

# Modern Experimental Particle Physics

FYST17

*February 22 and 23, 2017*

An introduction to the  
Quark Gluon Plasma  
and heavy-ion collisions

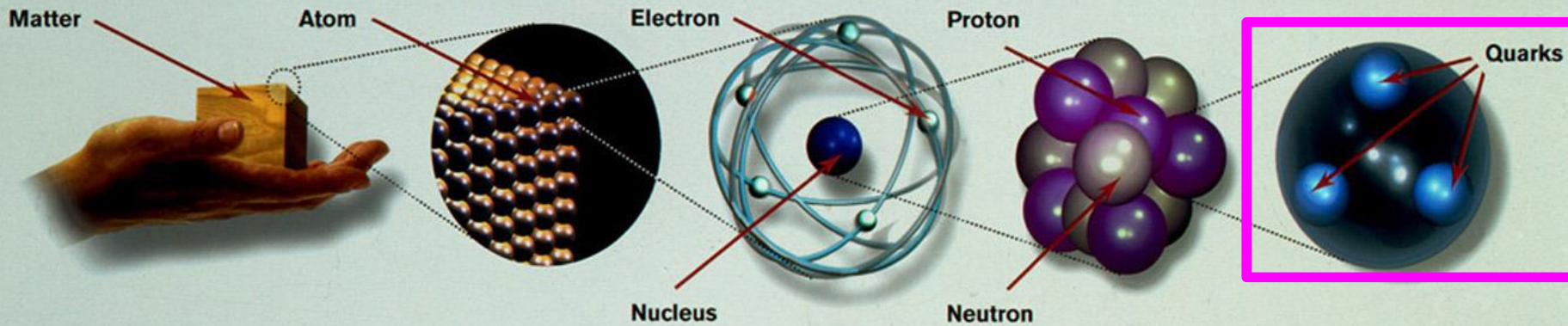
P. Christiansen (Lund University)





# Outline







- The medium temperature
- Hard probes
  - The standard candles
  - Quarkonium
  - Jets and high  $p_T$  particles
- Soft medium properties
  - Collective flow
- Quark Gluon Plasma (QGP) in small systems?



**Matter particles**







All ordinary particles belong to this group

**LEPTONS**

<b>FIRST FAMILY</b>	<p><b>Electron</b> Responsible for electricity and chemical reactions; it has a charge of -1</p> 	<p><b>Electron neutrino</b> Particle with no electric charge, and possibly no mass; billions fly through your body every second</p> 
<b>SECOND FAMILY</b>	<p><b>Muon</b> A heavier relative of the electron; it lives for two-millionths of a second</p> 	<p><b>Muon neutrino</b> Created along with muons when some particles decay</p> 
<b>THIRD FAMILY</b>	<p><b>Tau</b> Heavier still; it is extremely unstable. It was discovered in 1975</p> 	<p><b>Tau neutrino</b> not yet discovered but believed to exist</p> 

These particles existed just after the Big Bang. Now they are found only in cosmic rays and accelerators

**QUARKS**

<p><b>Up</b> Has an electric charge of plus two-thirds; protons contain two, neutrons contain one</p> 	<p><b>Down</b> Has an electric charge of minus one-third; protons contain one, neutrons contain two</p> 
<p><b>Charm</b> A heavier relative of the up; found in 1974</p> 	<p><b>Strange</b> A heavier relative of the down; found in 1964</p> 
<p><b>Top</b> Heavier still</p> 	<p><b>Bottom</b> Heavier still; measuring bottom quarks is an important test of electroweak theory</p> 

**Force particles**

These particles transmit the four fundamental forces of nature although gravitons have so far not been discovered

**Gluons**

Carriers of the strong force between quarks

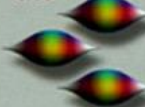


Felt by: quarks

The explosive release of nuclear energy is the result of the **strong force**

**Photons**

Particles that make up light; they carry the electromagnetic force



Felt by: quarks and charged leptons

Electricity, magnetism and chemistry are all the results of **electro-magnetic force**

**Intermediate vector bosons**  
Carriers of the weak force



Felt by: quarks and leptons

Some forms of radio-activity are the result of the **weak force**

**Gravitons**  
Carriers of gravity



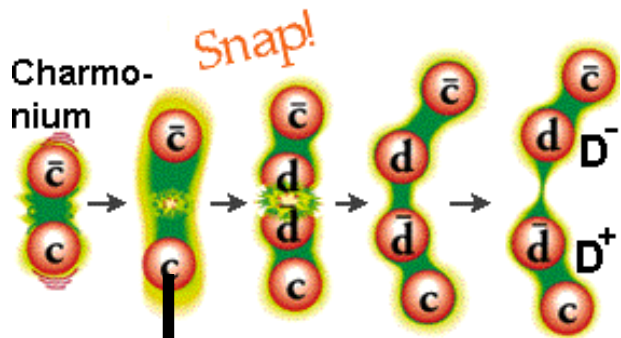
Felt by: all particles with mass

All the weight we experience is the result of the **gravitational force**

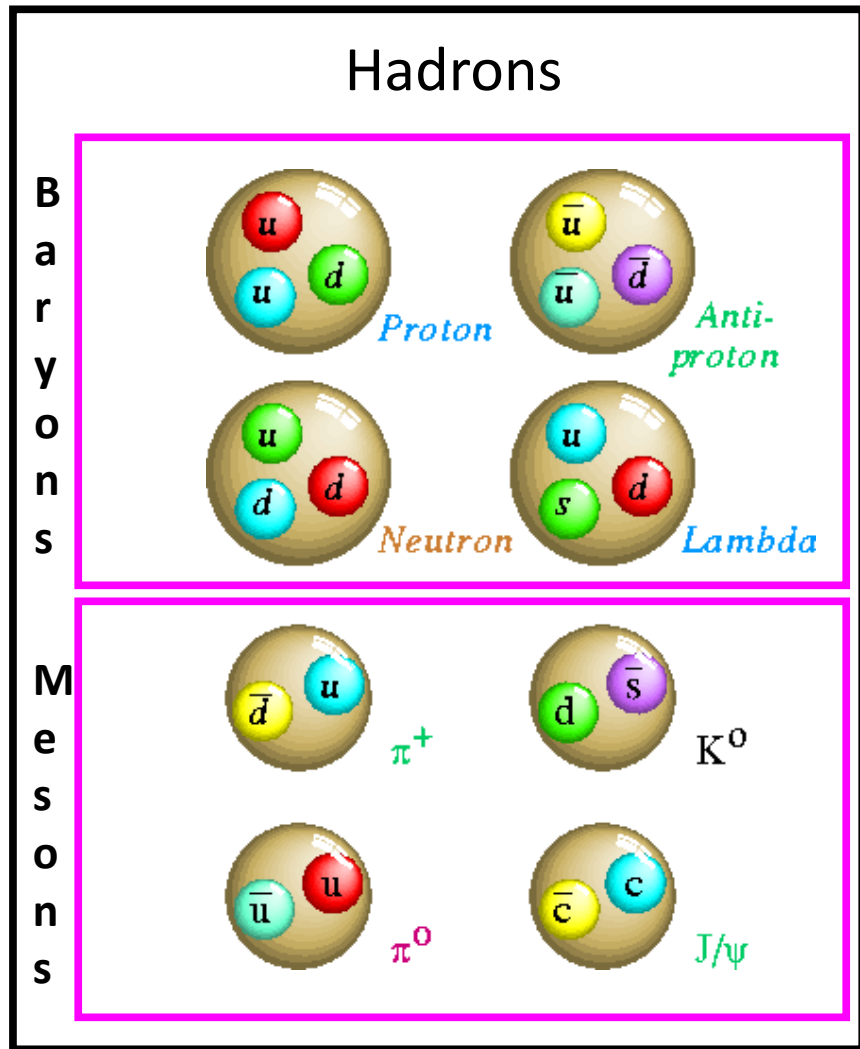


# The strong interaction: Quantum Chromo Dynamics (QCD)

3 strong charges (red, green, blue)  
 Particles in nature are color neutral  
 Quarks are “confined”



(Force  $\sim 1 \text{ GeV/fm}$ )



# Let us recall why it is so





# What is the Quark Gluon Plasma?

“QGP' is not a new particle but a paradigm-shift of how we understand matter in extreme conditions. Thus there is no clear criterion by which one can claim an experimental discovery. For this reason QGP was rediscovered again with new experimental results obtained at the order of magnitude higher RHIC collision energies. In addition to confirming the CERN results, RHIC produced new puzzling phenomena; some will be discussed below. The circumstance repeats for the third time today: LHC data confirms SPS and RHIC results, and is offering another rich field of new experimental results.

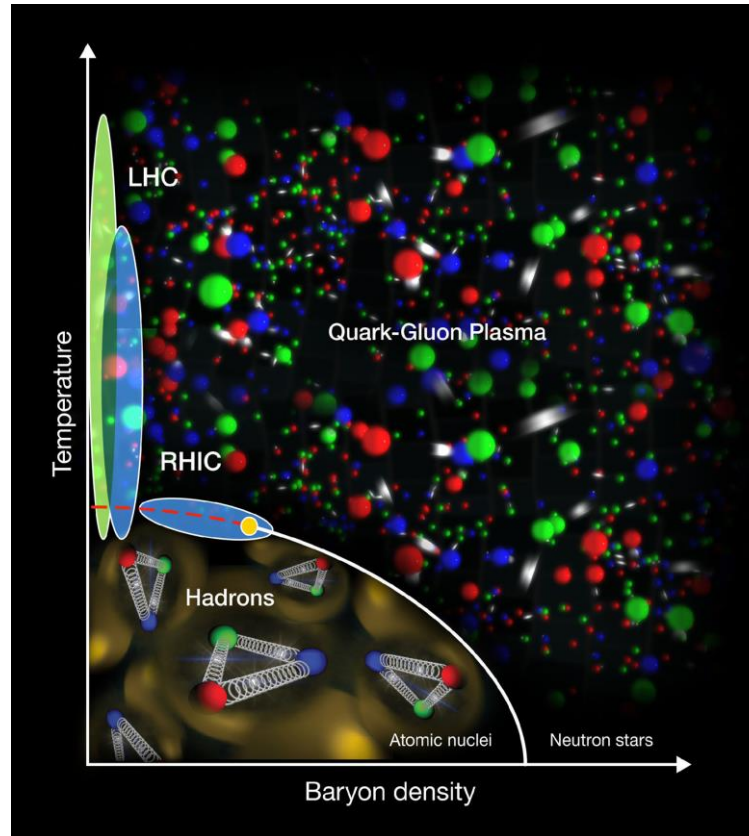
Since no one plans to announce the QGP discovery at the LHC, we conclude that QGP has gained considerable acceptance as a new form of matter.”

Johann Rafelski and Jeremey Birrell, arXiv:1311.0075.

They point out a big problem: our main understanding is experimental and very little of the theory is on a firm basis.

When we will have a better theoretical picture we can then really point out what the QGP properties are and when it was discovered.

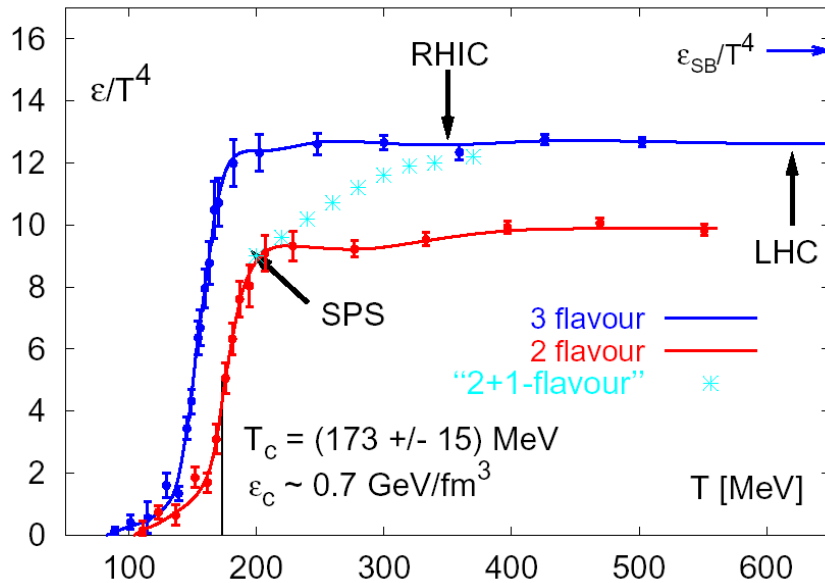
# The quark-gluon phase diagram



We will only talk about the high temperature transition which is what we probe in heavy-ion collisions



# Lattice QCD calculation of the energy density



$$\epsilon_{\text{Quark-Gluon gas}} = \frac{\pi^2}{30} (\text{bosonic dof} + \frac{7}{8} \text{ fermionic dof}) T^4$$

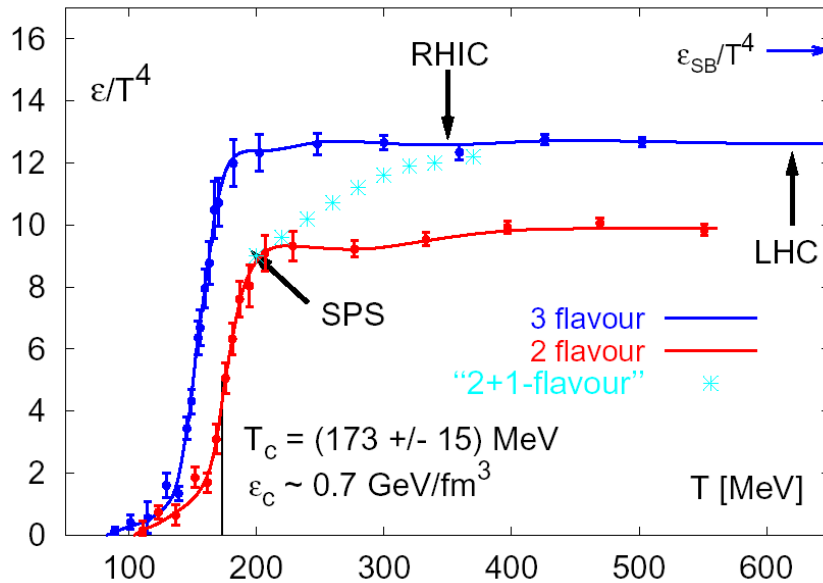
?

With lattice QCD one can study some aspects of QCD numerically. In this case the energy density shows that at a temperature of  $\sim 170$  MeV there is a phase transition. The phase transition is believed to be a Xover, meaning that for temperatures around  $T_c$  the hadronic and QGP phases coexist and that no entropy is produced in the phase transition.





# Lattice QCD calculation of the energy density



$$\epsilon_{\text{Quark-Gluon gas}} = \frac{\pi^2}{30} \left( 2 \times 8 + \frac{7}{8} 2 \times 2 \times 3 \times 3 \right) T^4$$

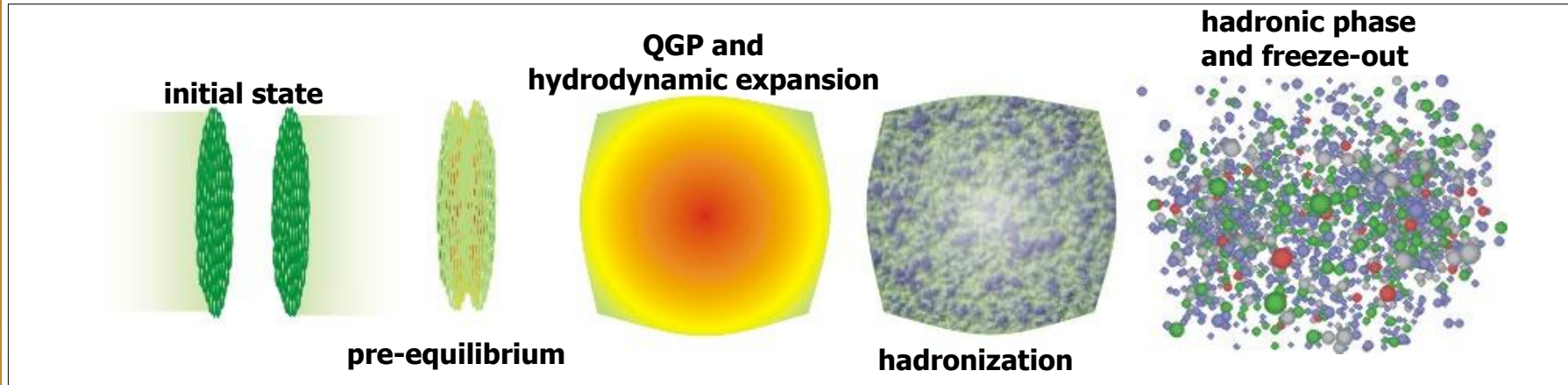
↑ Gluon spin and color  
 ↑ (Anti+)quark spin, color and flavor

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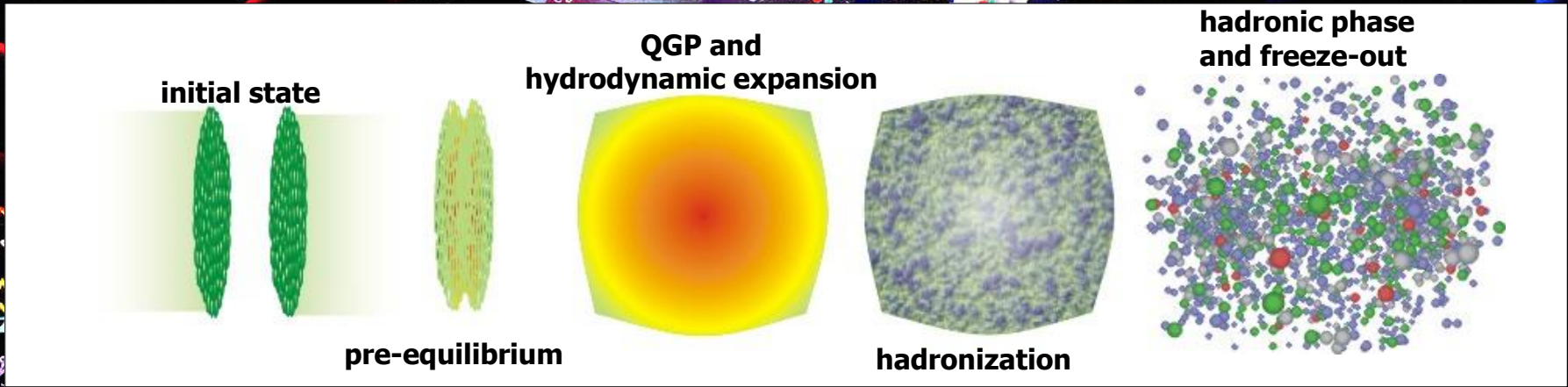
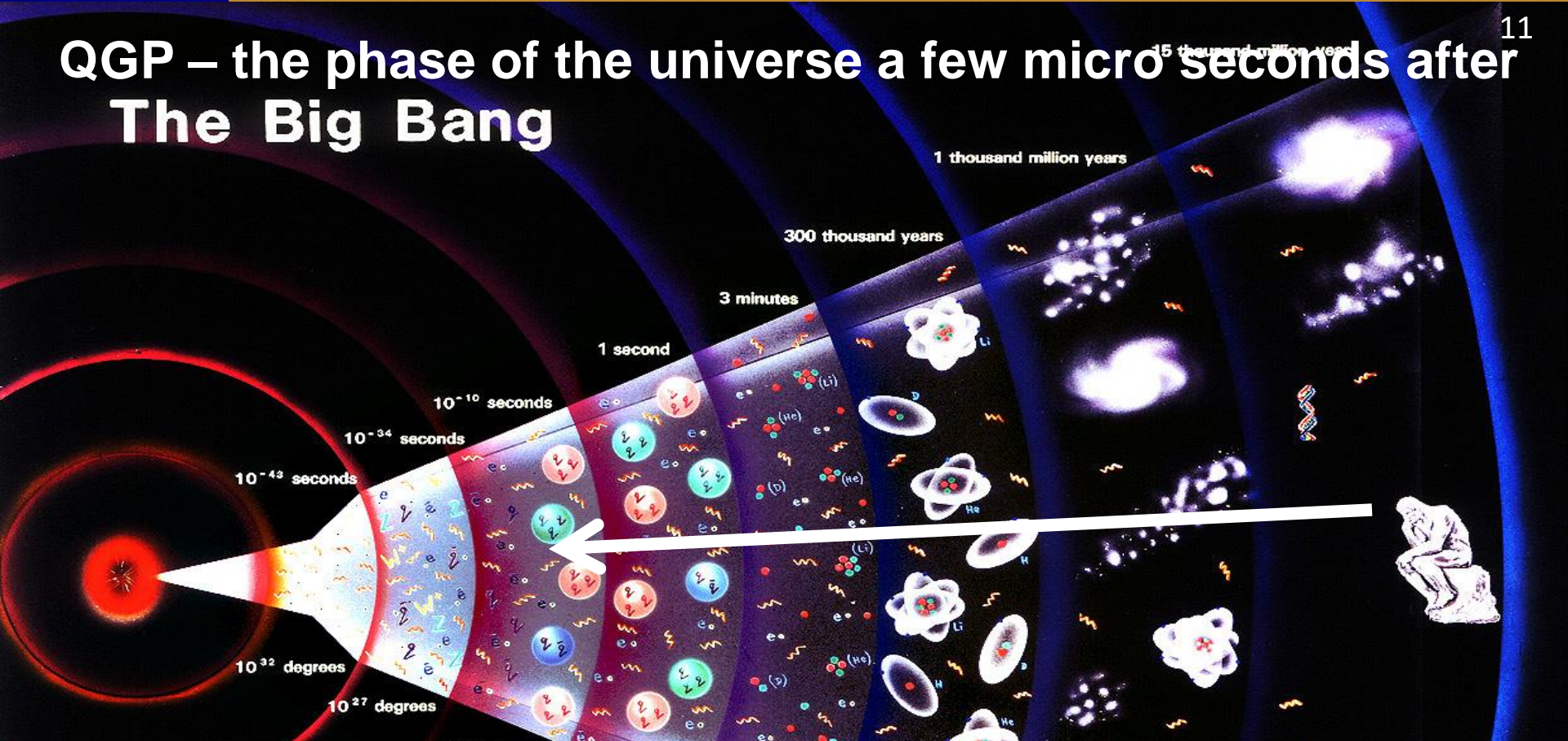


# Heavy ion collisions



- The only way we can create the QGP in the laboratory!
- By colliding heavy ions it is possible to create a large ( $\gg 1\text{fm}^3$ ) zone of hot and dense QCD matter
- Goal is to create and study the properties of the Quark Gluon Plasma
- Experimentally mainly the final state particles are observed, so the conclusions have to be inferred via models

