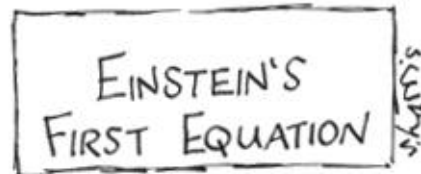


The representation of observables

...some simple formulae for operators



A hand-drawn scribble representing a piece of paper with two equations written on it: $3x = 12$ and $x = 6$.



A hand-drawn box containing the text "EINSTEIN'S FIRST EQUATION".

The recipe for cooking up operators



1. take classical variable
2. replace as follows:

$$x \rightarrow \mathbf{x}\psi = x\psi$$

$$p_x \rightarrow \mathbf{p}_x\psi = -i\frac{h}{2\pi}\frac{\partial\psi}{\partial x}$$

3. do the same for y and z
4. same applies to relations between variables

The position operator(s)

A vector* operator

$$\mathbf{R} = x\vec{u}_x + y\vec{u}_y + z\vec{u}_z$$

← operators act only on wavefunctions

where

$$x\psi = x\psi, \quad y\psi = y\psi \quad \text{and} \quad z\psi = z\psi$$

are the component operators

* Operators of this form are *not* ordinary 3d vectors!

The linear momentum operator(s)

A vector operator

$$\mathbf{p} = p_x \vec{u}_x + p_y \vec{u}_y + p_z \vec{u}_z$$

where

$$p_x \psi = -i \frac{h}{2\pi} \frac{\partial \psi}{\partial x}, \quad p_y \psi = -i \frac{h}{2\pi} \frac{\partial \psi}{\partial y} \quad \text{and} \quad p_z \psi = -i \frac{h}{2\pi} \frac{\partial \psi}{\partial z}$$

are the component operators

► *What is the operator for p^2 ?*

Multiplication of operators



What does it mean to ‘multiply’ operators?



It means that we apply the operator to a wavefunction more than once:

e.g. $p_x^2 \psi = p_x(p_x \psi)$ where $p_x \psi = -i \frac{h}{2\pi} \frac{\partial \psi}{\partial x}$

gives us

$$p_x^2 \psi = \left(-i \frac{h}{2\pi}\right) p_x \left(\frac{\partial \psi}{\partial x}\right) = \left(-i \frac{h}{2\pi}\right)^2 \frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial x}\right) = -\left(\frac{h}{2\pi}\right)^2 \frac{\partial^2 \psi}{\partial x^2}$$

The operator for orbital angular momentum

The classical representation (variable) is $\vec{L} = \vec{r} \times \vec{p}$, so the quantum mechanical representation (operator) will be $\mathbf{L} = \mathbf{r} \times \mathbf{p}$, which gives

$$L_x = (yp_z - zp_y), L_y = (zp_x - xp_z), L_z = (xp_y - yp_x)$$

For example,

$$L_x \psi = (yp_z - zp_y) \psi = yp_z \psi - zp_y \psi = y(p_z \psi) - z(p_y \psi)$$

$$L_x \psi = -i \frac{h}{2\pi} \left(y \frac{\partial \psi}{\partial z} - z \frac{\partial \psi}{\partial y} \right)$$

The kinetic energy operator

We use the formula of classical mechanics:

$$\mathbf{K} = \frac{1}{2m}\mathbf{p}^2 = \frac{1}{2m}(\mathbf{p}_x^2 + \mathbf{p}_y^2 + \mathbf{p}_z^2) = -\frac{1}{2m}\left(\frac{h}{2\pi}\right)^2 \tilde{\mathbf{N}}^2$$

where

$$\tilde{\mathbf{N}}^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

The total energy operator (when $E=K+V$)

We assume an object in motion (K) in a region of a potential V :

$$\mathbf{H} = \mathbf{K} + \mathbf{V}(x, y, z) = -\frac{1}{2m} \left(\frac{h}{2\pi} \right)^2 \tilde{\mathbf{N}}^2 + V(x, y, z)$$

Observe that the potential operator is given by the classical variable because it depends on x, y, z

The operator H is called the Hamiltonian operator ('the Hamiltonian of the system')

The operator E

$$E = i \frac{h}{2\pi} \frac{\partial}{\partial t}$$

It gives the time-dependence of energy eigenstates

By combining with the expression $E=K+V$, we obtain:

$$i \frac{h}{2\pi} \frac{\partial \Psi}{\partial t} = -\frac{1}{2m} \left(\frac{h}{2\pi} \right)^2 \tilde{N}^2 \Psi + V\Psi$$

the time-dependent Schrödinger equation

The commutator operator

Can we ‘multiply’ two operators A and B in any order?

We can, if their commutator $[A, B]$ is zero.

Definition of the commutator operator:

$$[A, B] = AB - BA$$

Physical meaning of $[A, B] = 0$:
the operators A and B have
common (simultaneous) eigenstates

Examples of commutators

1. position and momentum in the same direction:

$$[x, p_x] = i \frac{h}{2\pi} \text{ (same relation for y and z)}$$

2. angular momentum in different directions:

$$[L_x, L_y] = i \frac{h}{2\pi} L_z \text{ (and cyclic permutations) but}$$

$$[L^2, L_x] = [L^2, L_y] = [L^2, L_z] = 0$$

Common eigenstate problem in atoms

Assume that we have an object (atom) in a potential $V(r)$ and we are interested in states of constant total energy, E .

The following relations hold:

$$[x, H] \neq 0, [y, H] \neq 0, [z, H] \neq 0$$

$$[H, p_x] \neq 0, [H, p_y] \neq 0, [H, p_z] \neq 0$$

$$[L^2, H] = [H, L_z] = [L^2, L_z] = 0$$

So we can measure simultaneously: E, L^2, L_z .

Mean value and Expectation value

We measure the observable A n times: a_1, a_2, \dots, a_n
and we calculate the mean value as

$$\bar{a} = \frac{1}{n} \sum_{i=1}^n a_i$$

Postulate of Quantum Mechanics:

$$\bar{a} = \langle \mathbf{O} \rangle = \langle \mathbf{O} \rangle = \int \psi^* \mathbf{O} \psi dV$$

where

$\langle \mathbf{O} \rangle$ is the expectation value of the operator \mathbf{O} in the state ψ .

The standard deviation

Gives the spread of the measured values
around the mean value

$$\Delta O = \sqrt{\sum_k w_k (a_k - \bar{a})^2}$$

where w_k = relative frequency of value a_k
Quantum mechanically, we calculate

$$\Delta O = \sqrt{\int \psi^* (O - \langle O \rangle)^2 \psi dV}$$

The Heisenberg principle

Assume that we measure P and Q in the state ψ (which is not an eigenstate of their operators). The following holds for their standard deviations:

$$(\Delta P)^2 \cdot (\Delta Q)^2 \geq \left(-\frac{1}{4} \left[\int \psi^* (PQ - QP) \psi dV \right]^2 \right)$$

If the operators do *not* commute, the right-hand side is not zero. Examples:

$$\Delta p_x \cdot \Delta x = h/(4\pi), \quad \Delta E \cdot \Delta t = h/(4\pi)$$

Conserved Observables

An observable is conserved (i.e. its expectation value does not change with time) when its operator commutes with the Hamiltonian (of the system we study)

$$\frac{d}{dt} \langle \mathbf{O} \rangle = \frac{2\pi i}{h} \langle [\mathbf{H}, \mathbf{O}] \rangle$$

Application: the angular orbital momentum of atoms is a constant of the motion