pmatrix} 4v_1^2(\lambda_1 + \lambda_3) + v_2^2\lambda_5 & (4\lambda_3 + \lambda_5)v_1v_2 \\ (4\lambda_3 + \lambda_5)v_1v_2 & 4v_2^2(\lambda_2 + \lambda_3) + v_1^2\lambda_5 \end{pmatrix}$$

5 Higgs bosons! h, H, A, H^\pm

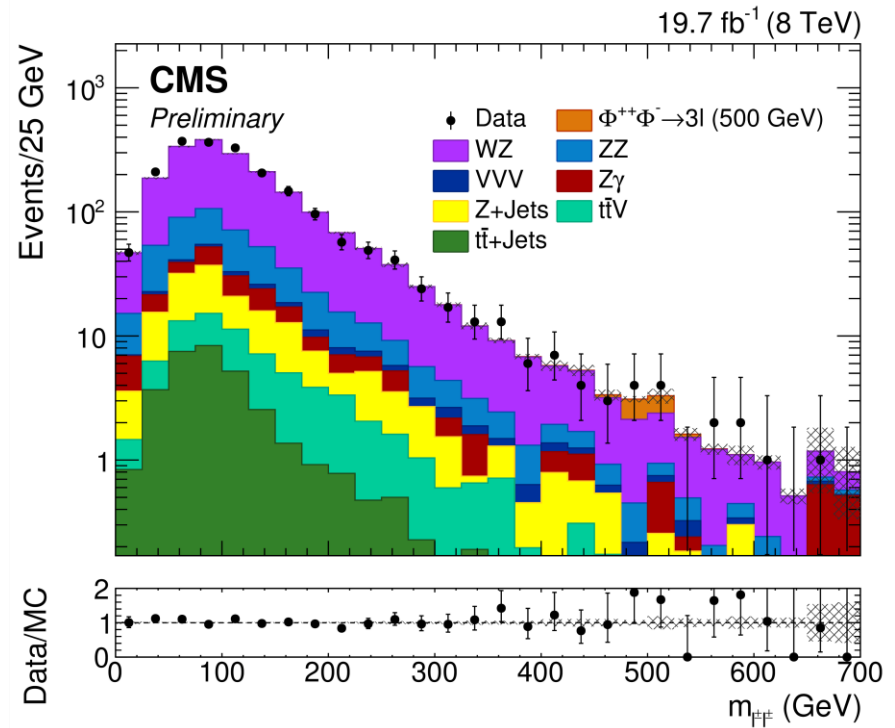
Examples of searches for extra Higgs bosons

Singly-charged



Limits around $\mathcal{O}(200 \text{ GeV})$

Doubly-charged



600 GeV

Parity Restoration: Signals

- ❖ Primary signals are (right-handed) W' (+ Z')
 - Dilepton resonances (Z') offer clean signals, well-understood backgrounds
 - At LHC, some concern about extrapolation of calibration from Z to very high energies
 - Electron/muon resolution improves/degrades with p_T
 - $t\bar{t}$ decays visible
 - ν_R is presumably heavy, W' may not decay to leptons
 - Only dijet or diboson
 - If ν_R lighter than W'/Z' , ν_R decays become important
- ❖ Note: many kinds of Z' - review by Langacker
 - W'/Z' would also require new fermions...