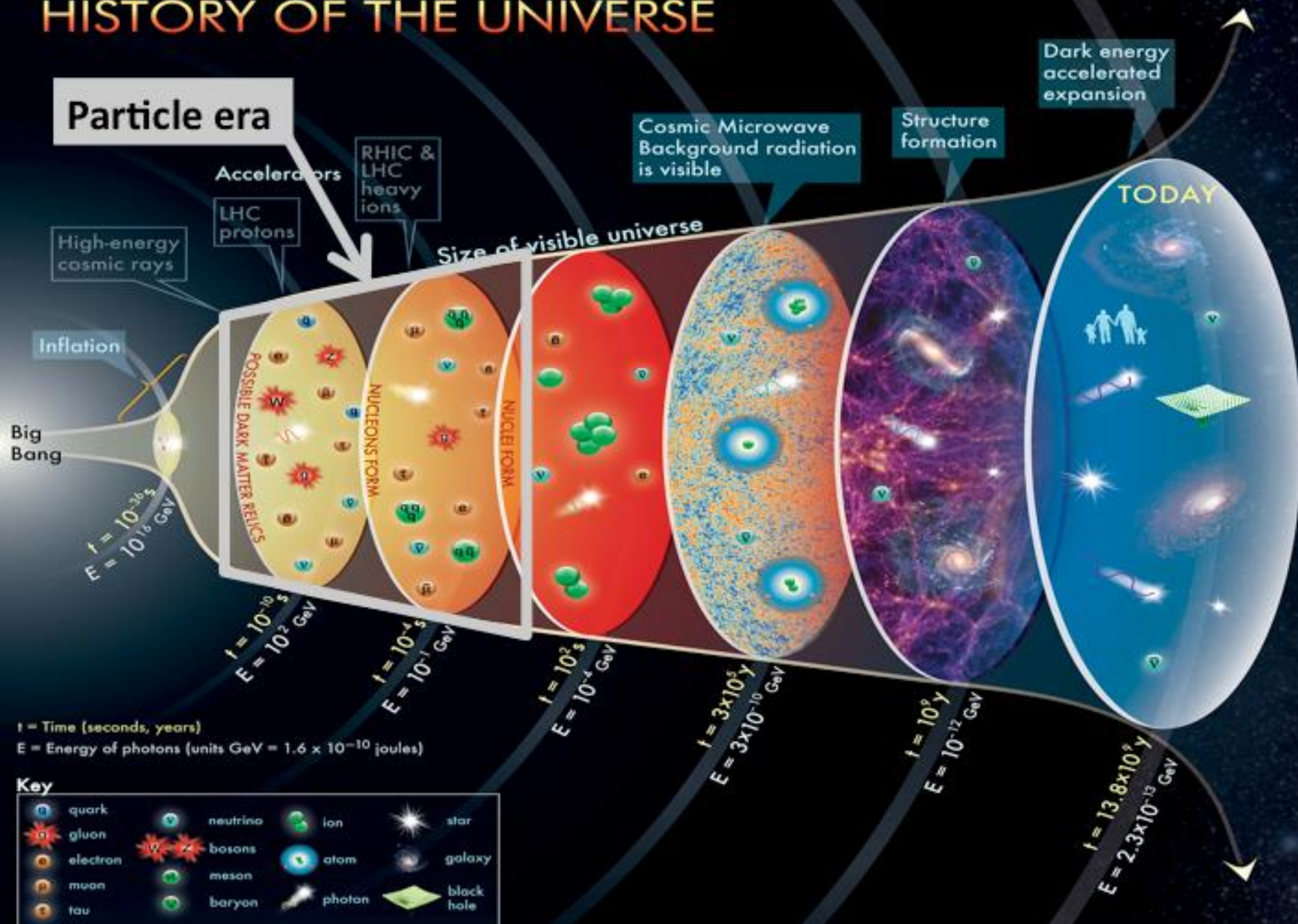
# QGP – the phase of the universe a few micro seconds after The Big Bang



anti-quark      He helium  
 e- electron      Li lithium

Ditzingen

# HISTORY OF THE UNIVERSE



t = Time (seconds, years)  
 E = Energy of photons (units GeV =  $1.6 \times 10^{-10}$  joules)

**Key**

	quark		neutrino		ion		star
	gluon		bosons		atom		galaxy
	electron		meson		photon		black hole
	muon		baryon				
	tau						

The concept for the above figure originated in a 1986 paper by Michael Turner.

# History of ultra-relativistic heavy-ion physics

- 1<sup>st</sup> generation  $\sqrt{s_{NN}} < 20 \text{ GeV}$  (fixed target)  
AGS (BNL, US), SPS (CERN) late 80s and 90s
  - Experiments: NA61, NA49, NA60, NA50, **NA44**, .....
- 2<sup>nd</sup> generation  $\sqrt{s_{NN}} < 200 \text{ GeV}$  (collider)  
RHIC (BNL, US) 2000-Now
  - Experiments: **PHENIX**, STAR, PHOBOS, **BRAHMS**
- 3<sup>rd</sup> generation  $\sqrt{s_{NN}} = 2760 \text{ GeV}$  in run 1 (collider),  
 $\sqrt{s_{NN}} = 5020 \text{ GeV}$  in run 2 (2015-)  
LHC (CERN) 2010-Now
  - Experiments: **ALICE**, CMS, and ATLAS.

I will mainly show results from LHC!

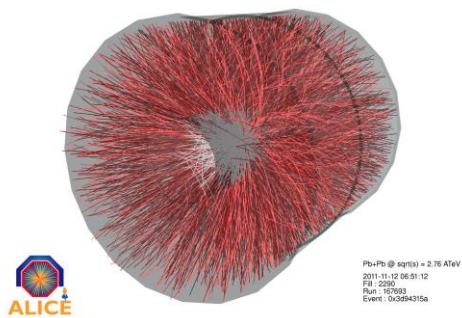
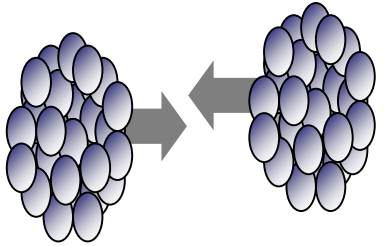




# The three systems

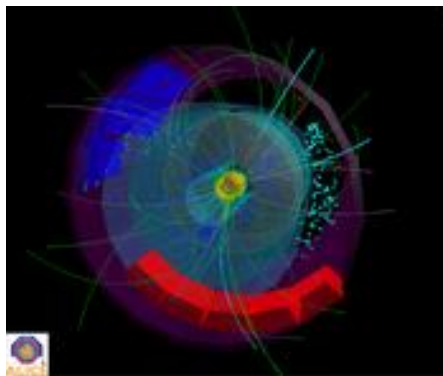
(understanding before 2012)

Pb-Pb



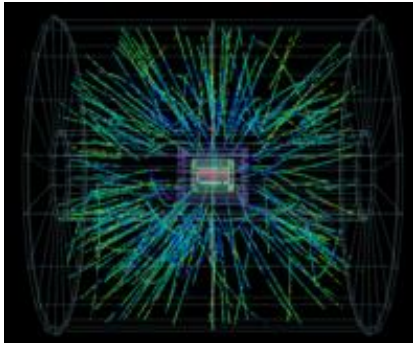
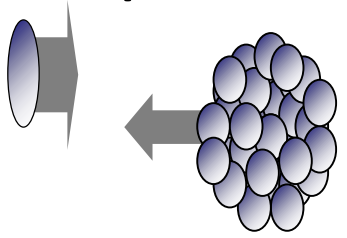
Hot QCD matter:  
This is where we expect the QGP to be created in central collisions.

pp



QCD baseline:  
This is the baseline for "standard" QCD phenomena.

p-Pb



Cold QCD matter:  
This is to isolate nuclear effects, e.g. nuclear pdfs.



# THERMAL PHOTONS

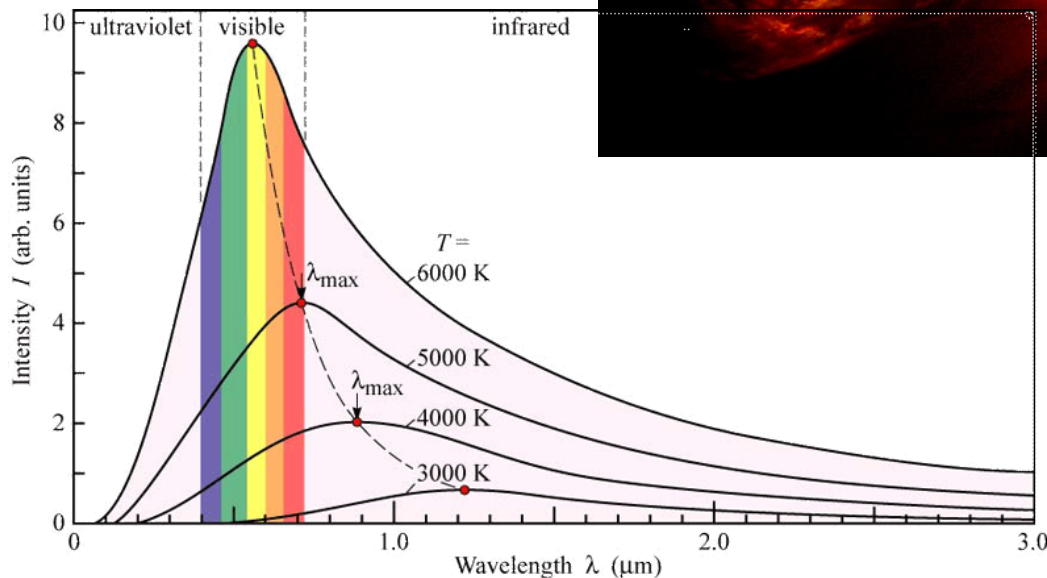
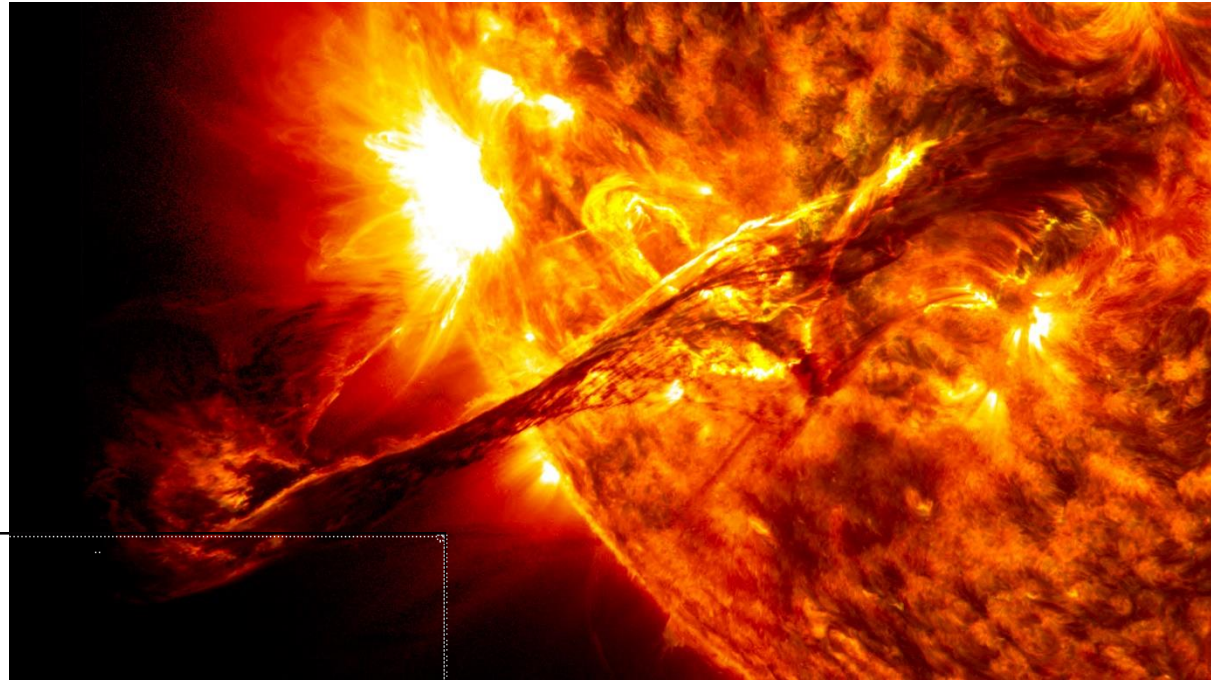
# Measuring the medium temperature

- To establish the necessary conditions for a medium one would like to measure the temperature/energy density
- There are in general three possibilities
  - Extracting the final temperature from the  $p_T$  slope and particle yields (this is the final temperature)
  - Bjorken has made a famous relation between the initial energy density and the final transverse energy (using hydrodynamics and assumptions on when a medium is formed) Phys. Rev. D27 (1983) 140-151 (2300 citations)
  - Measuring the thermal photon spectrum
    - Photons only interacts weakly with the QGP and so one can probe the early times directly (one measures the time integrated spectrum)





In a similar way as we measure the surface temperature of the sun...



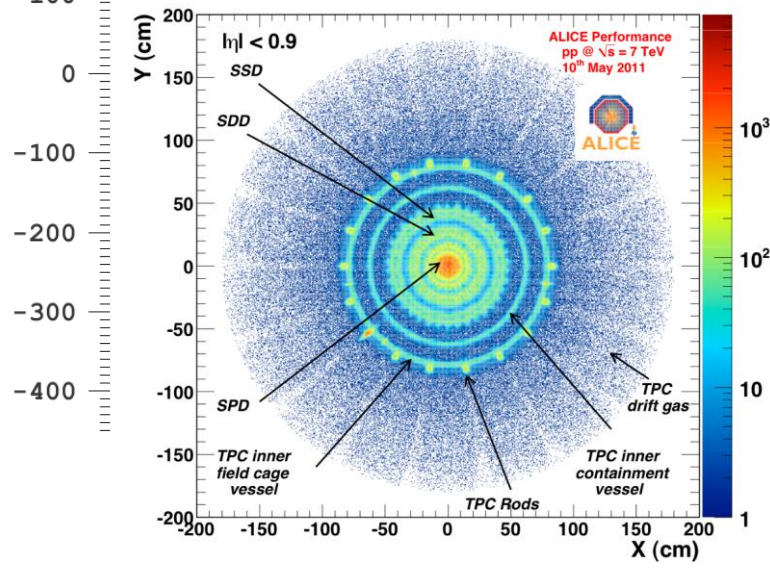
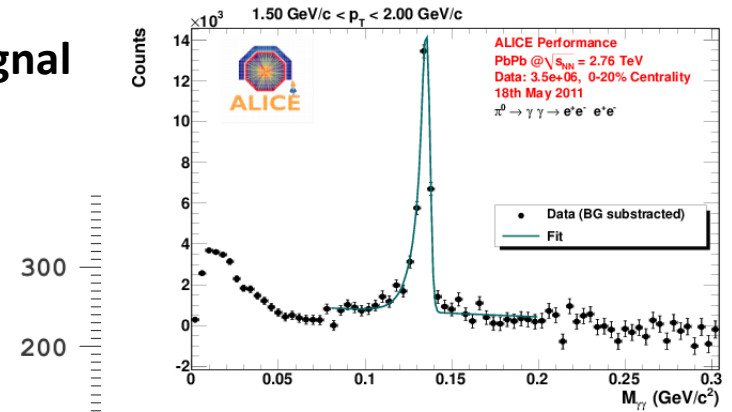
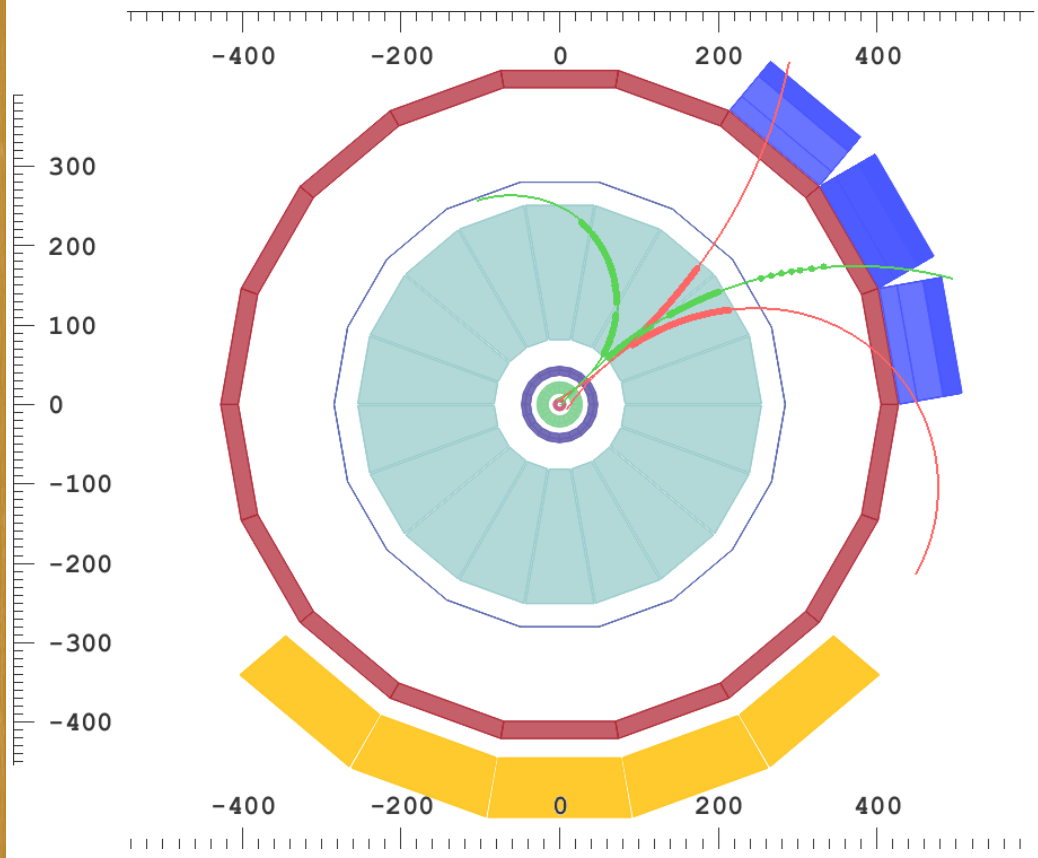
... we can measure the temperature of the QGP





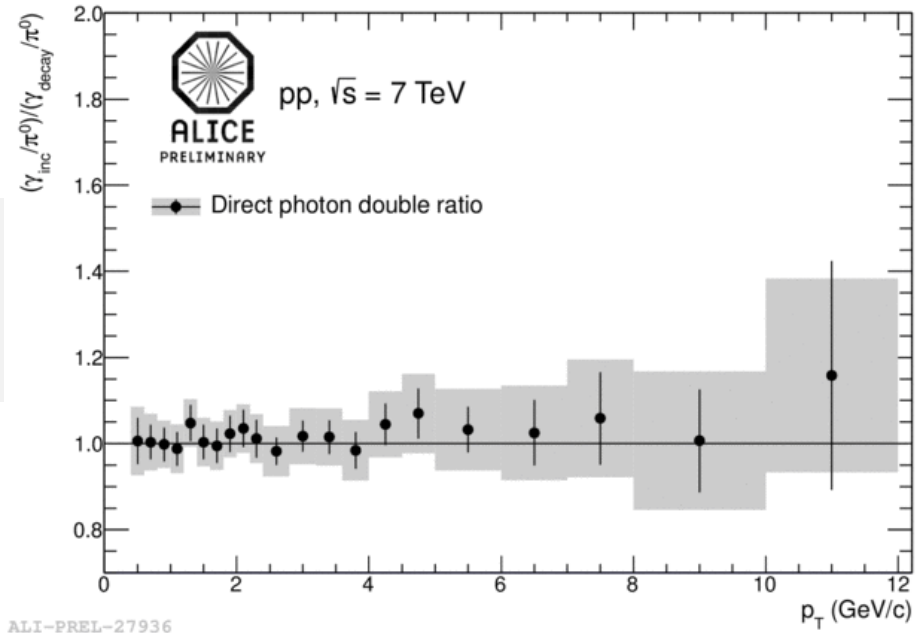
# Photon identification in the TPC via conversion

$\pi^0 \rightarrow 2\gamma$  converts to  $2*(e^- + e^+) =$  background for signal



# Direct photons in pp collisions

Double Ratio:  $\frac{\gamma_{inc}}{\pi^0} / \frac{\gamma_{decay}}{\pi^0_{param}} \approx \frac{\gamma_{inc}}{\gamma_{decay}}$   
 → cancellation of uncertainties

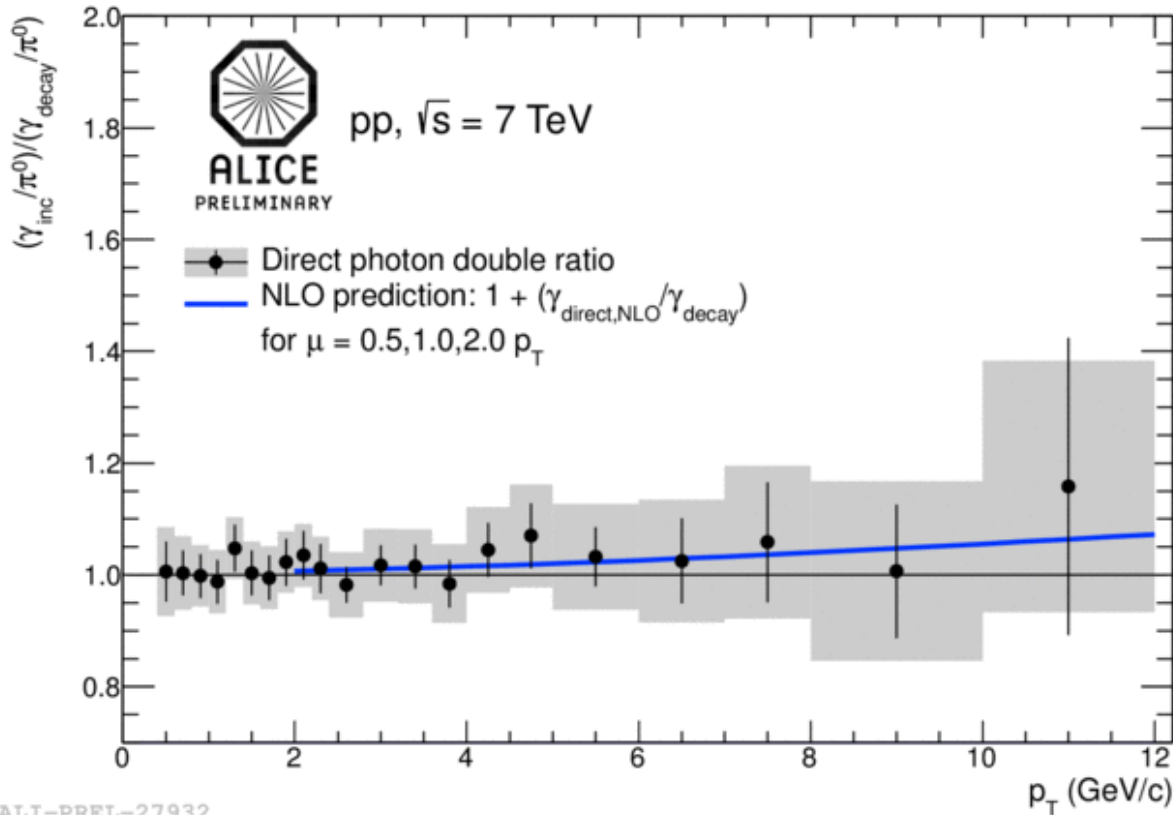


- Construct double ratio to eliminate/reduce systematics
- **Numerator** is the actual measurement
- **Denominator** is from a cocktail calculation

$$\gamma_{direct} = \gamma_{inc} - \gamma_{decay} = \left(1 - \frac{\gamma_{decay}}{\gamma_{inc}}\right) \cdot \gamma_{inc}$$