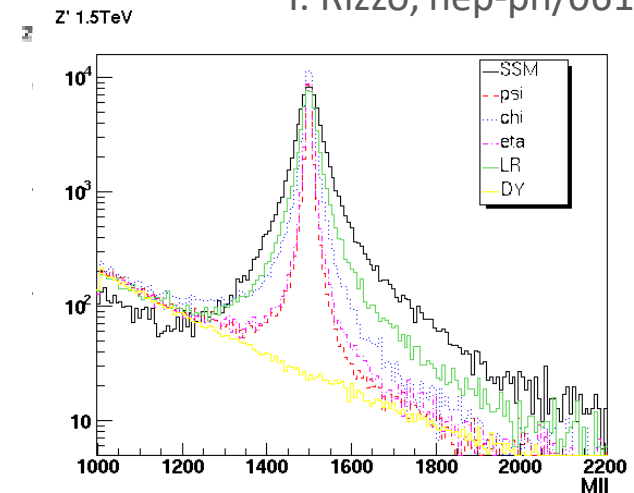
arXiv:0801.1345

Z' Production and Decay

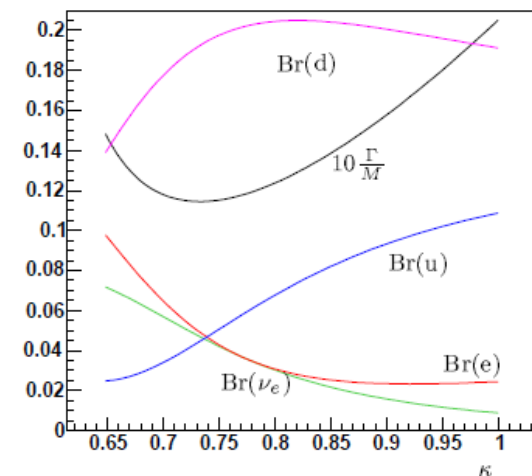
T. Rizzo, hep-ph/0610104

- ❖ Production from u, d quarks is dominant at LHC
 - Couplings vary by model
 - E.g. for LR symmetric models, $\kappa = g_R/g_L$ drives production cross-section (convolute with PDFs) and branching ratios

- ❖ Decays somewhat similar to Z (but almost no BR to light neutrinos, decays to top open up), plot assumes ν_R heavier



ATL-PHYS-PUB-2005-010

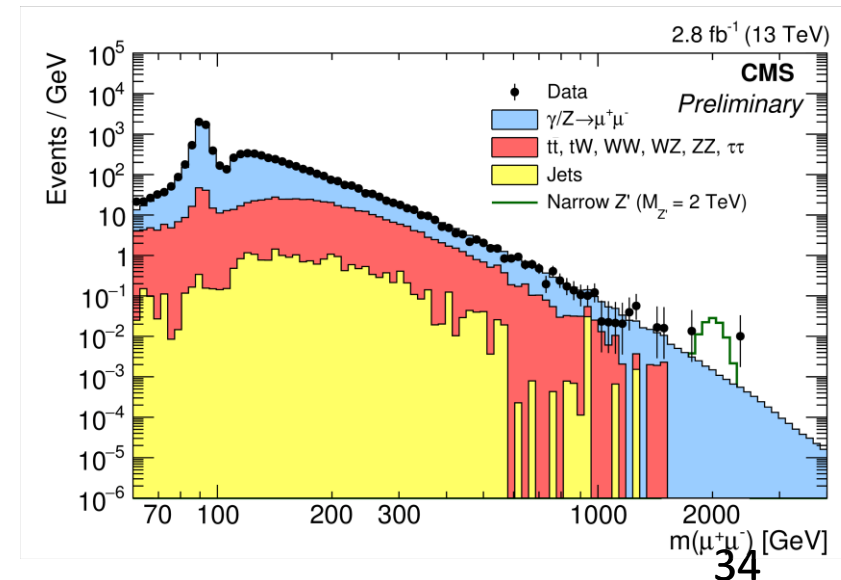
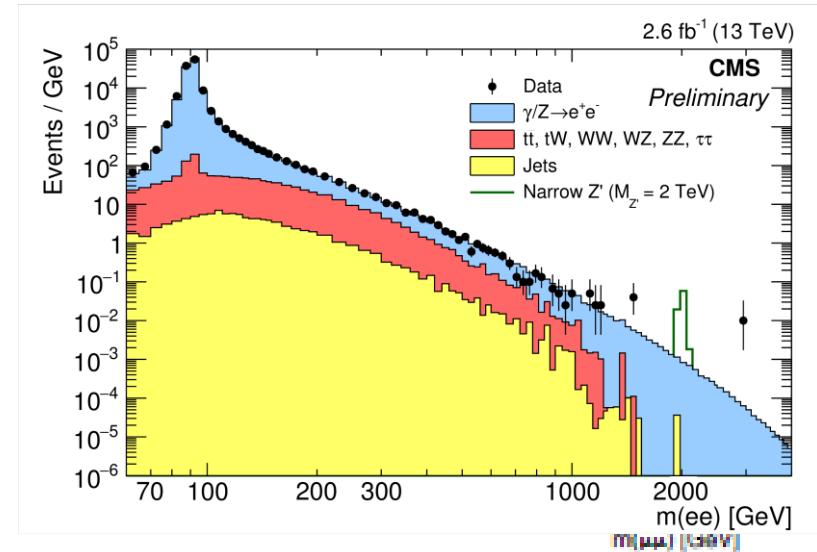


$Z' \rightarrow ee/\mu\mu$

❖ Most promising channels:

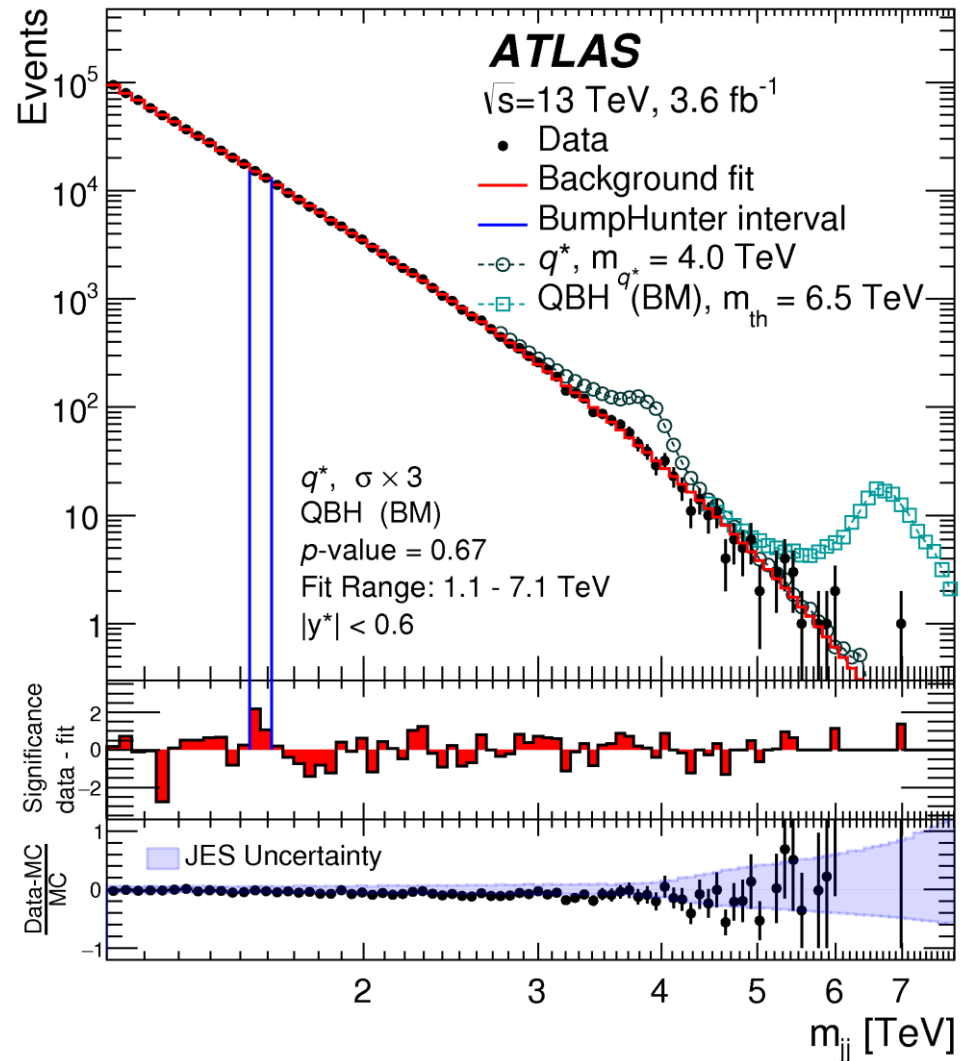
- Backgrounds very low!
- “Self-calibrating”
- In ee , at high masses, energy resolution dominated by constant term
- 10 GeV for 1.5 TeV electron
- Could measure width!

❖ LHC extended Tevatron reach immediately!