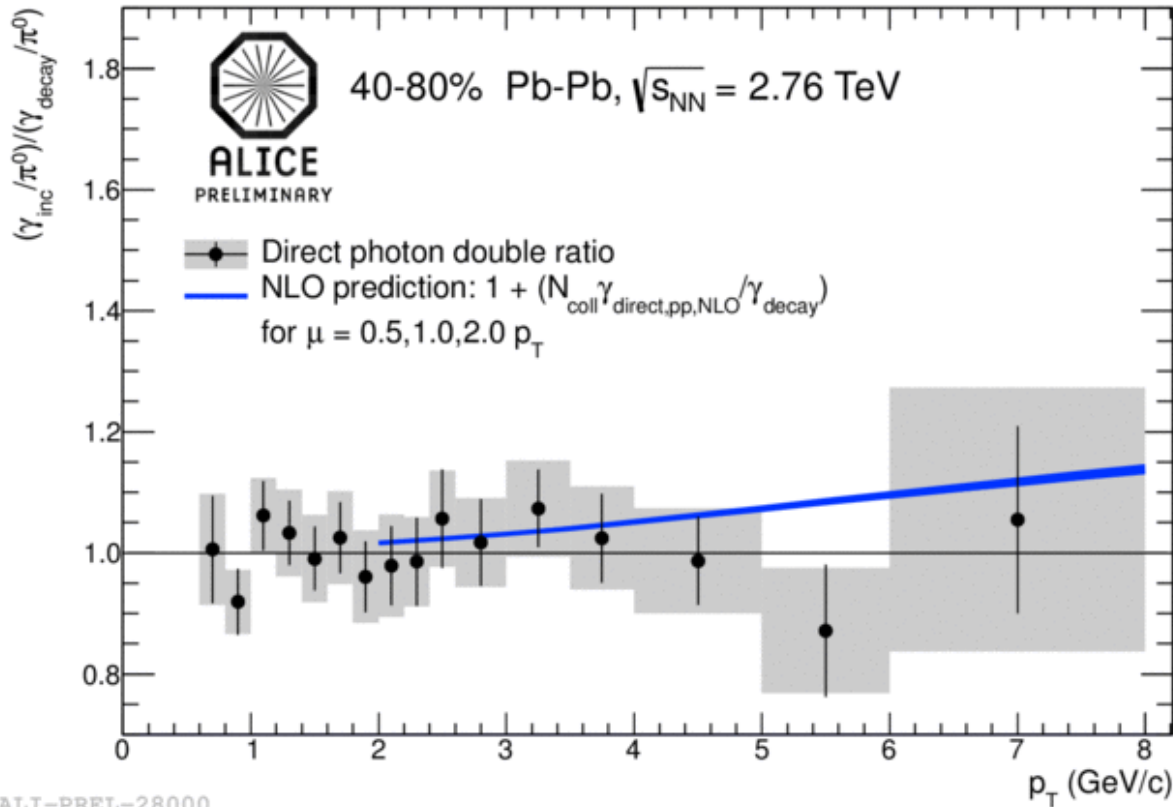
# Direct photons in pp collisions



- Comparison to pQCD NLO calculation



# Direct photons in peripheral Pb-Pb collisions

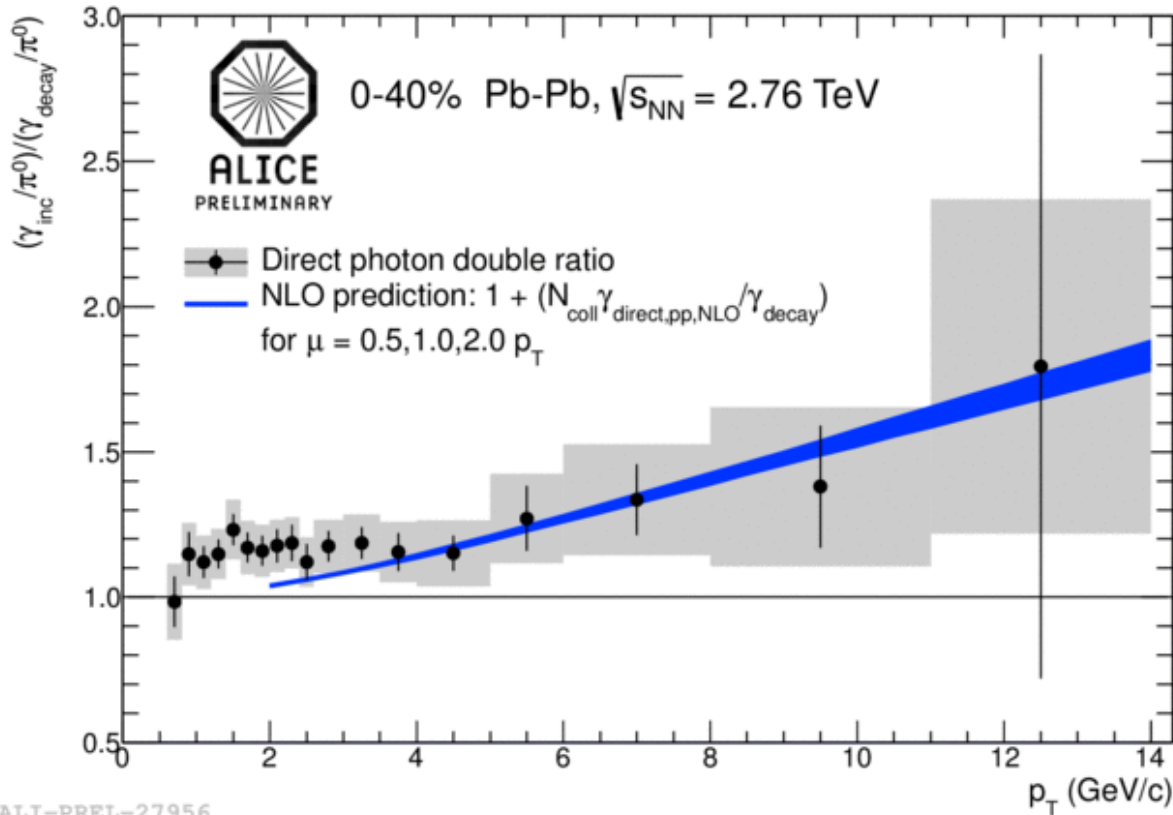


- Peripheral Pb-Pb

Consistent with only direct and decay photons



# Direct photons in central Pb-Pb collisions



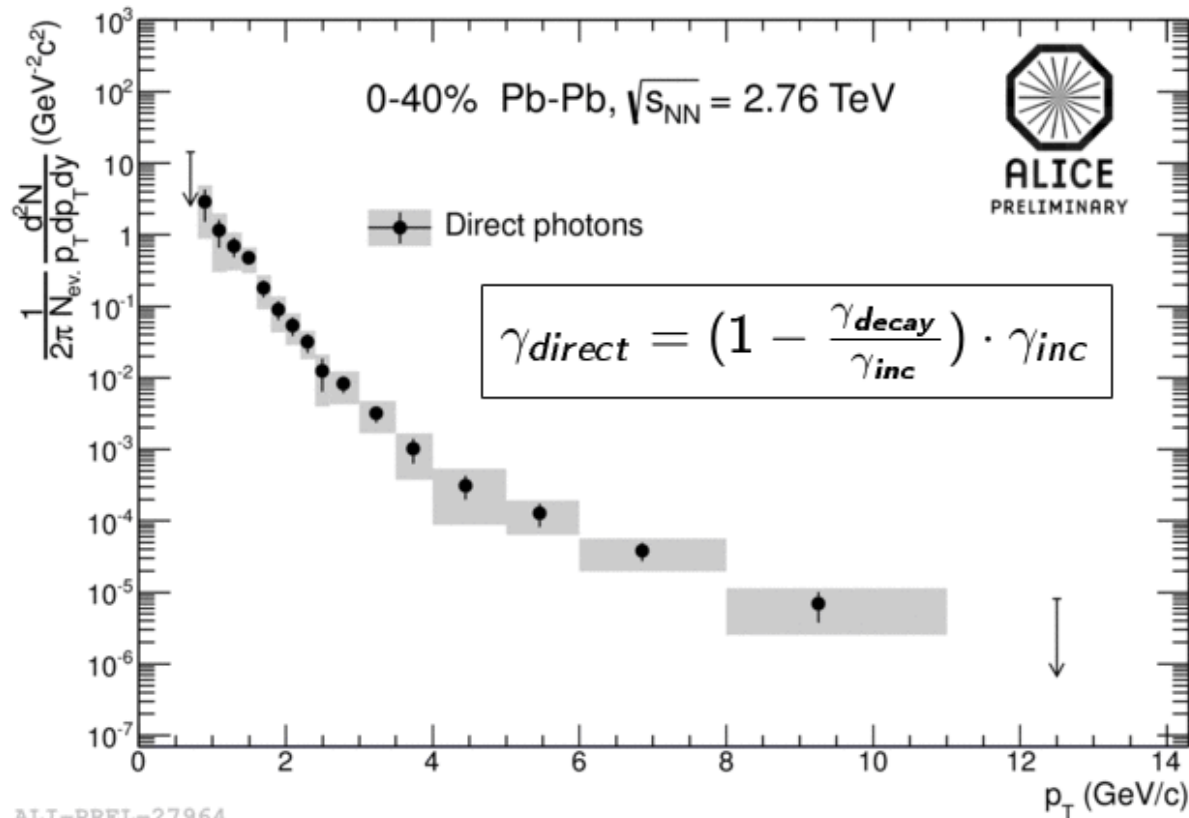
- Central Pb-Pb

Surplus of photons at high  $p_T$  is expected from hard production

Surplus of photons at low  $p_T$  are from thermal radiation from the QGP



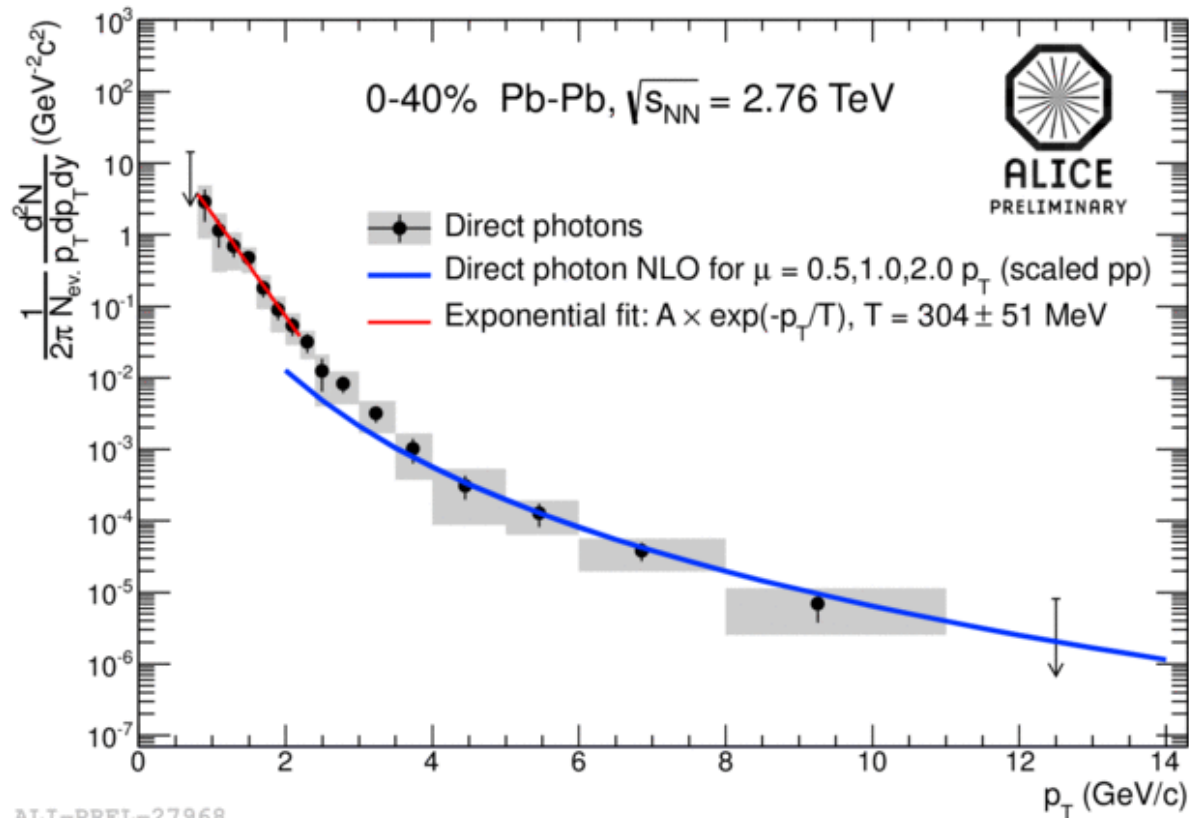
# Direct photon spectrum in central Pb-Pb



- Obtain the direct  $\gamma$  spectrum by scaling with the inclusive  $\gamma$  spectrum



# Direct photon spectrum in central Pb-Pb



- The temperature of the low  $p_T$  direct  $\gamma$  spectrum is of order 300 MeV (recall that this is a time average over the lifetime of the medium)







# The highest man-made temperature

## Highest man-made temperature

Share   



Who

**CERN, LARGE HADRON COLLIDER**

What

**$5 \times 10^{12}$  DEGREE(S) KELVIN**

Where

**SWITZERLAND**

When

**13 AUGUST 2012**

On 13 August 2012 scientists at CERN's Large Hadron Collider, Geneva, Switzerland, announced that they had achieved temperatures of over 5 trillion K and perhaps as high as 5.5 trillion K. The team had been using the ALICE experiment to smash together lead ions at 99% of the speed of light to create a quark gluon plasma – an exotic state of matter believed to have filled the universe just after the Big Bang.





# HARD PROBES

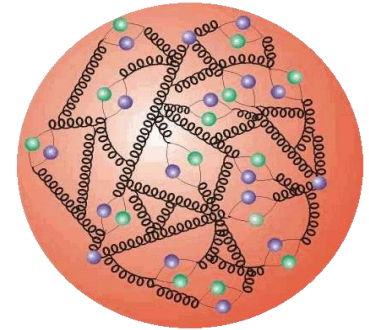
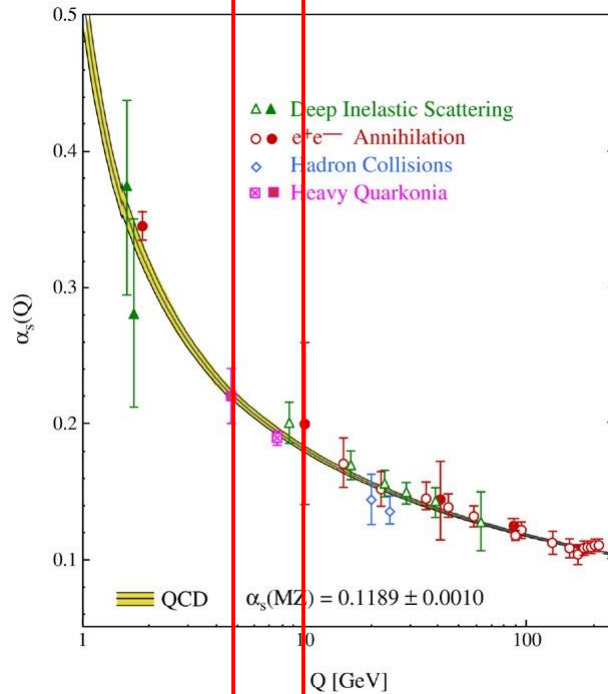
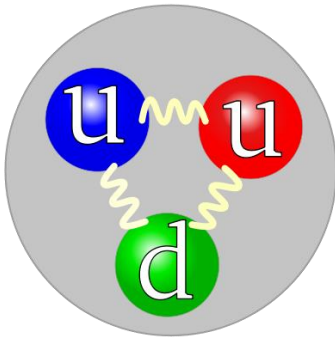
# What happens when we collide pp and Pb-Pb?

- Think for a few minutes
- Discuss with your neighbors for a few minutes



# What happens when we collide pp and Pb-Pb

- 2 answers!



Non-perturbative physics  
 (knows the equations but not how to  
 Solve them)  
 Bulk properties (=medium)

Perturbative physics  
 (theoretical predictions)  
 Rare jets (=probes)





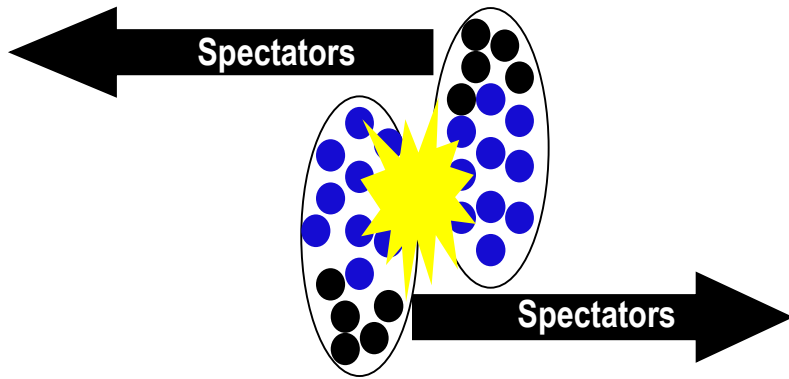
# Hard probes in heavy-ion collisions

- The cross sections for Hard probes (HP) can in principle be calculated using pQCD (+ pdfs and FFs). In practice they need to be measured to achieve the best systematic precision.
- The cross sections can easily be extrapolated from pp to Pb-Pb collisions using binary scaling (next slide) unless there are “nuclear” effects
  - HP allow the precise study of nuclear effects
- Penetrating probes
- Strongly interacting probes
  - Heavy quarks (quarkonia), jets

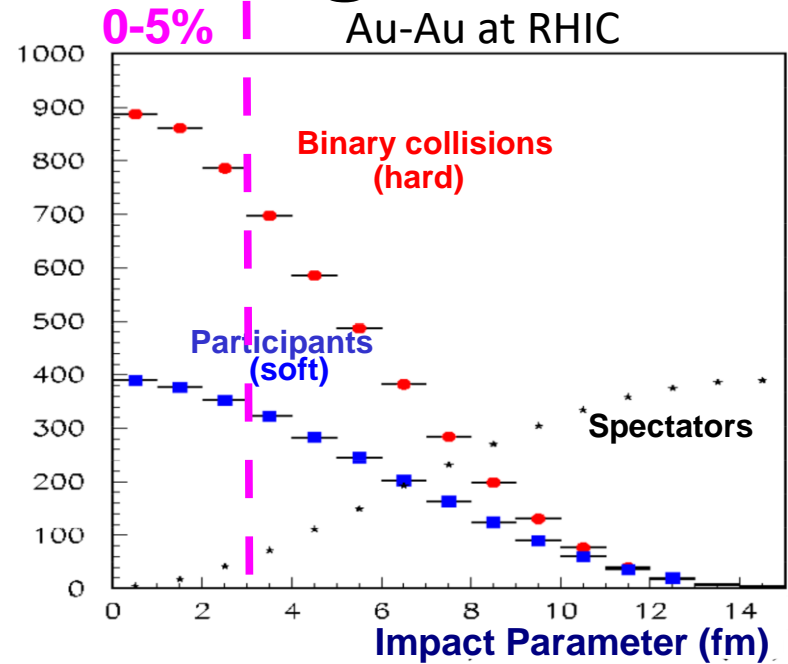


# The collision geometry and binary scaling

- Centrality (ex. for Au+Au):

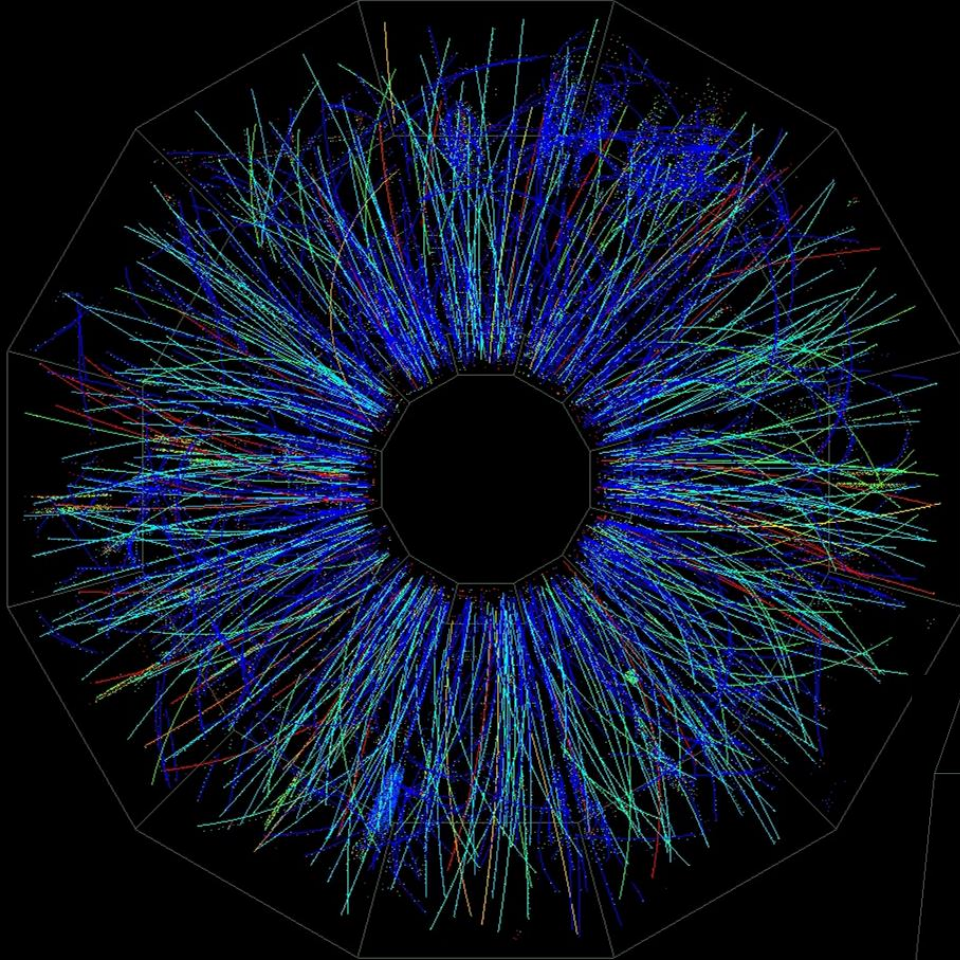


**Participants** =  $2 \cdot 197 - \text{Spectators}$



- The “medium” energy is proportional to the # of participant (Npart)
- The number of parton-parton (quark-quark, quark-gluon, gluon-gluon) is proportional to the # of binary collisions (Nbin)
- Example:
  - 6 participant
  - 8 binary collisions
  - (pp has 2 participant and 1 binary collision)