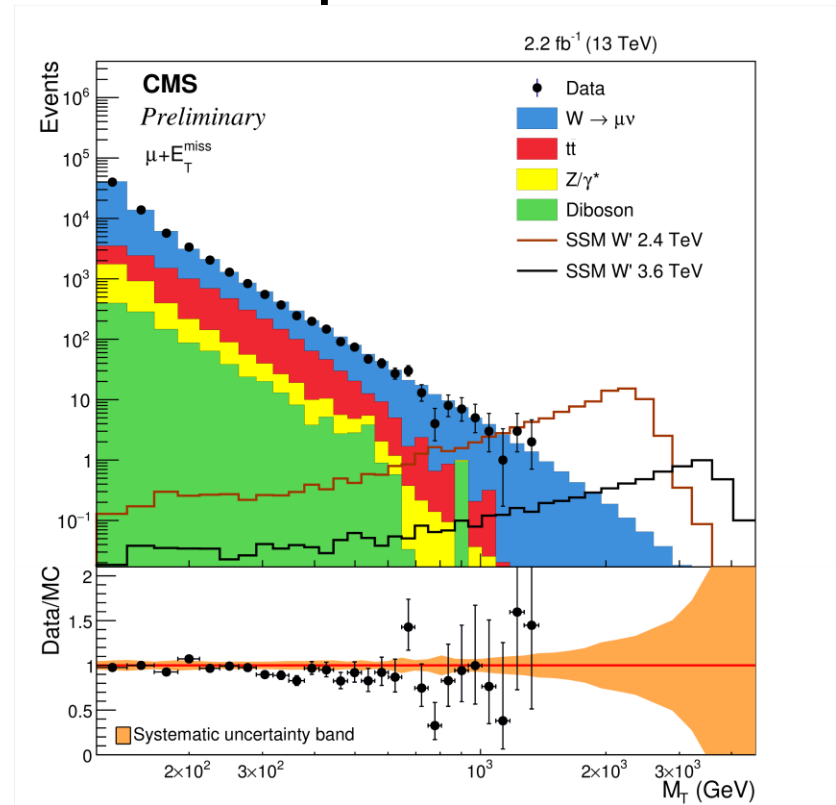
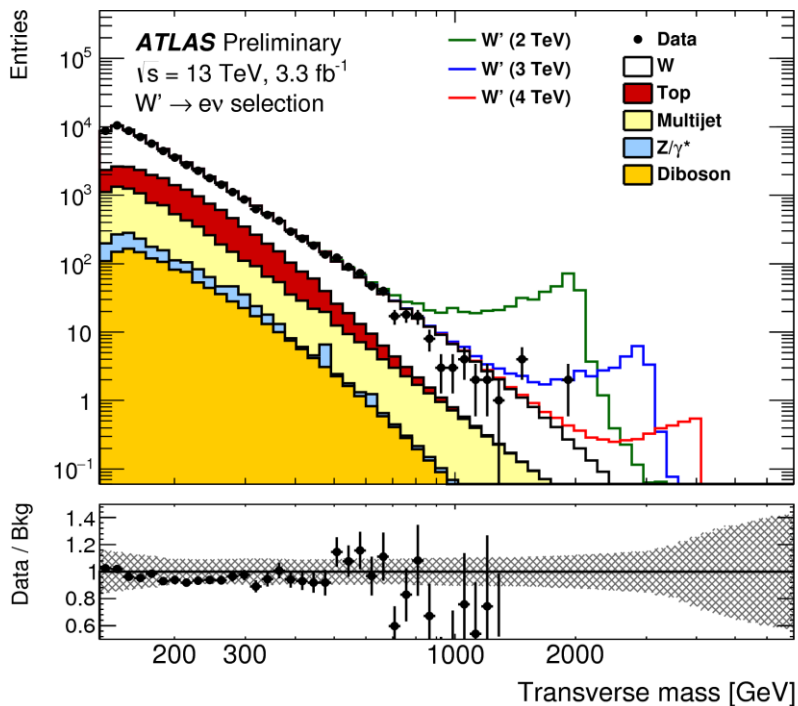
Dijets

- SM background obviously much larger
 - But single source
 - And opens the door to strongly interacting objects



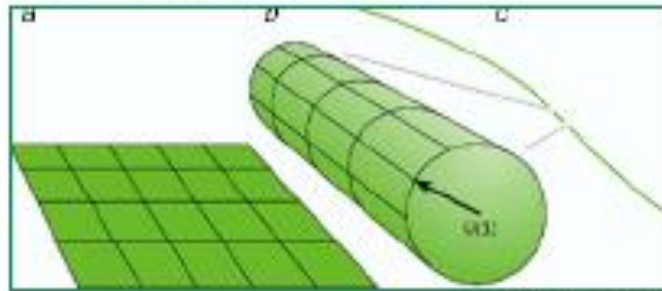
$$W' \rightarrow \mu\nu / e\nu$$

Another very simple selection: lepton + MET



Extra Dimensions

- ❖ A promising approach to quantum gravity consists in adding extra space dimensions: string theory
 - Additional space dimensions are hidden, presumably because they are compactified