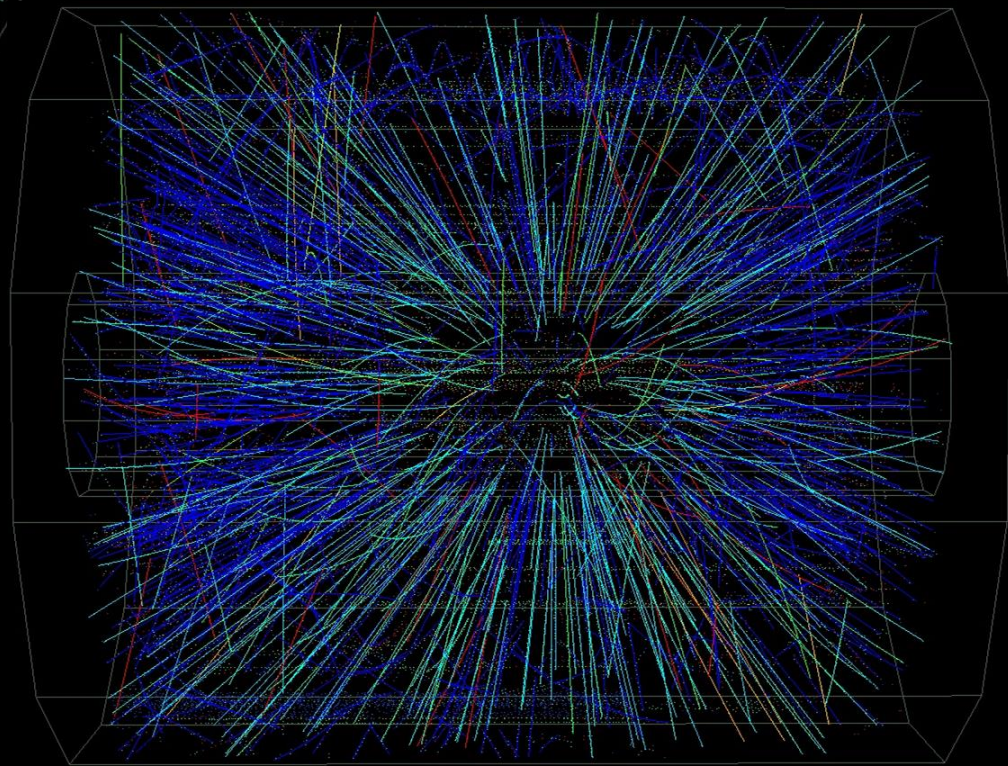




## Peripheral Event

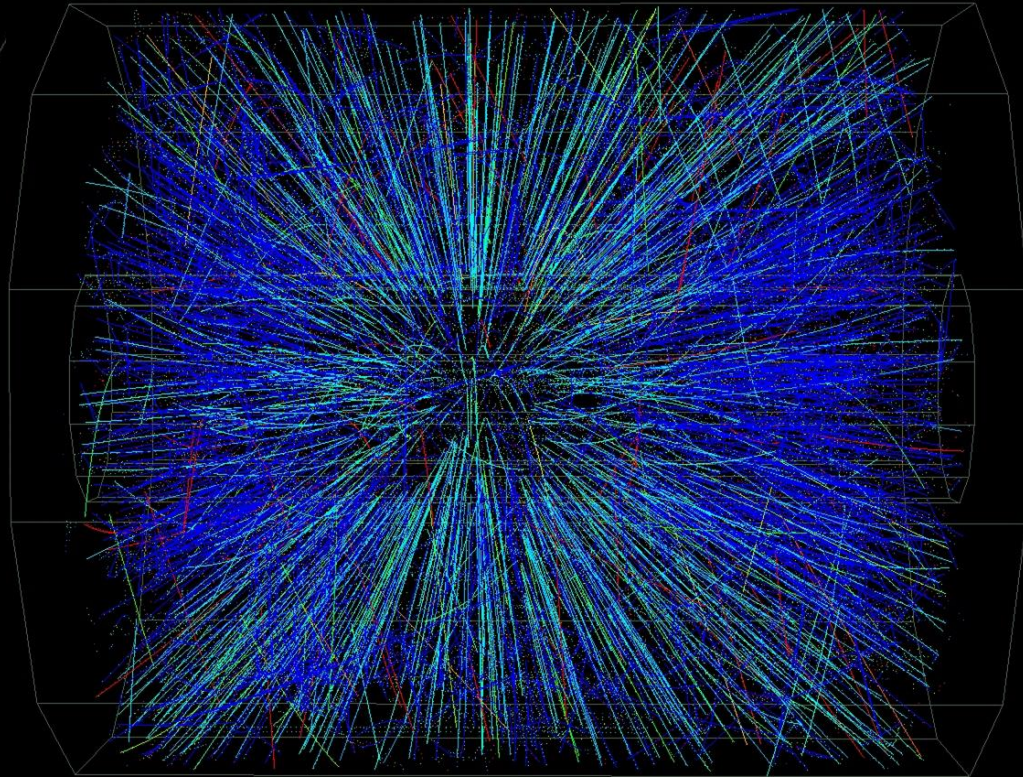
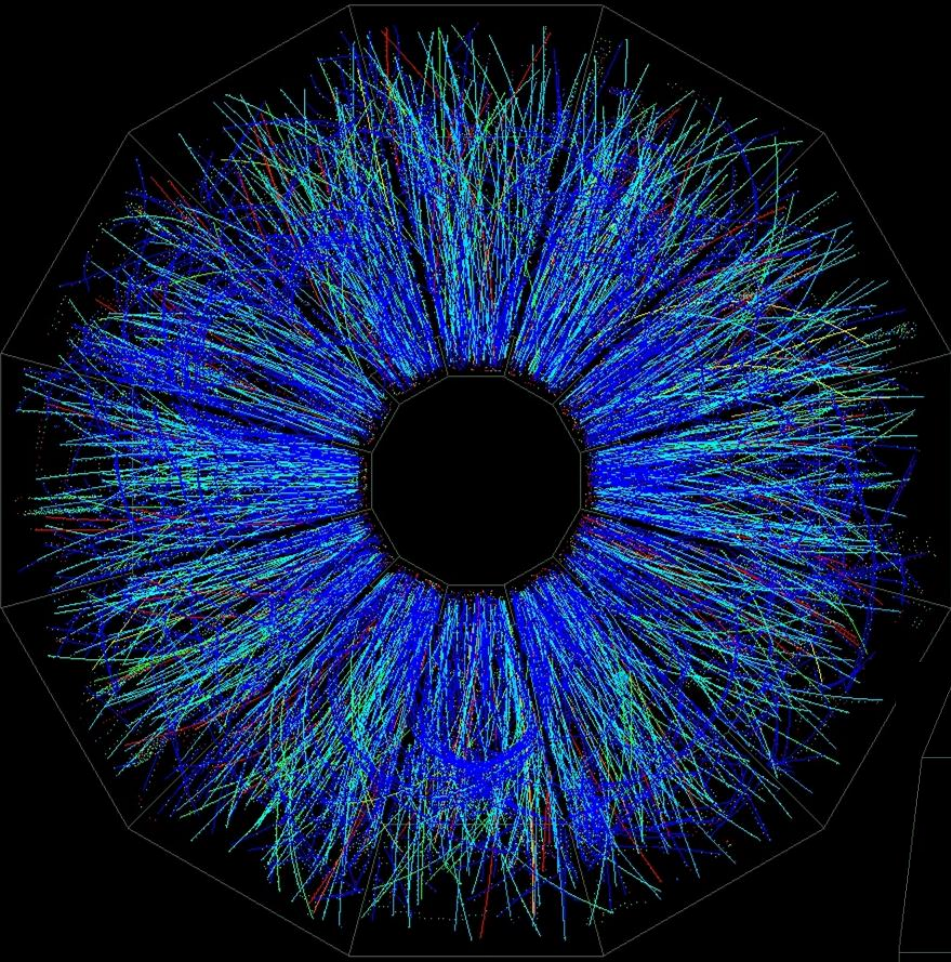
From real-time Level 3 display.

color code  $\Rightarrow$  energy loss

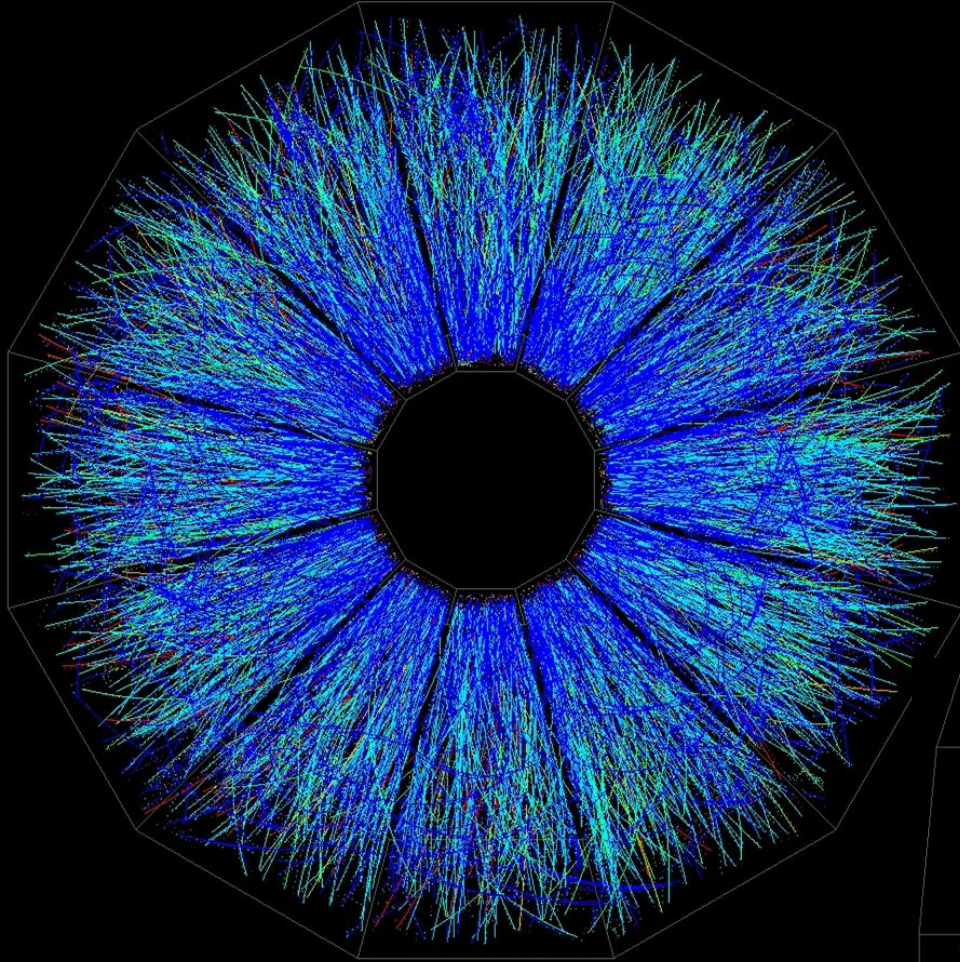


## Mid-Central Event

From real-time Level 3 display.

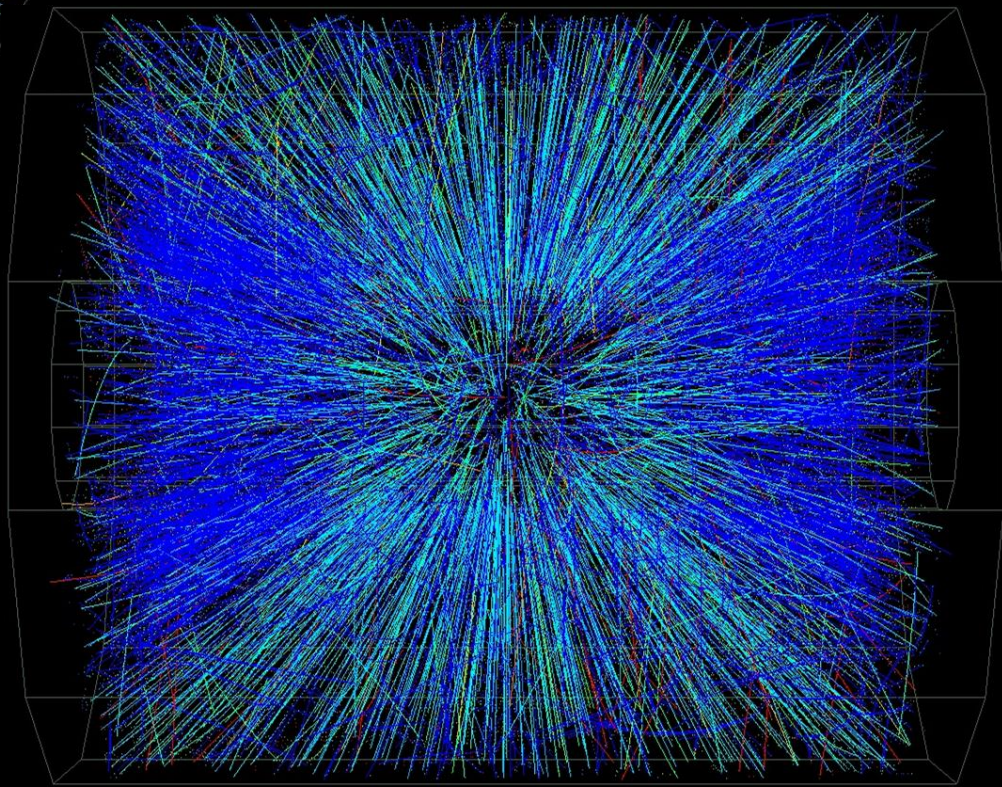






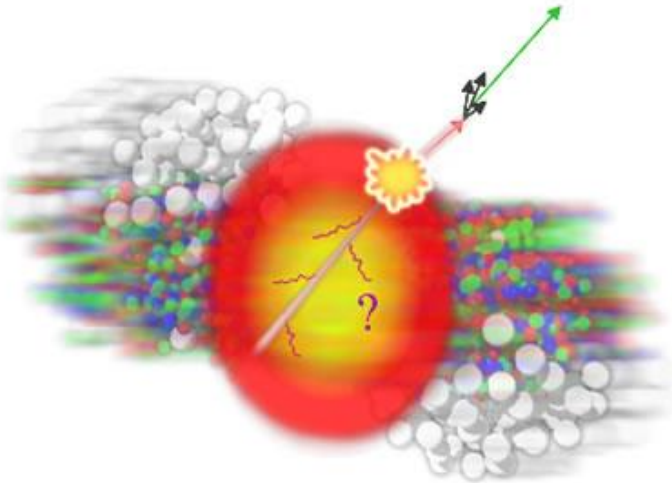
## Central Event

From real-time Level 3 display.



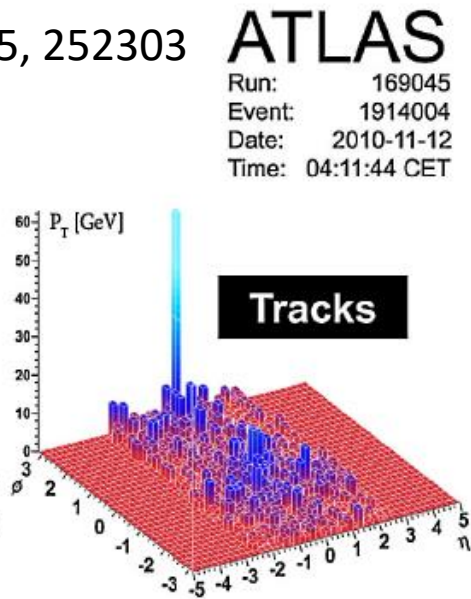
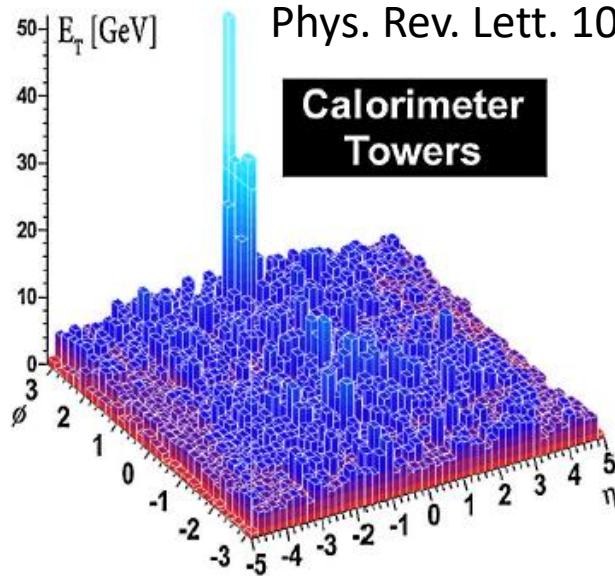
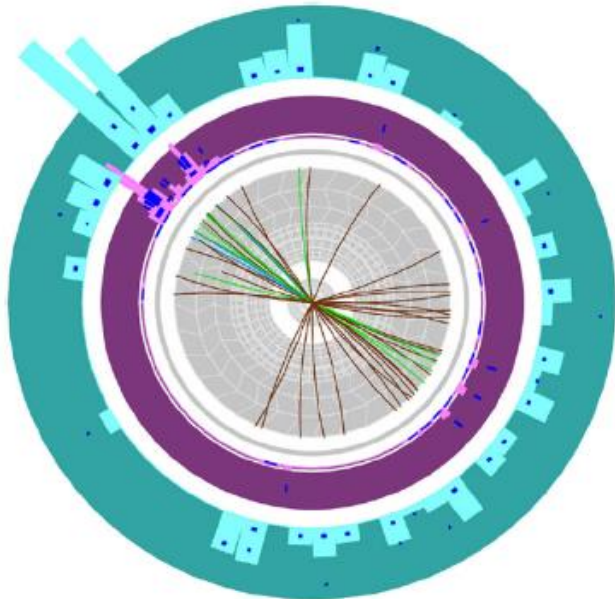


# Jet quenching in the QGP



For color charges one expects energy losses of GeV/fm vs MeV/cm for electric charges in normal atomic matter.  
 10,000,000,000,000,000 ( $10^{16}$ ) times larger energy loss.

Heavy-ion physics and the QGP (P. Christiansen, Lund)



**ATLAS**  
 Run: 169045  
 Event: 1914004  
 Date: 2010-11-12  
 Time: 04:11:44 CET

# The nuclear modification factor $R_{AA}$ (1/2)

$$R_{AA} = \frac{d^2 N^{AA} / dp_T d\eta}{\langle T_{AA} \rangle d^2 \sigma^{pp} / dp_T d\eta}$$

$$\langle T_{AA} \rangle \sigma^{pp} = \langle N_{\text{coll}} \rangle$$

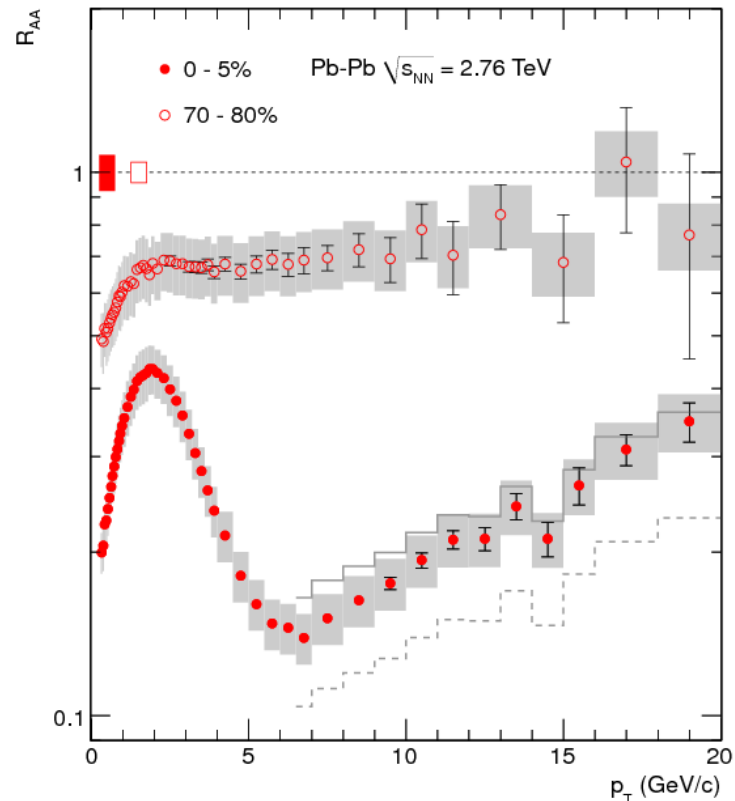
$N_{\text{coll}}$  is the number of binary collisions

For perturbative QCD processes:

$R_{AA} < 1$ : suppression

$R_{AA} = 1$ : no nuclear effects

$R_{AA} > 1$ : enhancement



Phys. Lett. B 696 (2011)

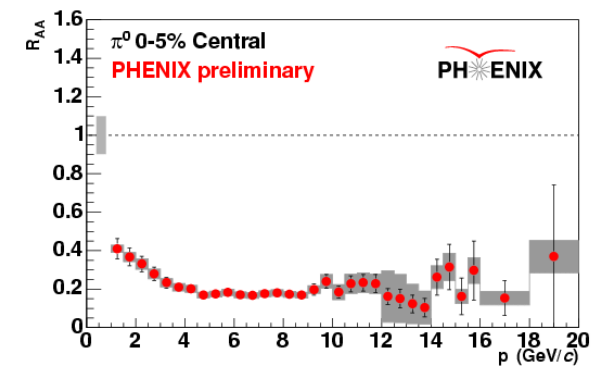
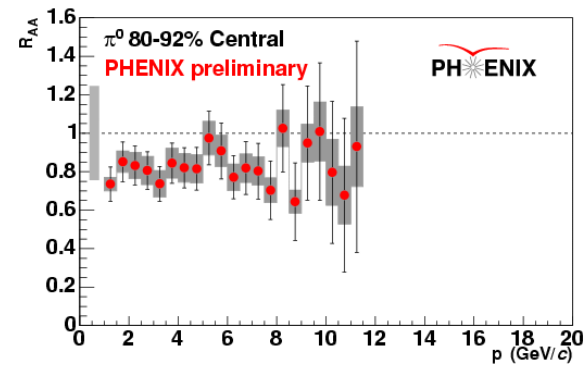
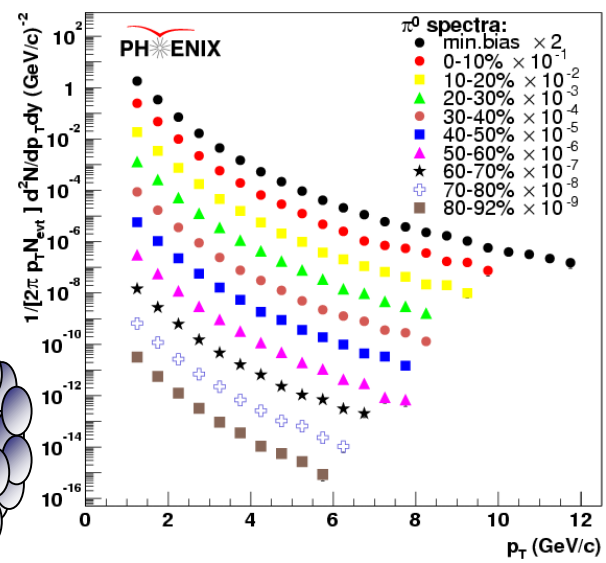
So  $R_{AA}$  is a way to quantify the nuclear effects!