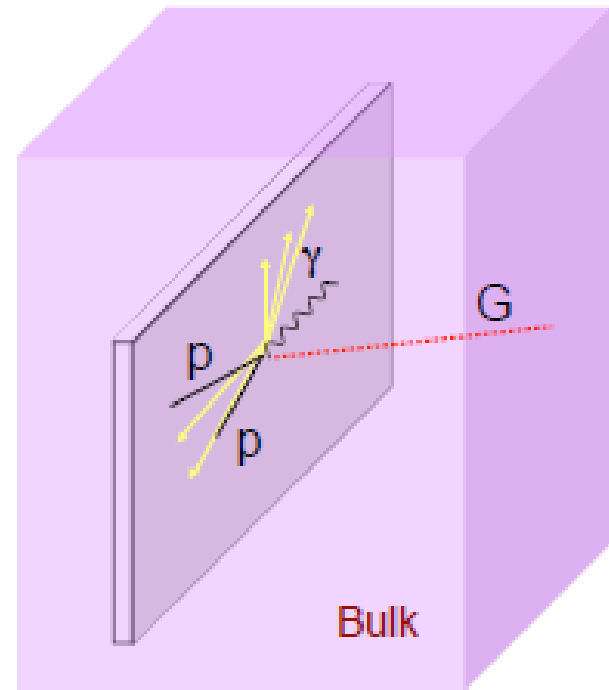
Source: PhysicsWorld

- ❖ Radius of compactification usually assumed to be at the scale of gravity, i.e. 10^{18} GeV
 - In '90 Antoniadis realized they may be much larger...

Phys Lett B246:387-384 1990

ADD extra dimensions

- ❖ “Large extra dimension” scenario (developed by Arkani-Hamed, Dimopoulos and Dvali): Phys.Lett. B429 (1998) 263-272
 - Standard model fields are confined to a 3+1 dimensional subspace (“brane”)
 - Gravity propagates in all dimensions
 - Gravity appears weak on the brane because only felt when graviton “goes through”

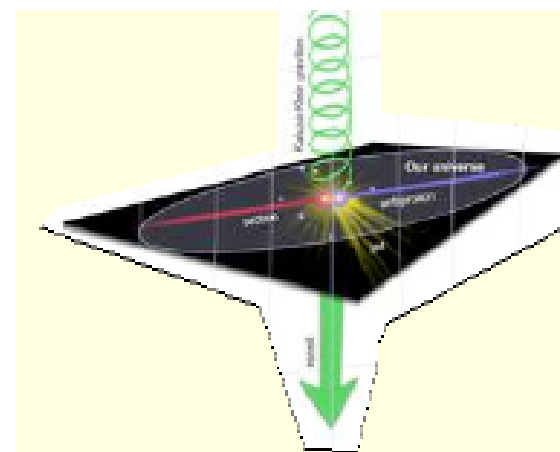


Drawing by K. Loureiro

ADD signatures

- ❖ Edges of extra dimensions identified
 - ➔ Boundary conditions
 - ➔ Momentum along extra dimension is quantized
 - Looks like mass to us
 - Very small separations → looks like continuum
 - Called Kaluza-Klein tower
- ❖ Coupling to single graviton very weak, but there are *lots* of them!
 - Large phase space → observable cross-section
 - Impacts all processes (graviton couples to energy-momentum)

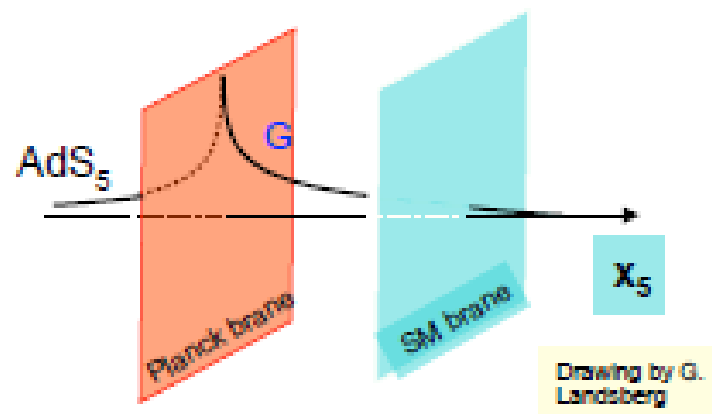
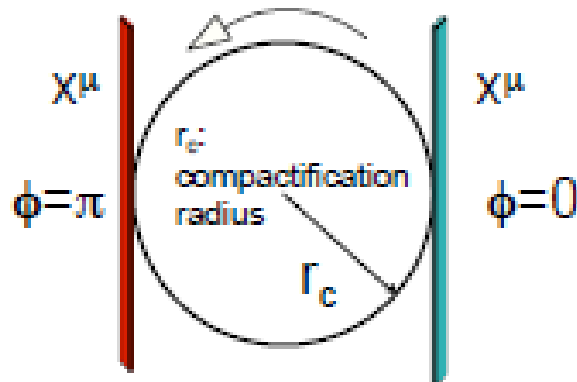
- ❖ Consider processes that involve the bulk (i.e. gravitons)
 - Translational invariance is broken
 - ➔ Momentum is not conserved ...
 - ... because graviton disappears in bulk right away
- ❖ Look for $p p \rightarrow \text{jet/photon} + \text{nothing}$ (i.e. \cancel{E}_T), or deviations in high mass/angular behavior in standard model processes
 - Graviton has spin 2, couples to energy-momentum!



Warped extra dimensions

❖ “Simple” Randall-Sundrum model:

- SM confined to a brane, and gravity propagating in an extra dimension
- As opposed to the original ADD scenario, the metric in the extra dimension is “warped” by a factor $\exp(-2kr_c\phi)$
- (Requires 2 branes)



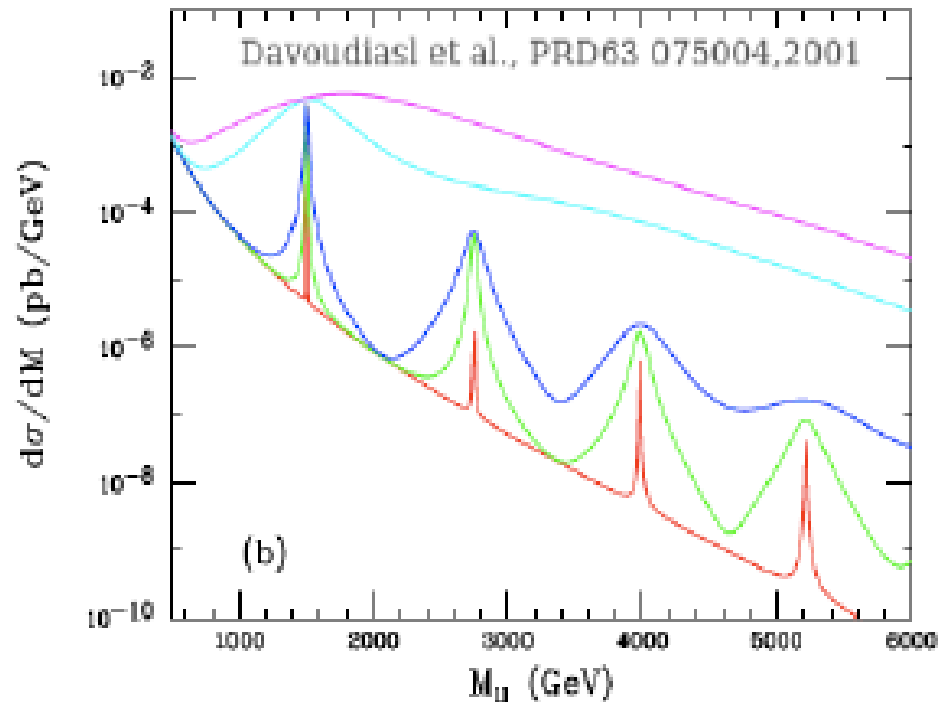
Graviton excitations

❖ In RS, get a few massive graviton excitations