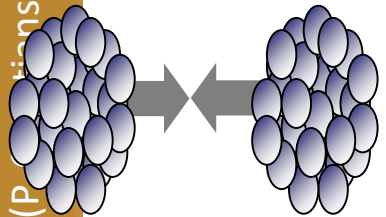


# The nuclear modification factor $R_{AA}$ (2/2)

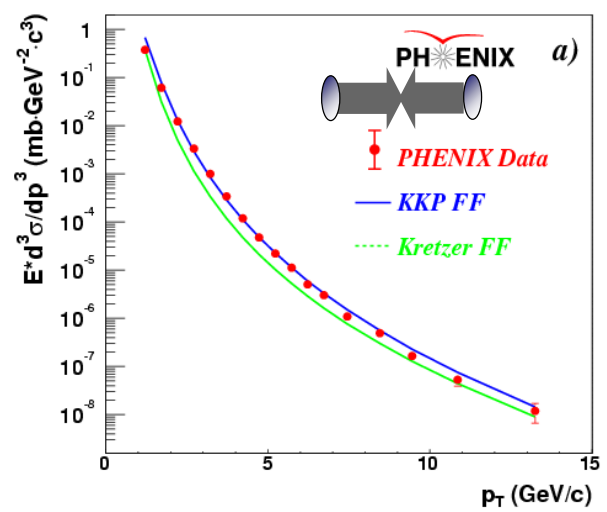
$$R_{AA} = \frac{d^2 N^{AA} / dp_T d\eta}{\langle T_{AA} \rangle d^2 \sigma^{pp} / dp_T d\eta}$$



Heavy-ion physics and the QGP (P. Thorenson, Lund)



## $N_{coll}$





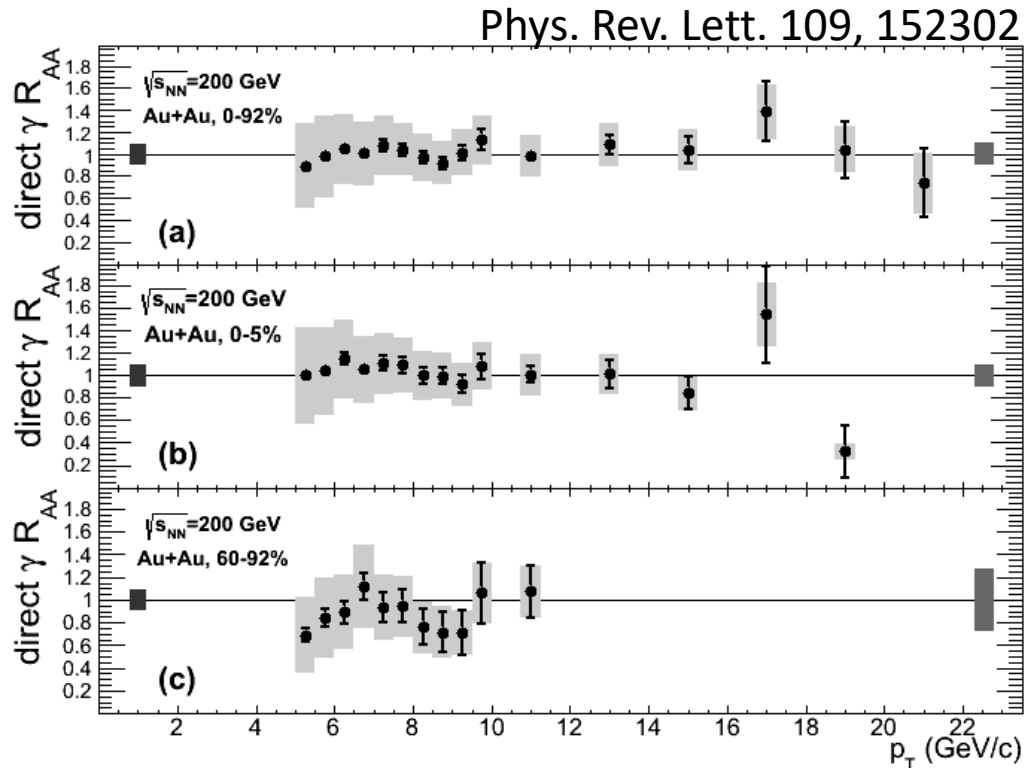
# Standard candles



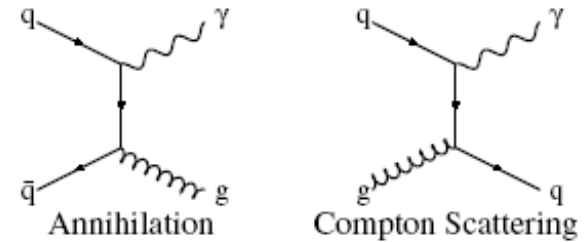
We need penetrating hard probes to validate the binary scaling

What could that be?

# Ncoll scaling for direct photons



## Source of direct photons



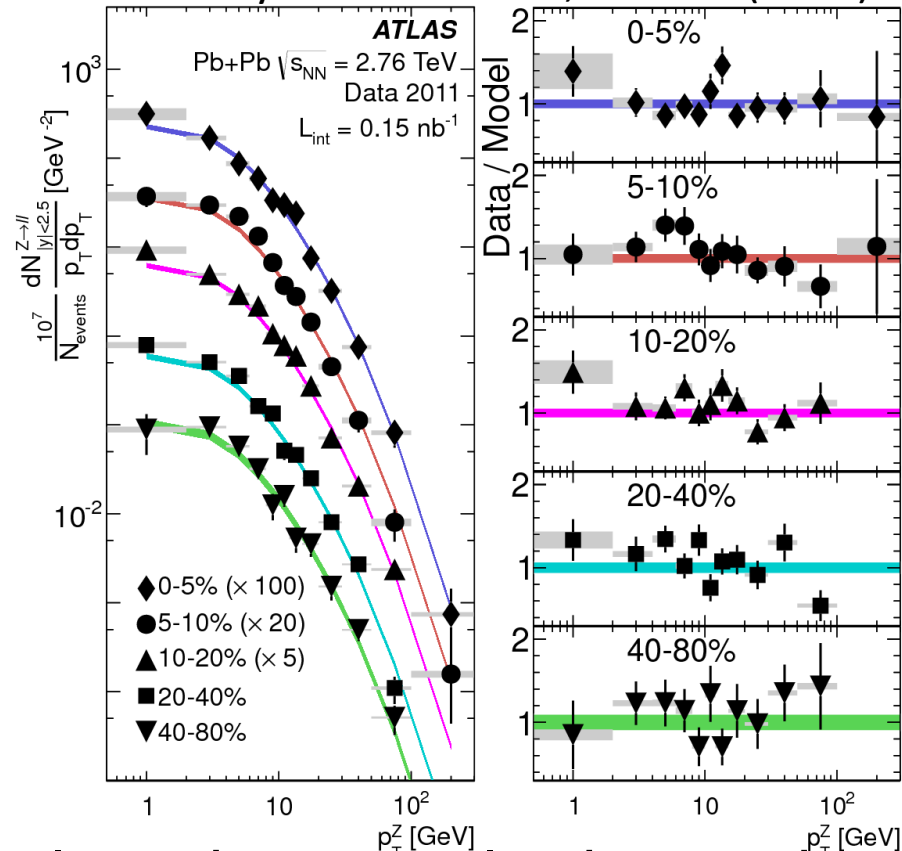
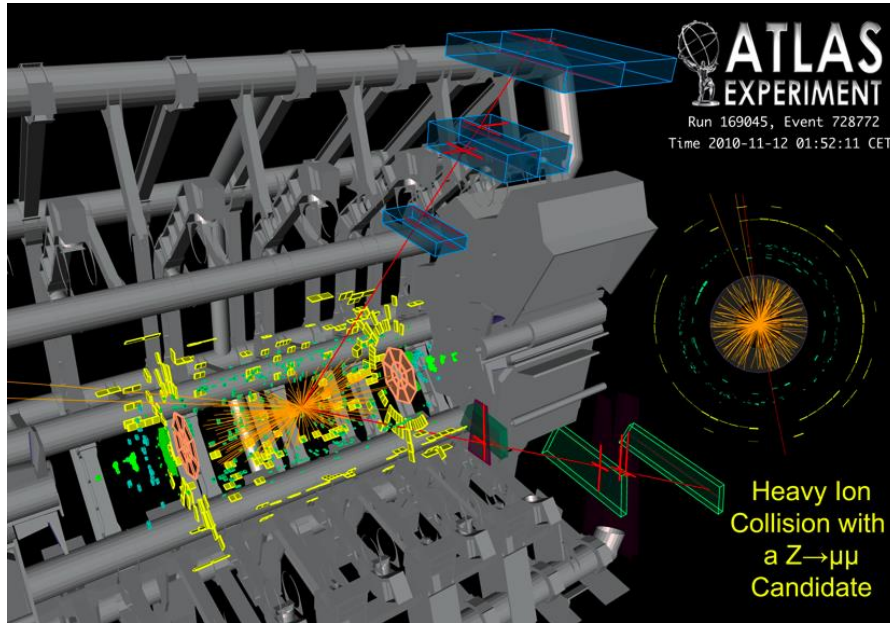
- Direct photons does not interact with final state hadronic matter and the results confirm binary scaling of hard processes!
- The best standard candle at RHIC





# New “standard candle” at LHC: ATLAS measures Z bosons

Phys. Rev. Lett 110, 022301 (2013)

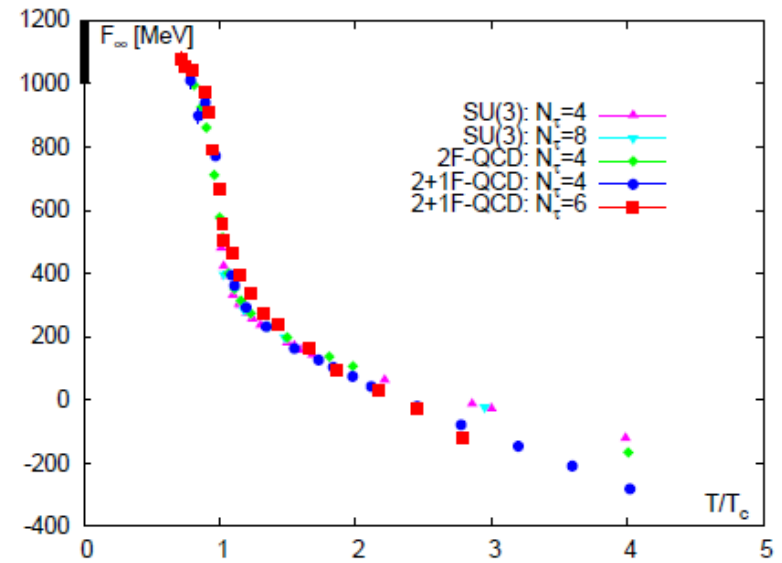
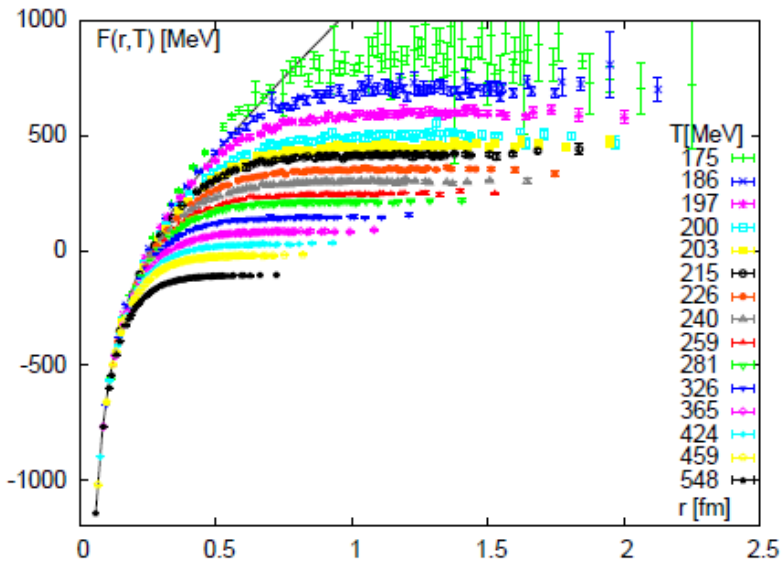


- The Z does not interact strongly and so can also be used to check binary scaling at LHC
- Also the W have been used at LHC



# Heavy quarks: I will only focus on Quarkonia

Lattice QCD results for the heavy quark potential (free energy): arXiv:0710.0498



Lattice QCD predicts that the long range force will be screened in the plasma  
Bound state properties will change and some states will disappear / melt

Original idea:

“J/ $\psi$  Suppression by Quark-Gluon Plasma Formation” by T. Matsui and H. Satz

Phys.Lett. B178 (1986) 416

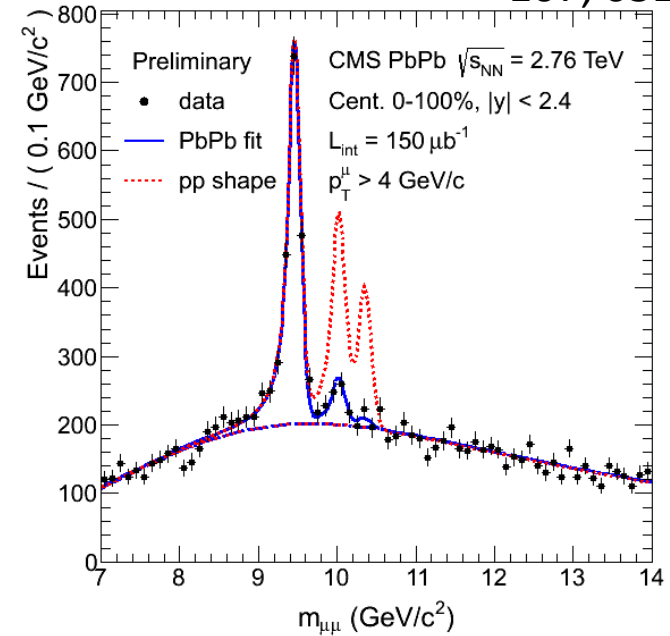
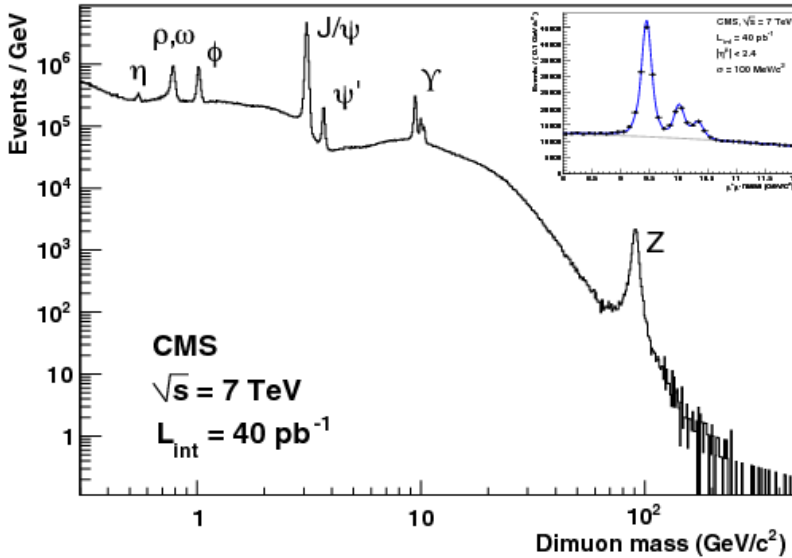
Cited by 2108 records



# LHC has delivered

$\Upsilon$ : 1S 2S 3S

Published in PRL  
107, 052302



CMS has proven to be a marvelous detector for bottomonium  
ALICE can complement by going to lower  $p_{\text{T}}$  for charmonium

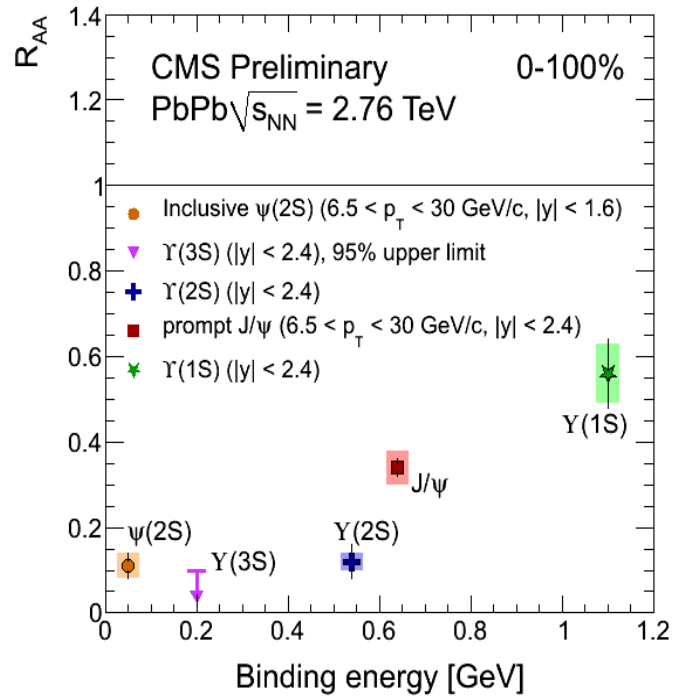
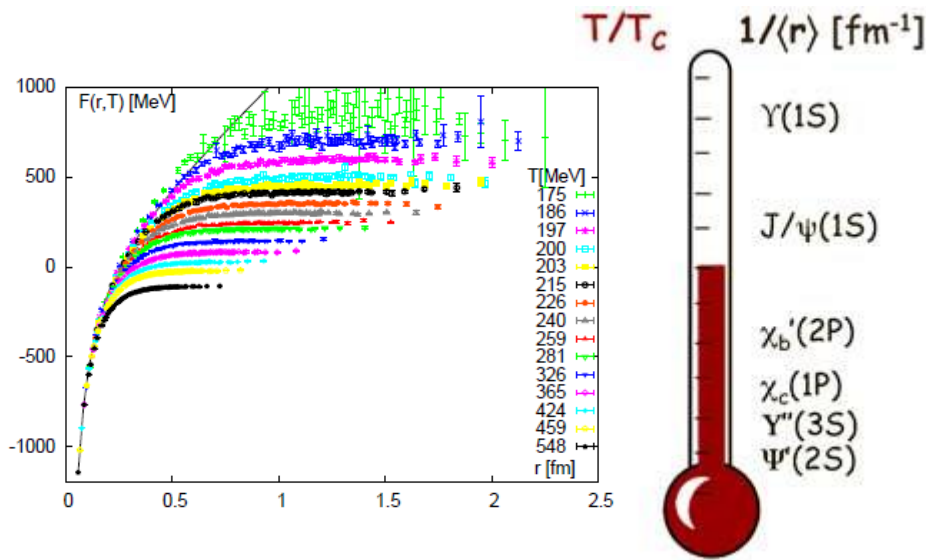
These textbook results demonstrate the amazing capabilities  
of LHC detectors





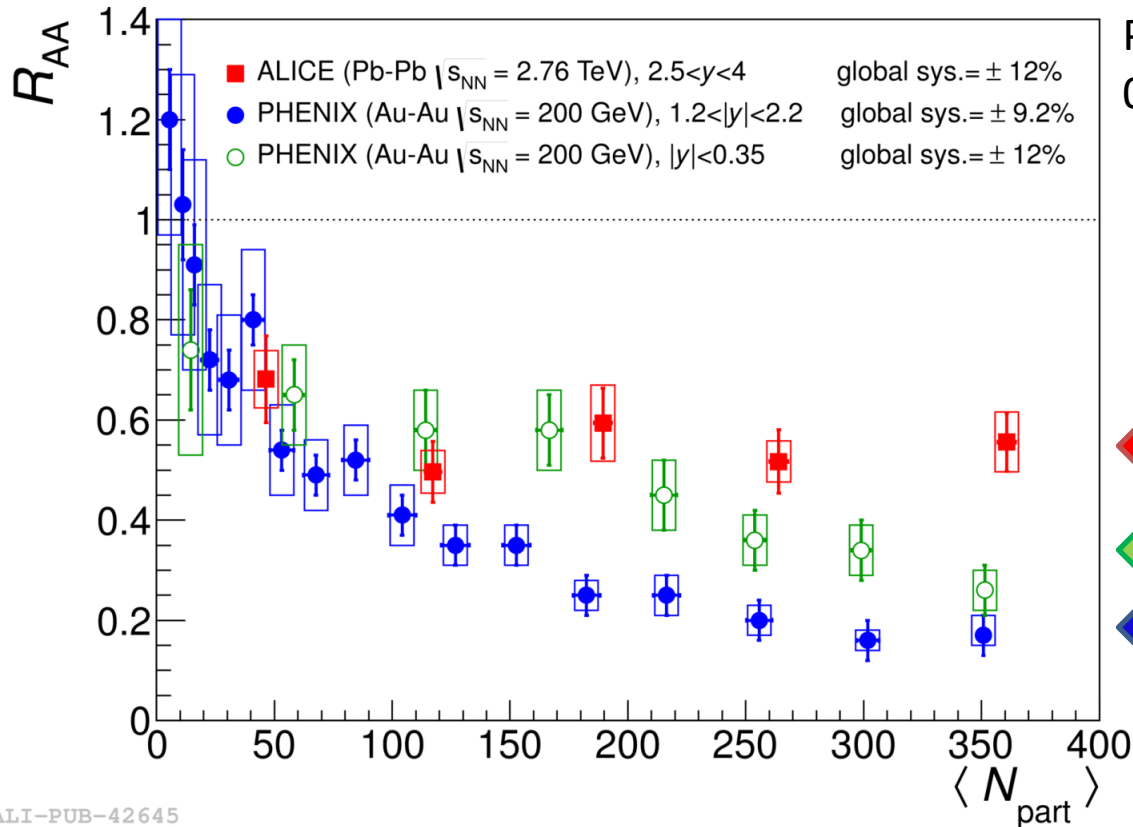
# Suppression of heavy quarkonia can work as a thermometer

Note:  $6.5 < p_T < 30$  GeV for  $J/\psi$  and  $\psi(2s)$






Suppression qualitatively depends on binding energy as predicted

# J/ $\psi$ results from ALICE: Evidence for regeneration



Phys. Rev. Lett. 109,  
072301 (2012)

ALICE can go  
down to  $p_T \sim 0$ .