- Widths depend on warp factor k
- Mass separation = zeros of Bessel function

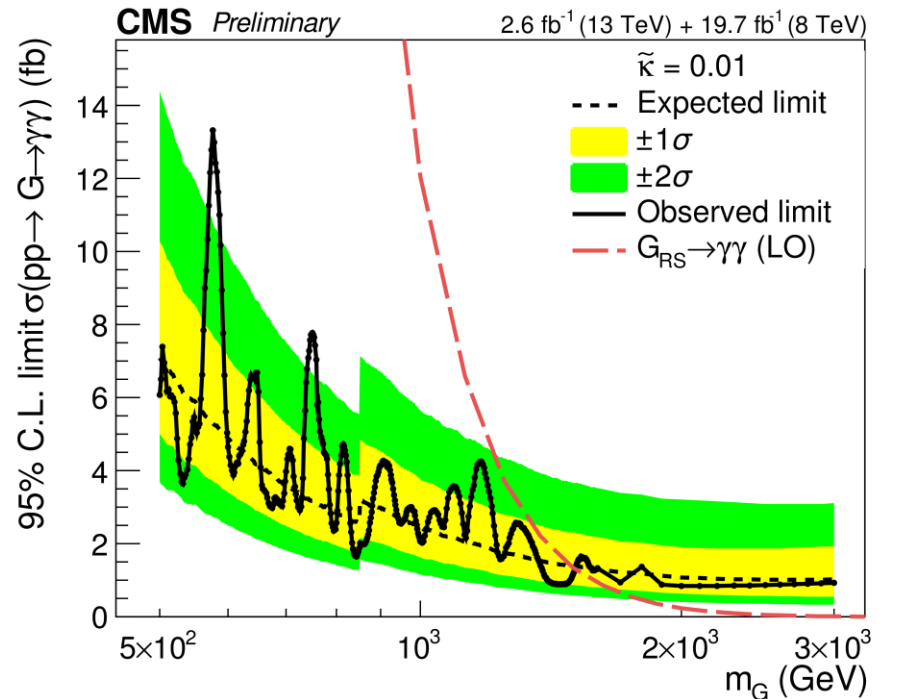
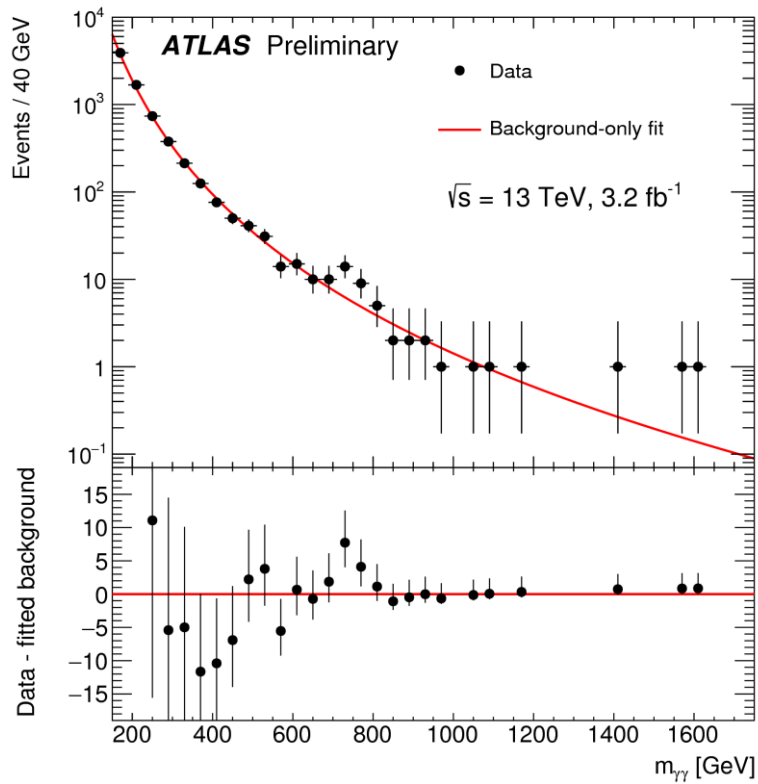
➔ Smoking gun!

(BRs also different than Z' :
e.g. $\gamma\gamma$ allowed)



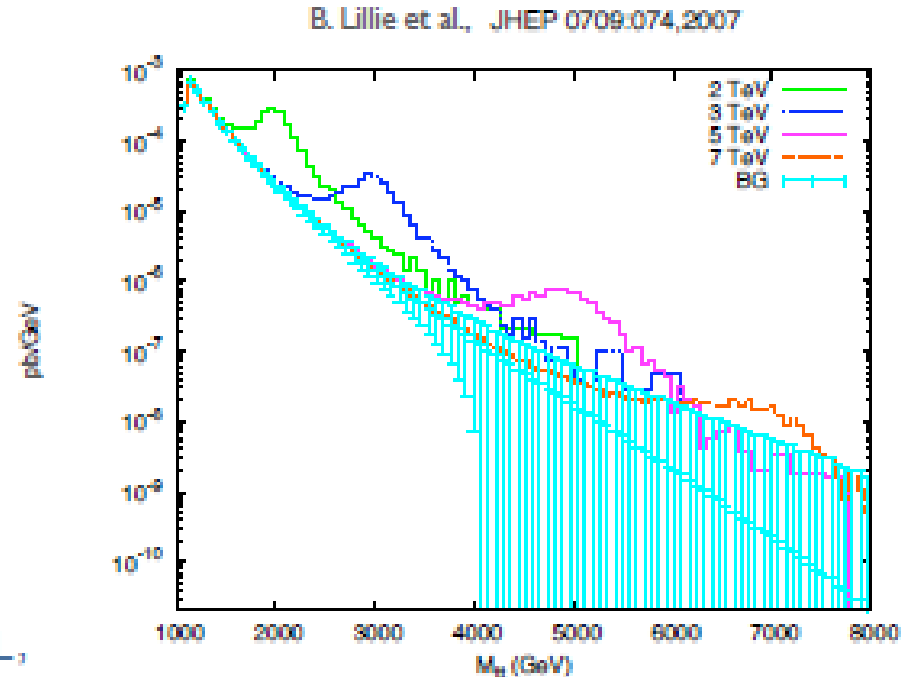
Example

The infamous $\gamma\gamma$ bump is an example of a search for RS gravitons:



Gauge boson excitations

- ❖ Excitations of the gauge bosons are very promising channels for discovery
 - Couplings to light fermions are small
 - Small production cross-sections
 - Large coupling to top, W_L , Z_L
 - Look for $t\bar{t}$, WW , ZZ resonances (that can be wide)

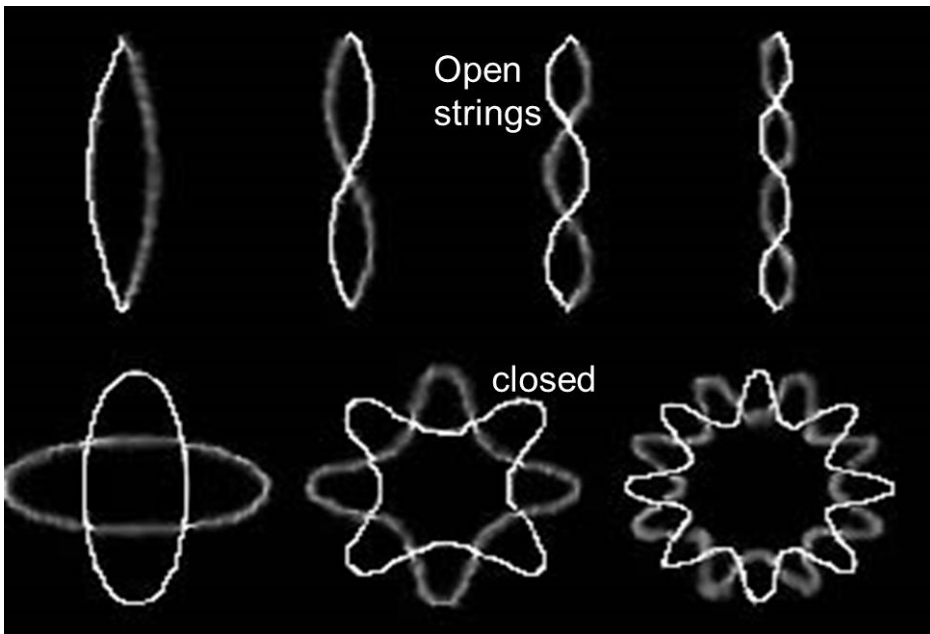
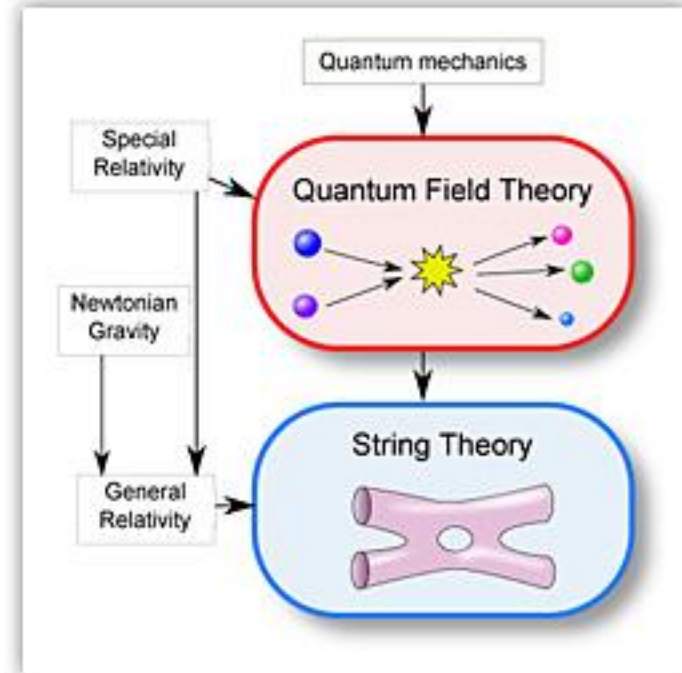


(super)Strings

Avoid infinities from point-like particles

Different vibration modes = different particles

One fundamental parameter: *string size*



Great idea but we have not yet understood how to test it at current "low" energies

Extra dimensions a must
Supersymmetry a plus