 LHC  
 RHIC (mid)  
 RHIC (fwd)

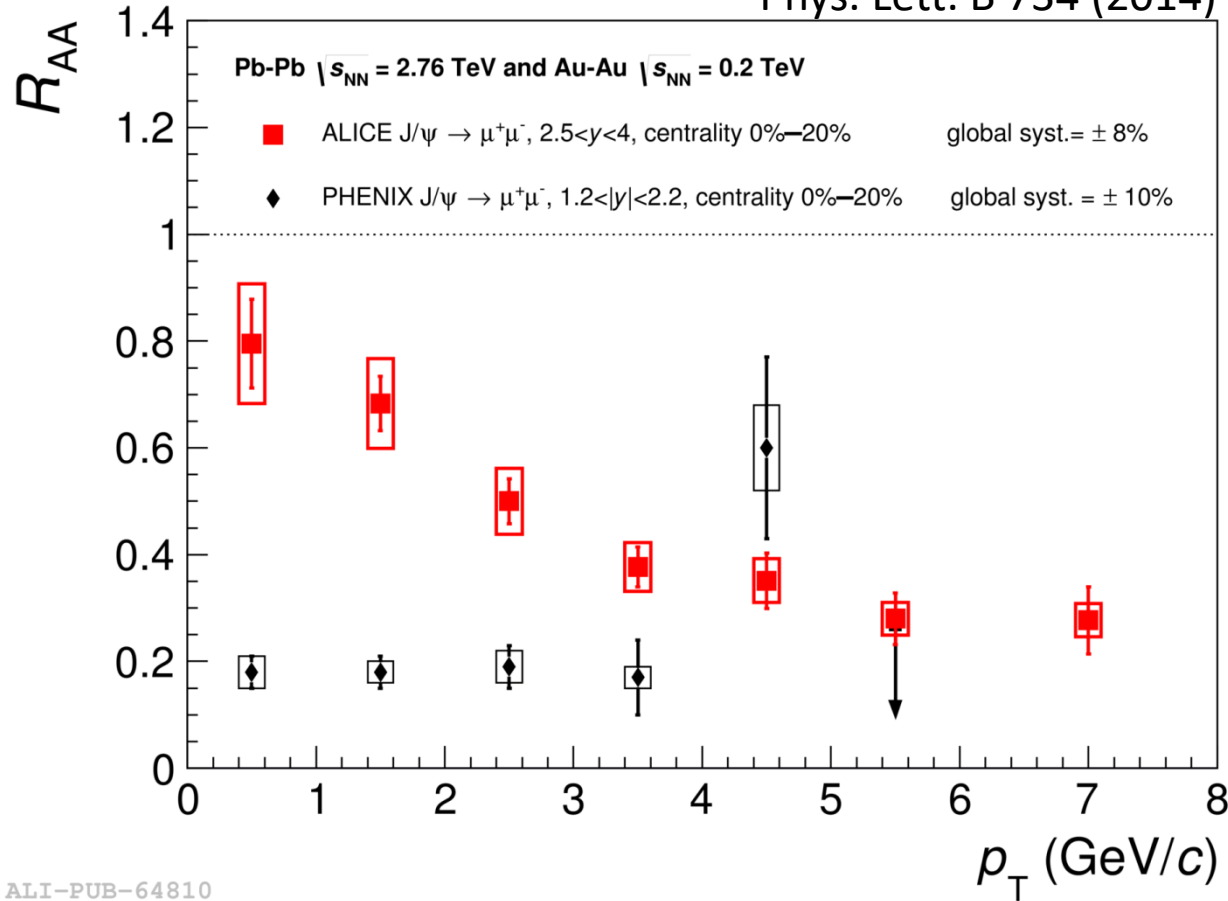
The suppression of J/ $\psi$  at LHC is less than at RHIC!

Understanding: due to the large cross section "random" charm and anti-charm quarks combine (see next slides)



# J/Psi differential results

Phys. Lett. B 734 (2014) 314-327

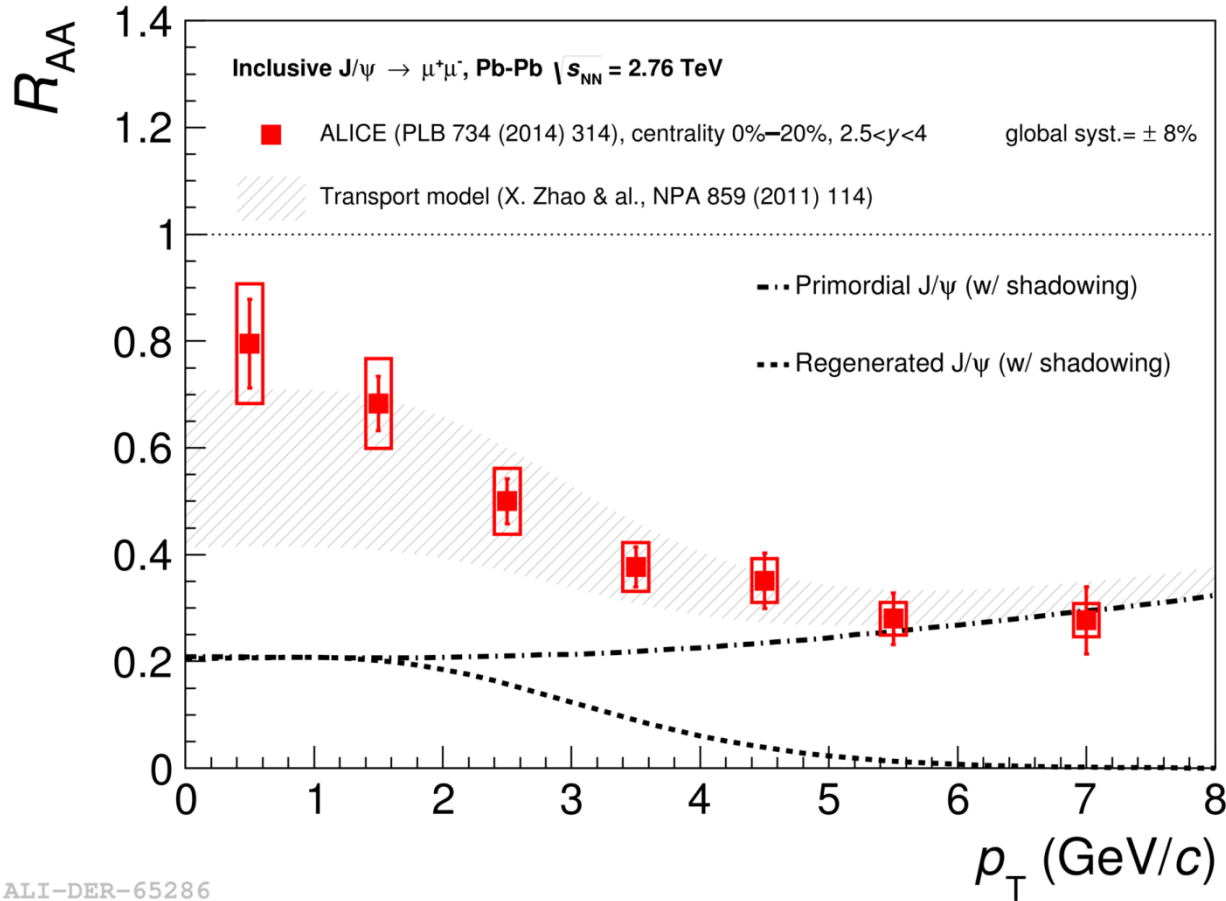


The difference in suppression is visible at low  $p_T$





# J/ $\Psi$ differential results

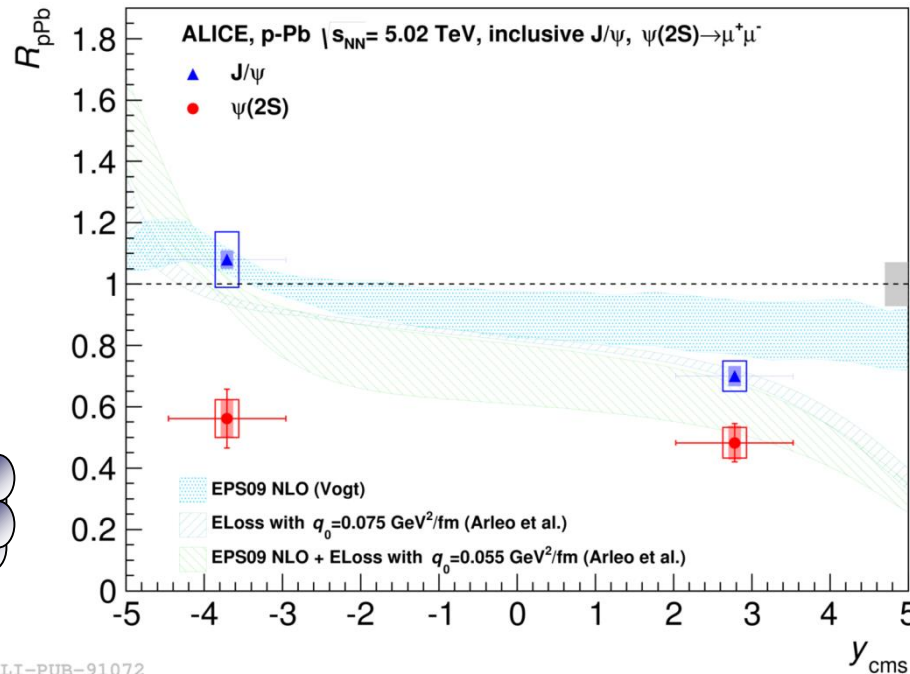
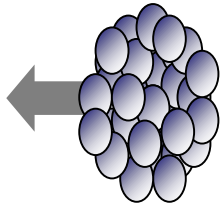


The difference can be explained by  $J/\psi$  regeneration (much larger charm Xsection at LHC)

(This effect was already proposed before the RHIC results came out.)

# Caveat: cold nuclear matter effects are not understood

lead  
rapidity



JHEP12(2014)073

proton  
rapidity

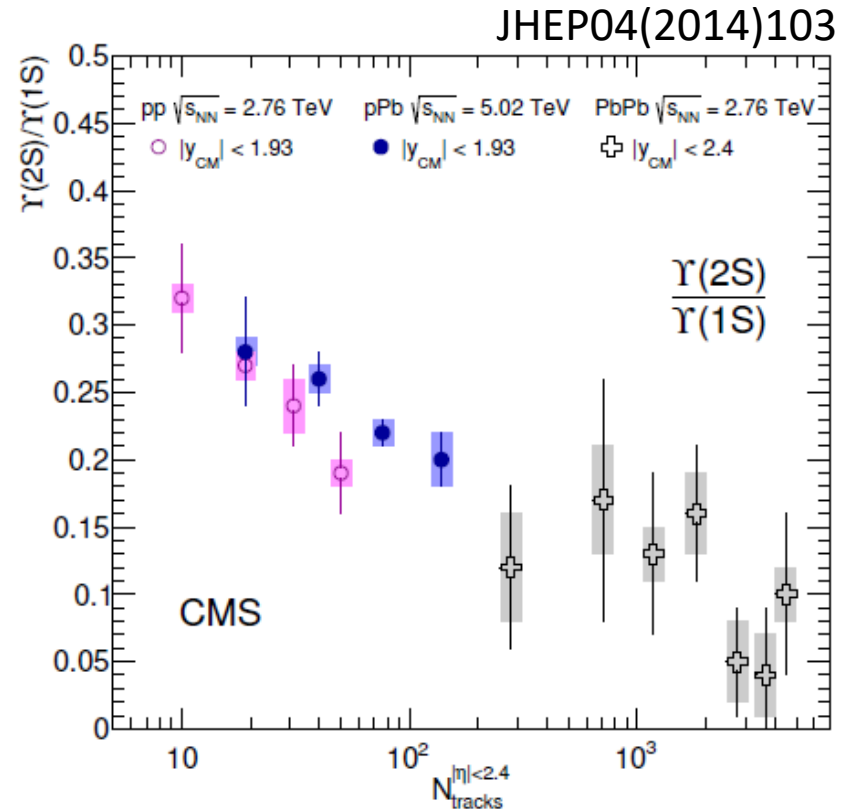
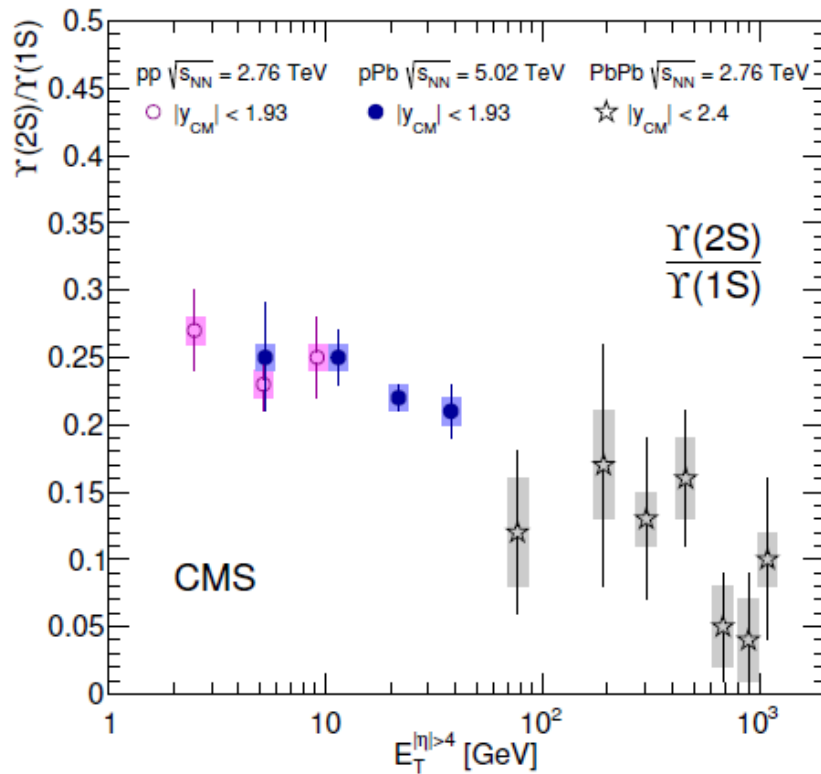


ALI-PUB-91072

For  $J/\psi$  we have reasonable good description: energy loss really just means that due to collisions with other nucleons the rapidity of the  $J/\psi$  is shifted towards the Pb-nuclei. But why is the 2S state extremely suppressed in p-Pb (at the rapidity) where the 1S state is a bit enhanced?



# The relative production of $\Upsilon$ states is affected by event activity (?)

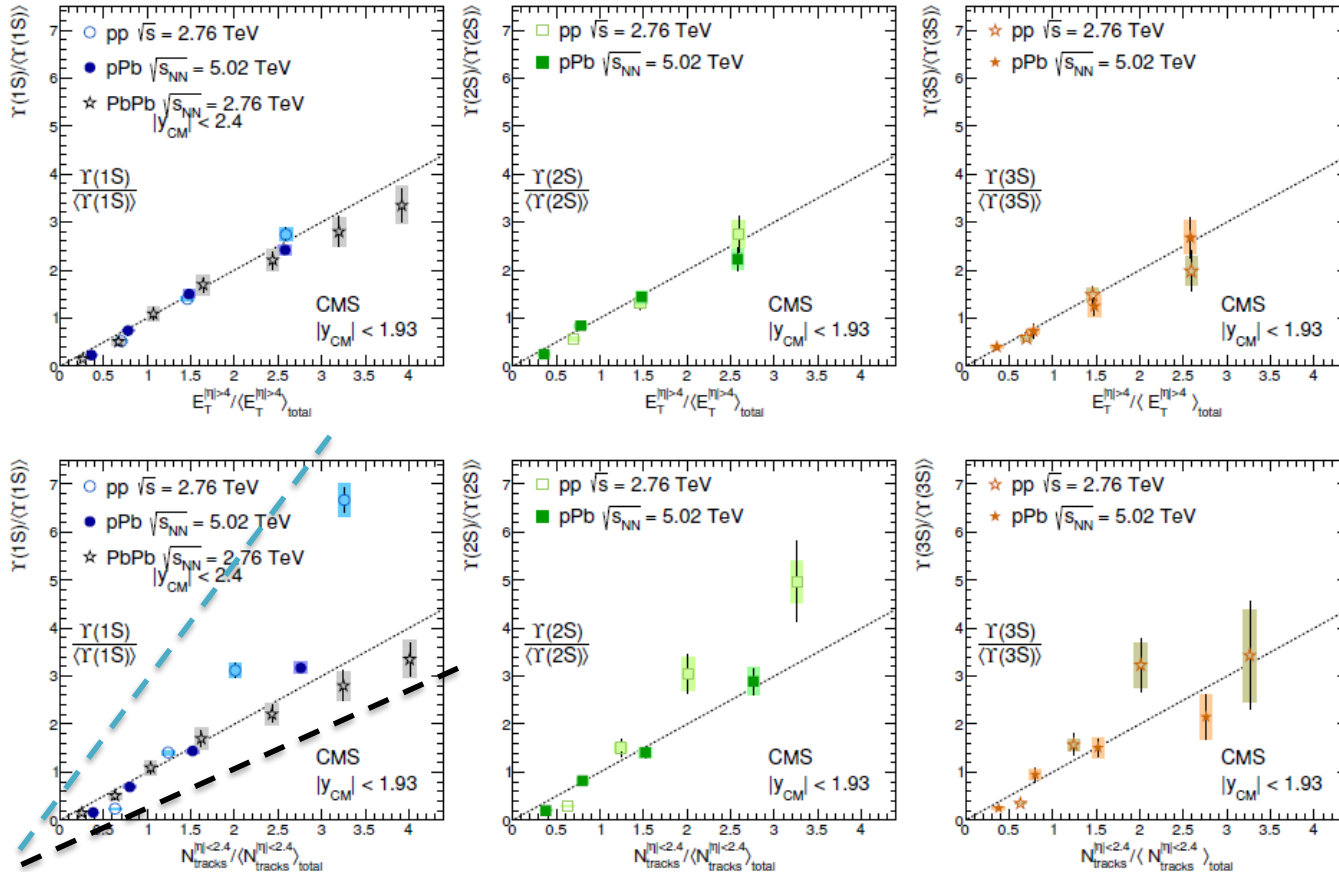


Also for  $\Upsilon$  we observe varying trends for the different states even for pp collisions!  
 (More info on next slide)





# In small systems we observe enhancement with event activity



JHEP04  
(2014)103

In this case we observe different trends with event activity  
 Caveat: should study binary scaling (but this does not work for pp)



# Summary of quarkonia results

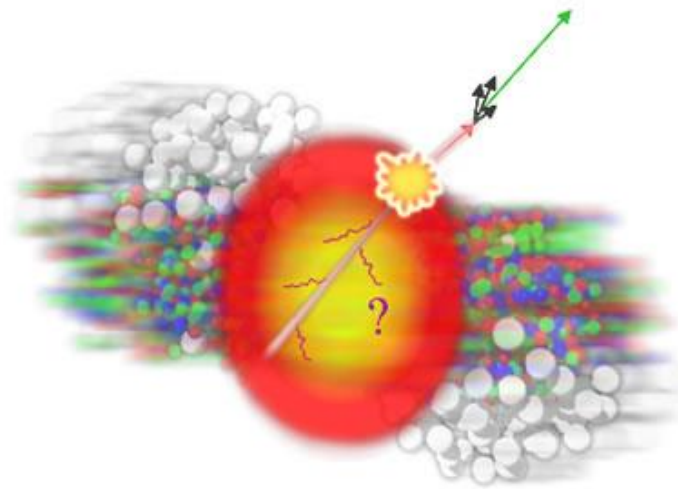
- The qualitative results are in agreement with what one expects for a Quark Gluon Plasma
- The advantage of LHC is that all quarkonia states have been measured making the measurements less sensitive to assumptions about feed-down corrections
- The quantitative picture is complicated and there are many issues that are even interesting for pp phenomenology





# JETS

# Jets as probes of the QGP

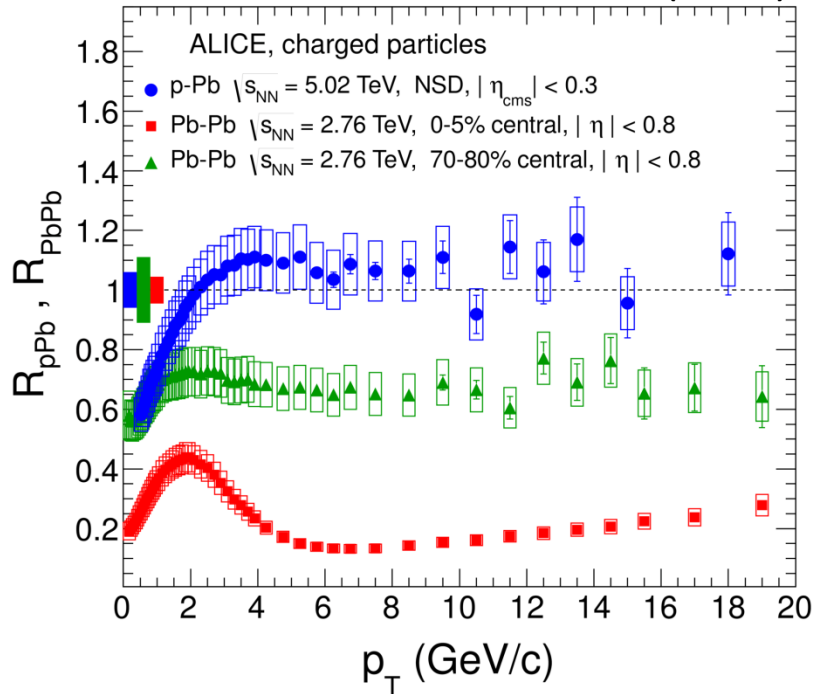


- The idea to use jets to probe the QGP goes back to Bjorken. He made a first study of collisional energy loss but never published this as he realized that radiative energy loss was much more important
- The phenomenology is quite difficult and has many variants so I focus on the experimental results



# High $p_T$ particles: a proxy for jets

ALICE, PRL 110, 082302 (2013)



At RHIC very few studies have been done with jets. Instead high  $p_T$  particles have been used to study jet quenching.

At LHC high  $p_T$  results are also used but additionally there are many jet studies.

Results at LHC are qualitatively similar to results from RHIC with a large suppression in central collisions

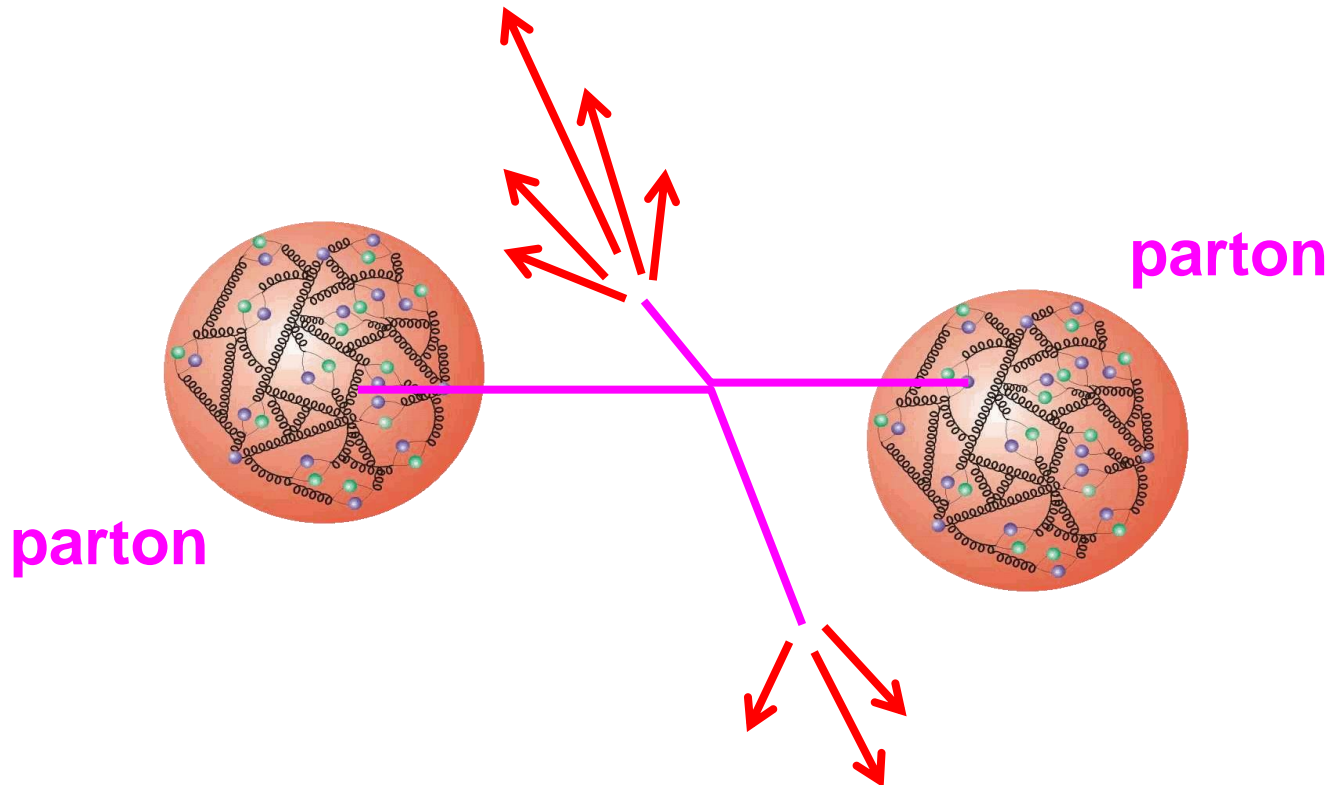
New things were the clear rise with  $p_T$

LHC has huge advantage due to the larger jet Xsection.

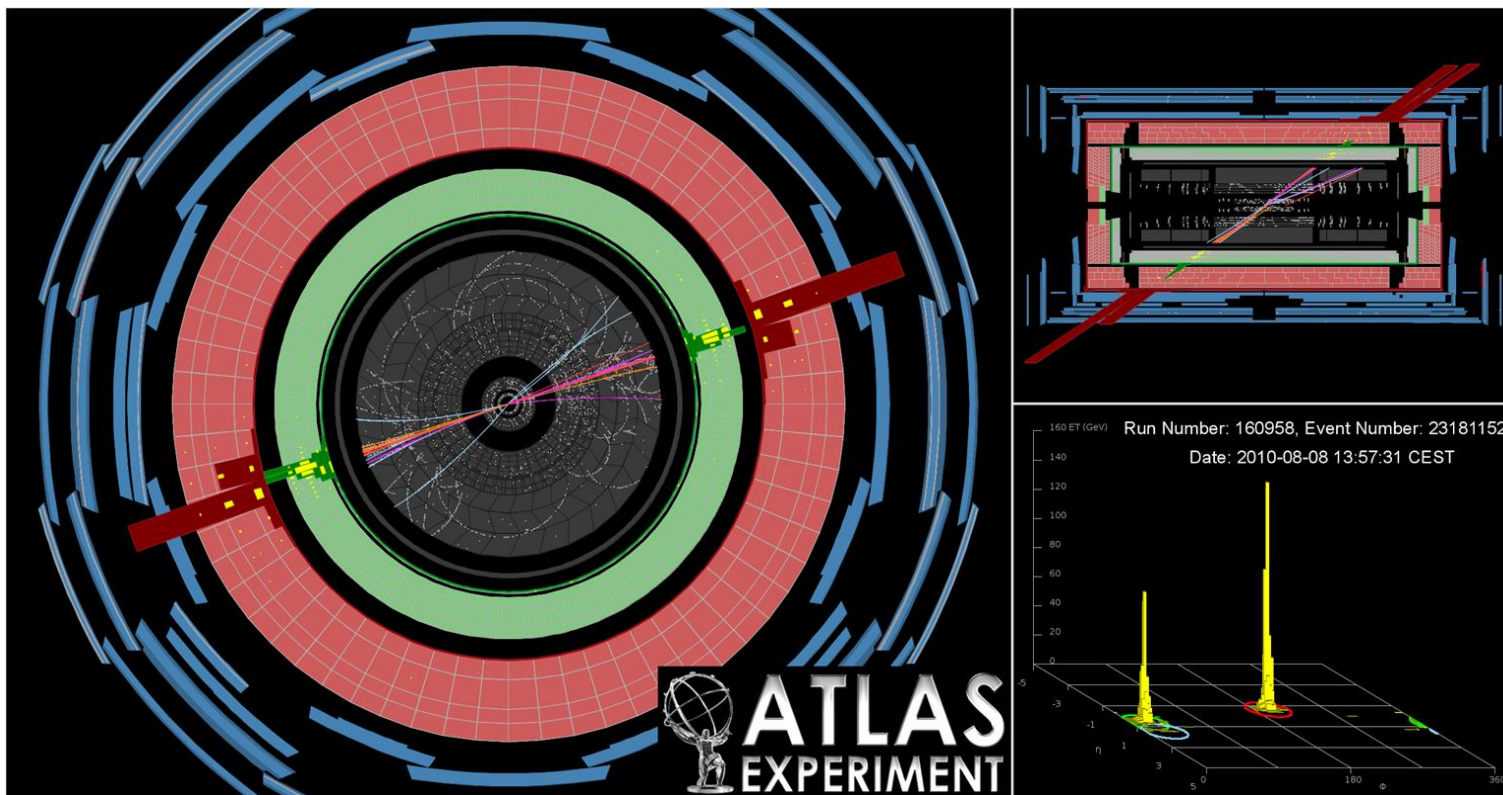


# Jets in pp

hadrons from jet fragmentation



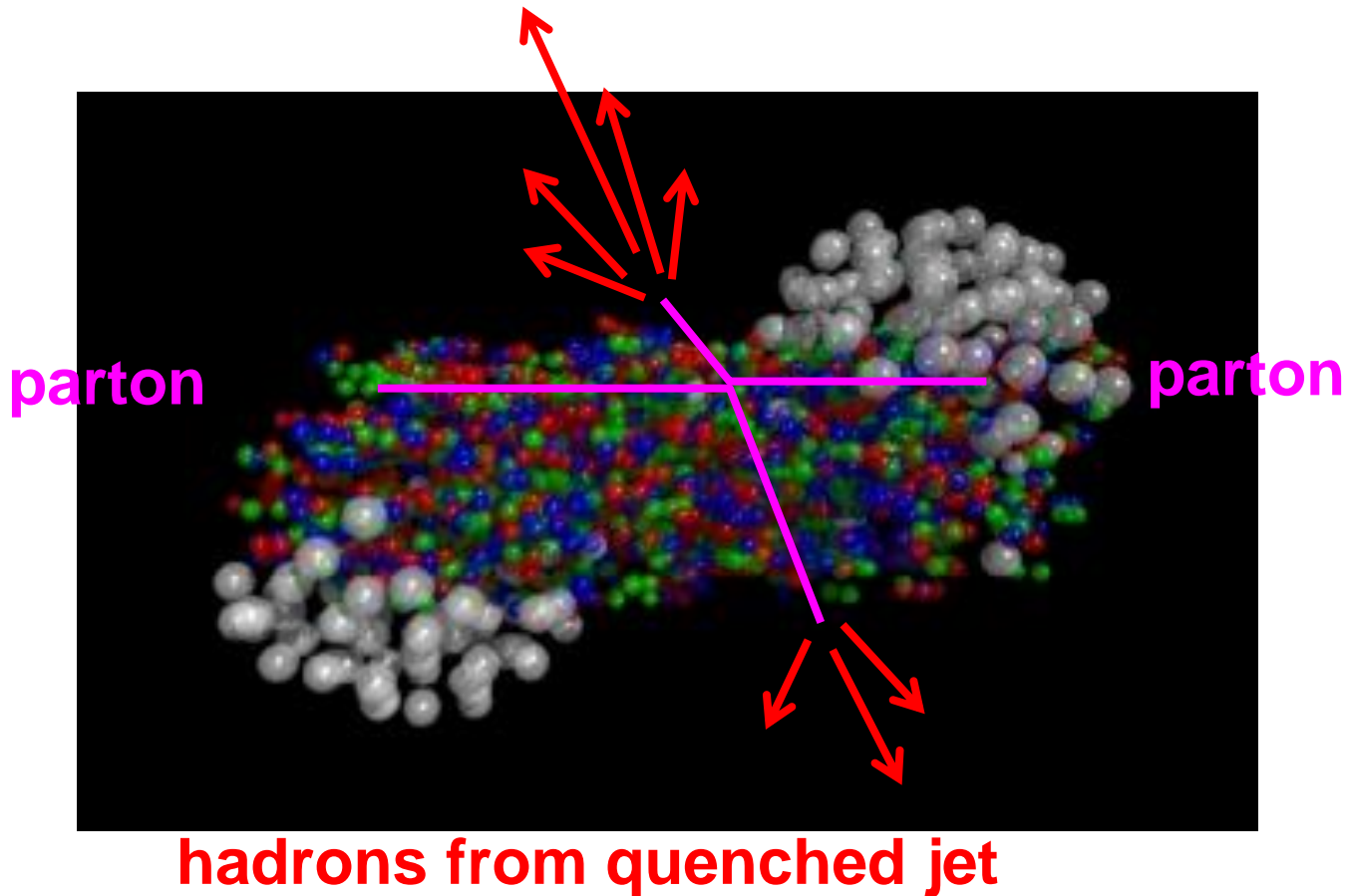
# Jets in pp





# Jets in Pb-Pb

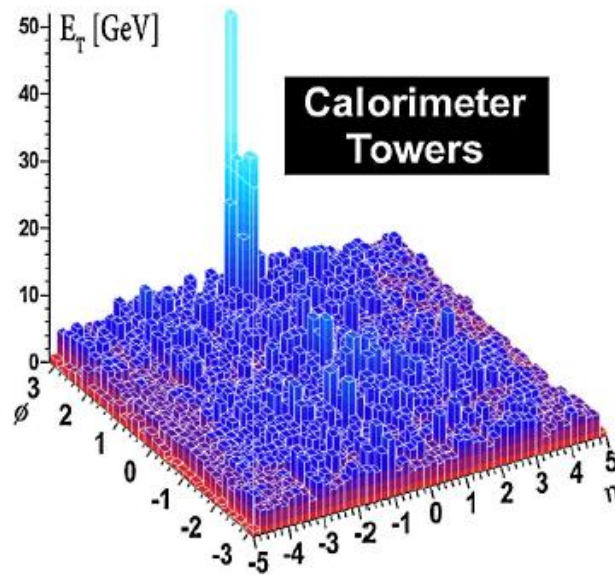
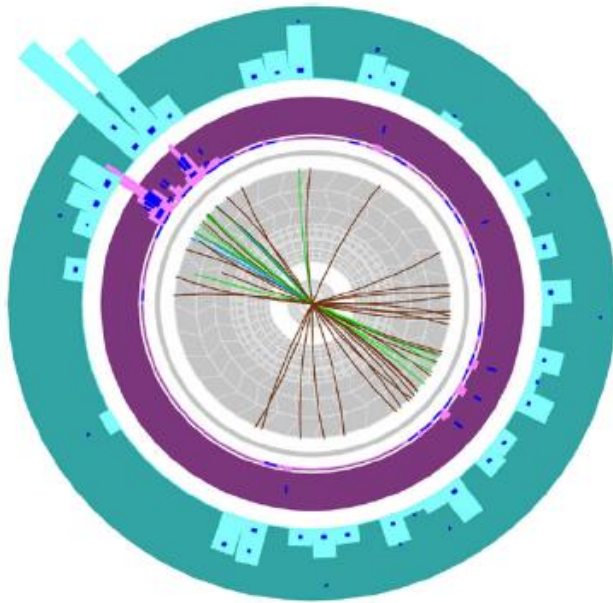
hadrons from leading jet





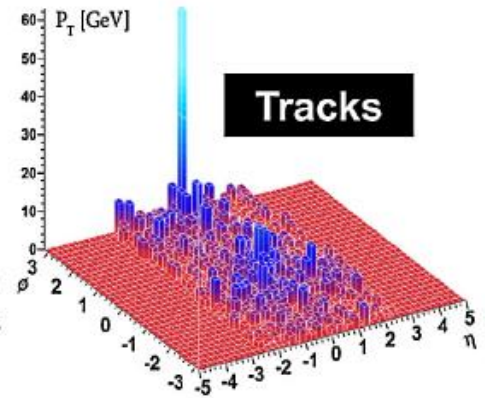
# Jets in Pb-Pb (ATLAS)

Phys. Rev. Lett. 105, 252303



## ATLAS

Run: 169045  
Event: 1914004  
Date: 2010-11-12  
Time: 04:11:44 CET

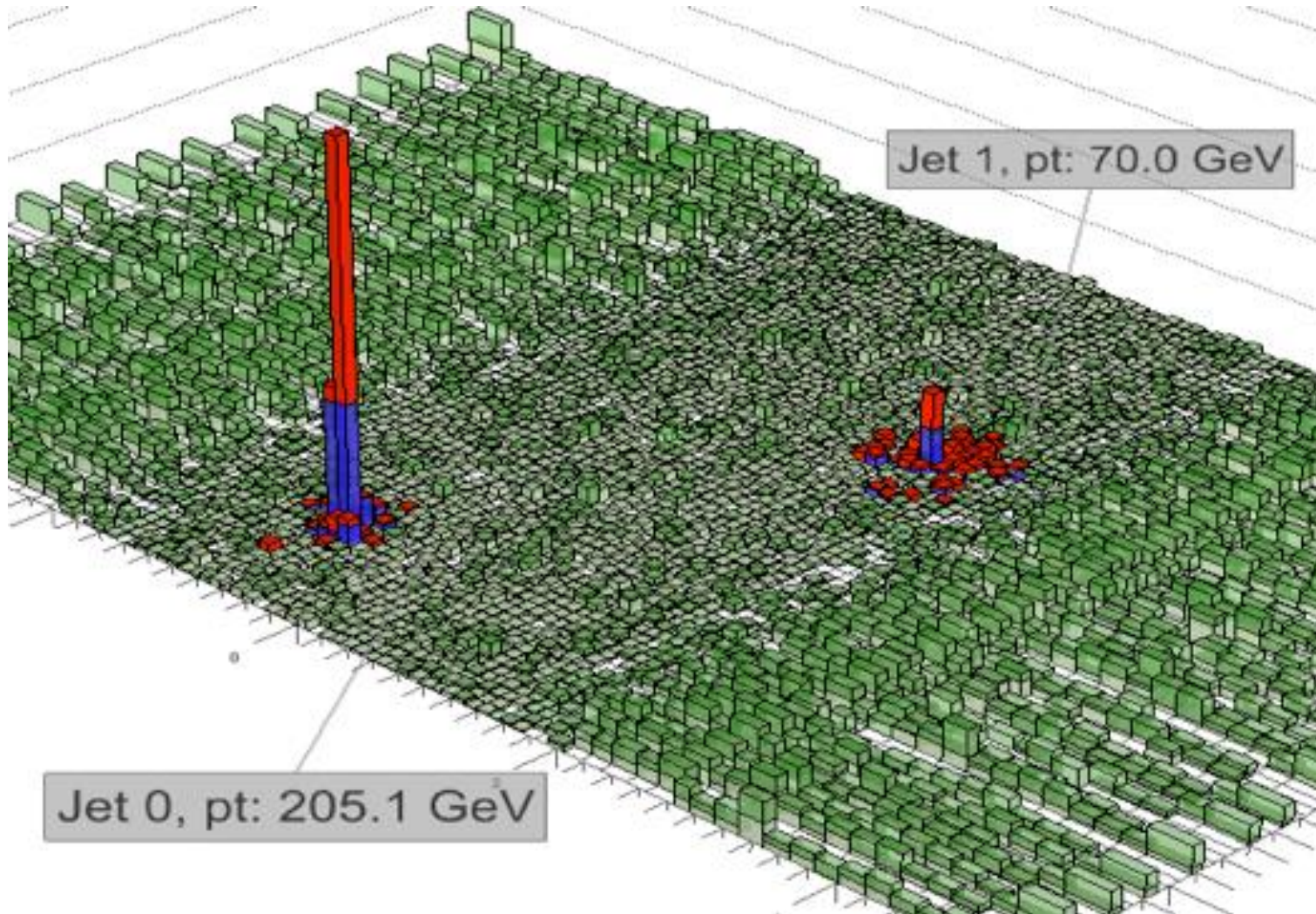


- Jet asymmetry – away side jet is absorbed/modified by the medium  
Advantage of jets is that they “map” onto the QCD degrees of freedoms: quarks and gluons (more complicated in heavy ion collisions)  
Advantage of ATLAS and CMS is larger acceptance



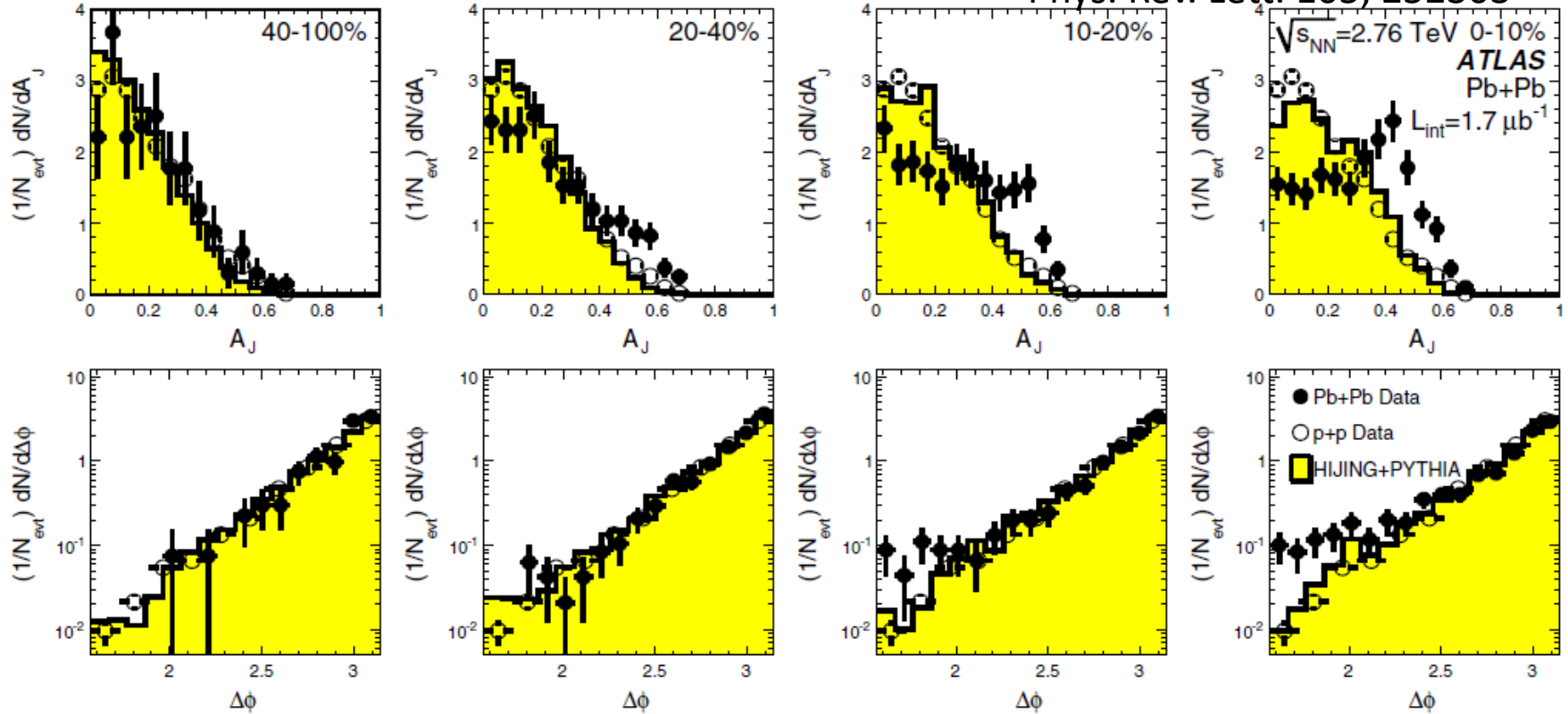


# Jets in Pb-Pb (CMS)



# Quantifying the dijet asymmetry

Phys. Rev. Lett. 105, 252303



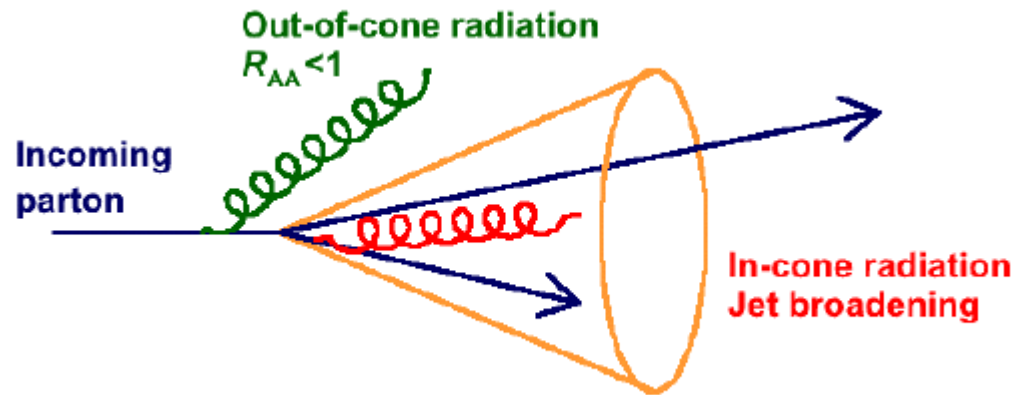
$$A_J = \frac{E_{T1} - E_{T2}}{E_{T1} + E_{T2}}$$

Where  $E_{T1}$  ( $E_{T2}$ ) is the transverse energy of the leading (subleading) jet ( $E_{T1} > 100$  GeV and  $E_{T2} > 25$  GeV).

Notice that the jets are still back-to-back!



# Looking into the jet structure



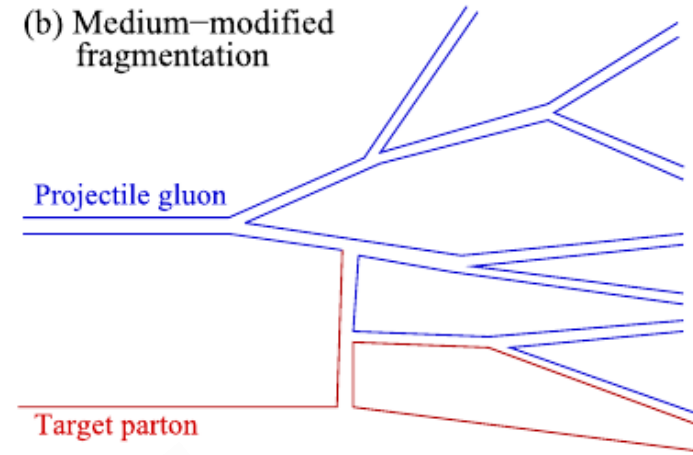
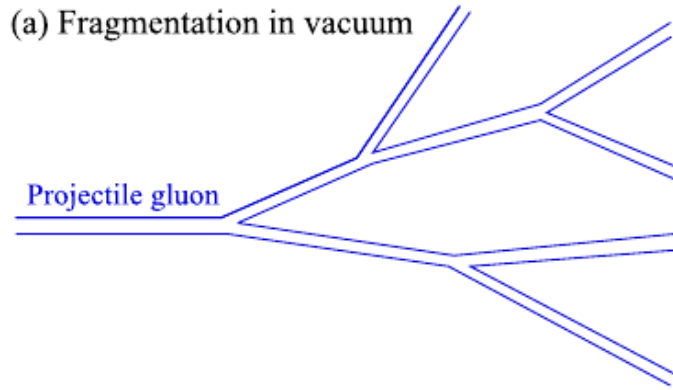
The motivation for these studies was:

- to recover some of the radiated energy (in principle jets did not have to be suppressed!)
- to study how the FF is modified

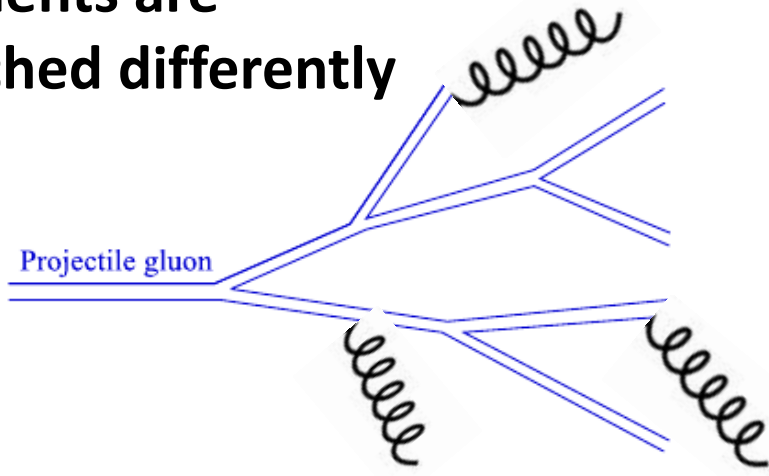


# Two examples of ideas for modified FFs

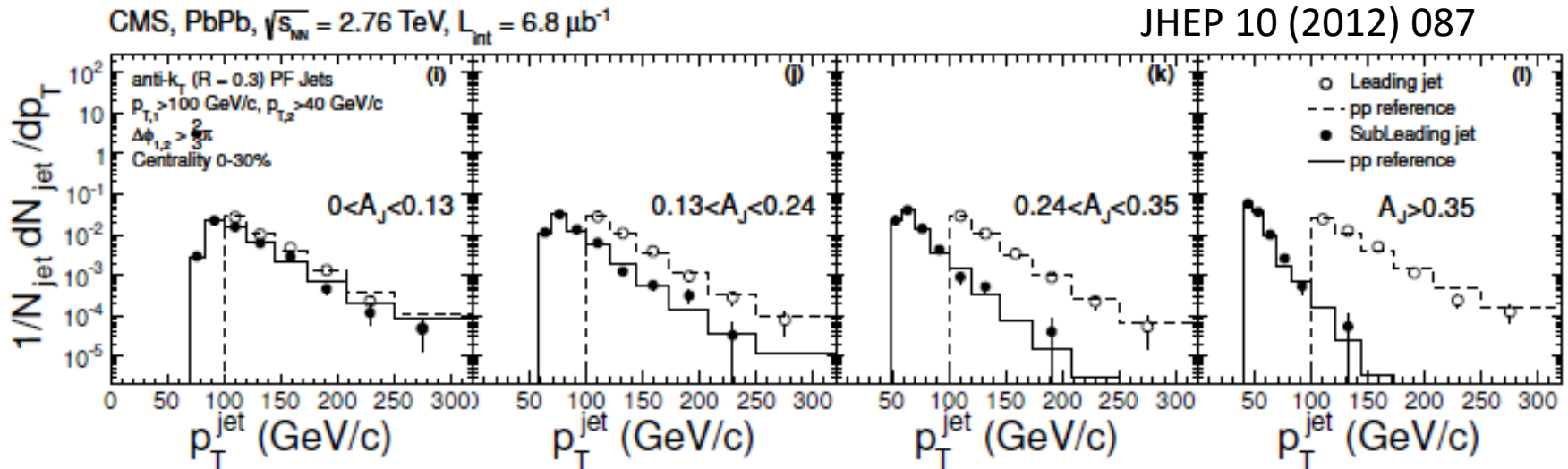
## Direct interaction with medium



**Fragments are quenched differently**



# Selecting dijet events and comparing to pp



The  $A_J$  selection introduces the same bias on the dijet samples

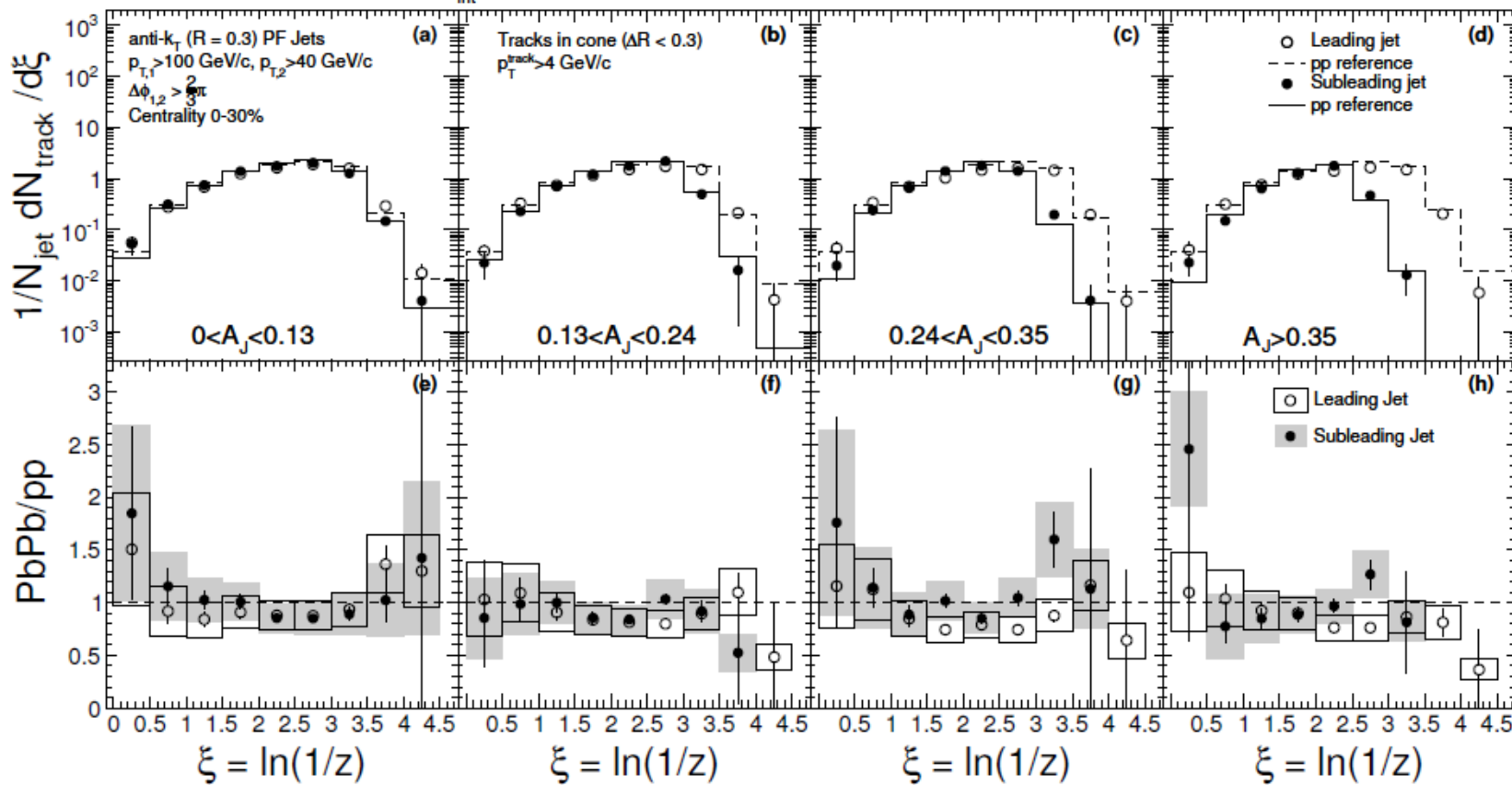




# Surprisingly the jet structure is the same! (tracks with $p_T > 4 \text{ GeV}/c$ )

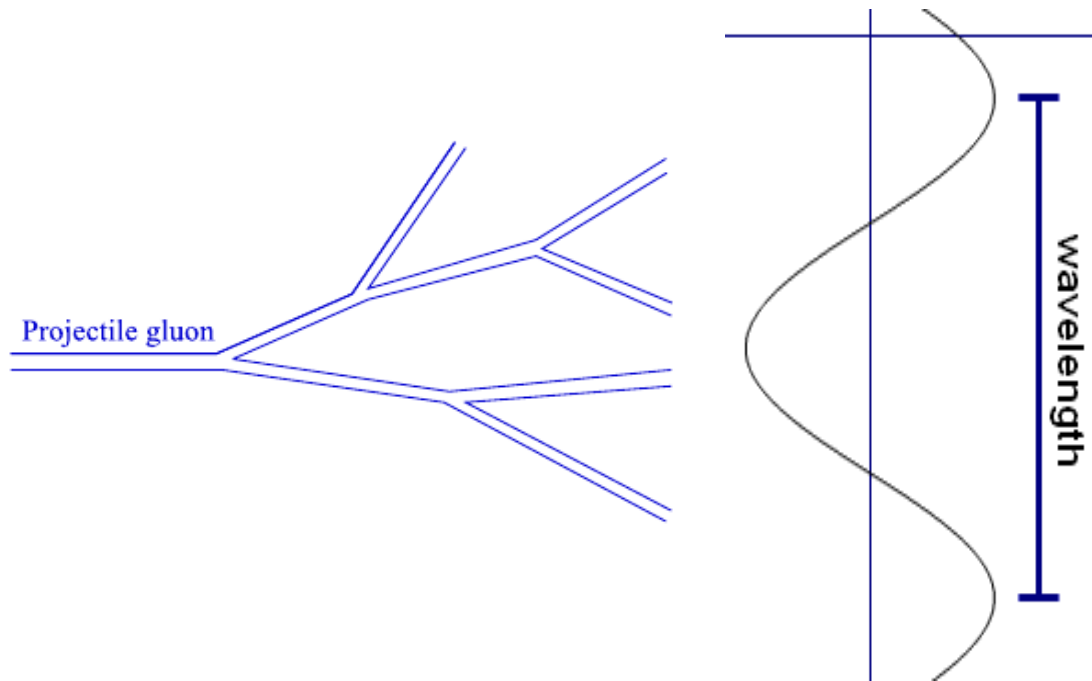
JHEP 10 (2012) 087

CMS, PbPb,  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ ,  $L_{int} = 6.8 \mu\text{b}^{-1}$



The result shows that quenched jets looks like pp (vacuum) jets!  
 Even in the case where  $A_J$  is large and for the subleading jet!

# The results show that quenching is coherent!?

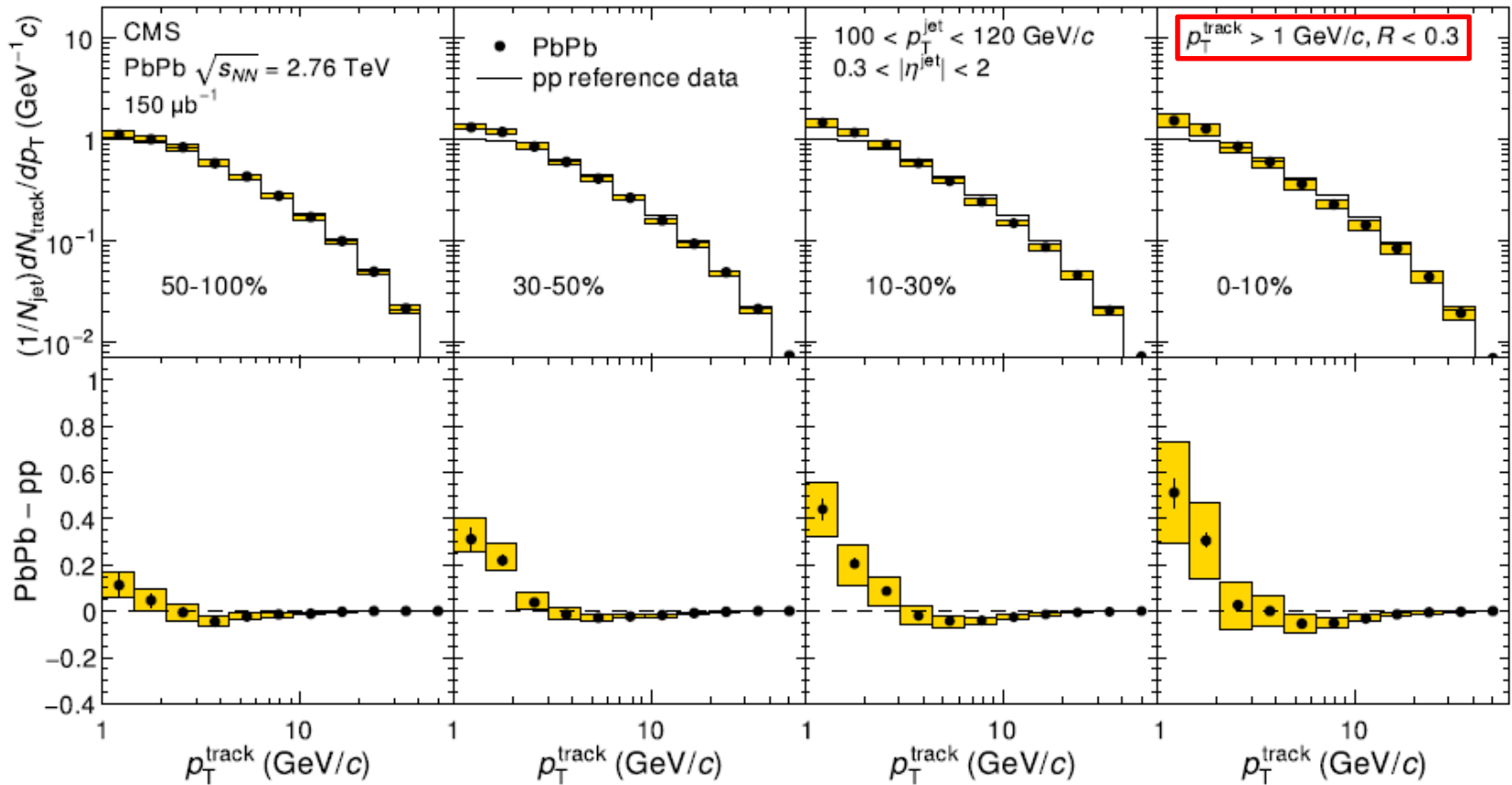


One proposed solution is that the medium cannot resolve the jet constituents (Phys.Lett. B725 (2013) 357-360).

So for the medium the jet will look essentially as a single parton and so all fragments are quenched coherently!

# The finer structure of the fragmentation function

Phys. Rev. C 90, 024908



This analysis considers tracks down to  $p_T = 1$  GeV/c

The modifications are sitting at low  $p_T$ , mainly  $p_T < 3$  GeV/c

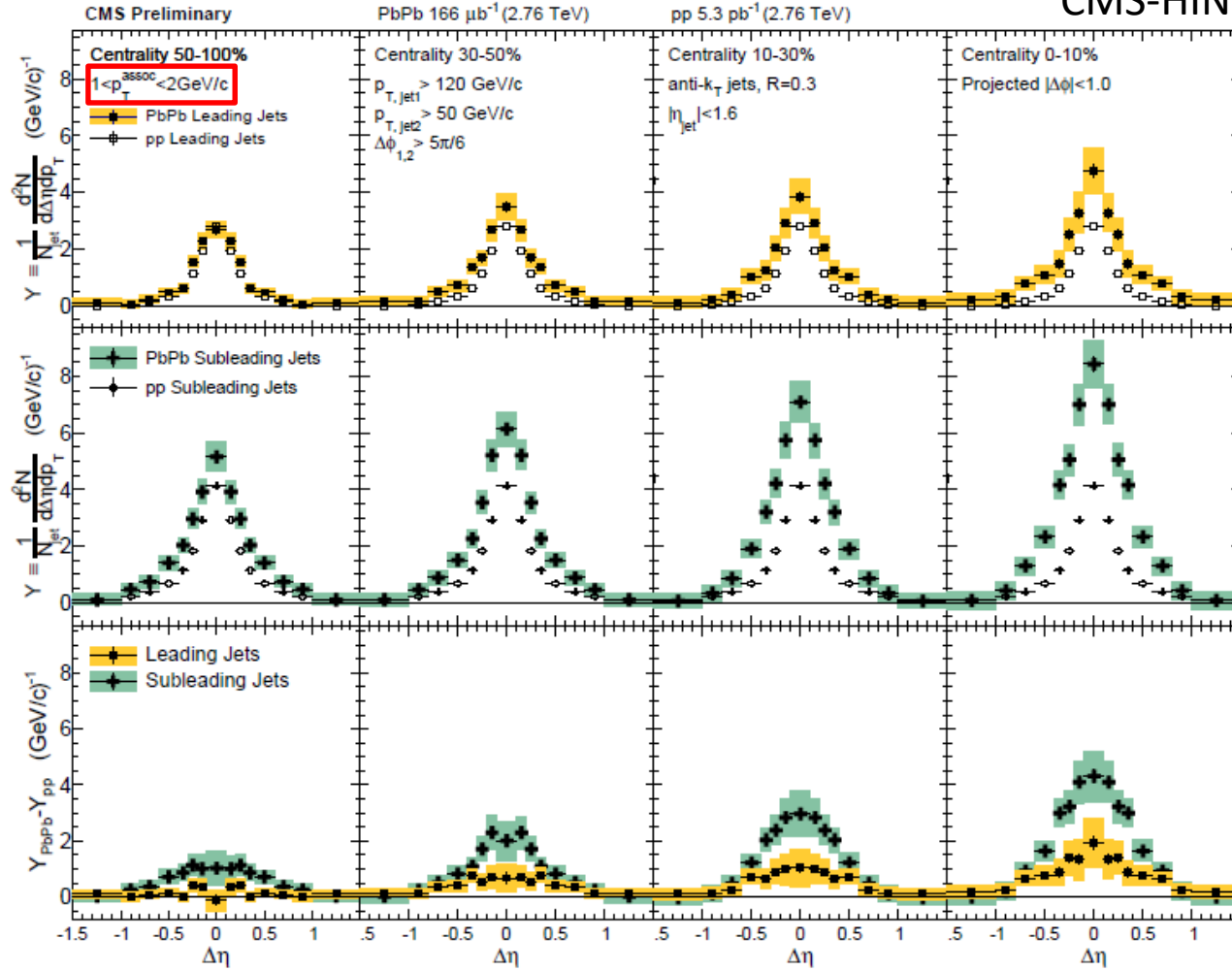






# Tracking the energy loss

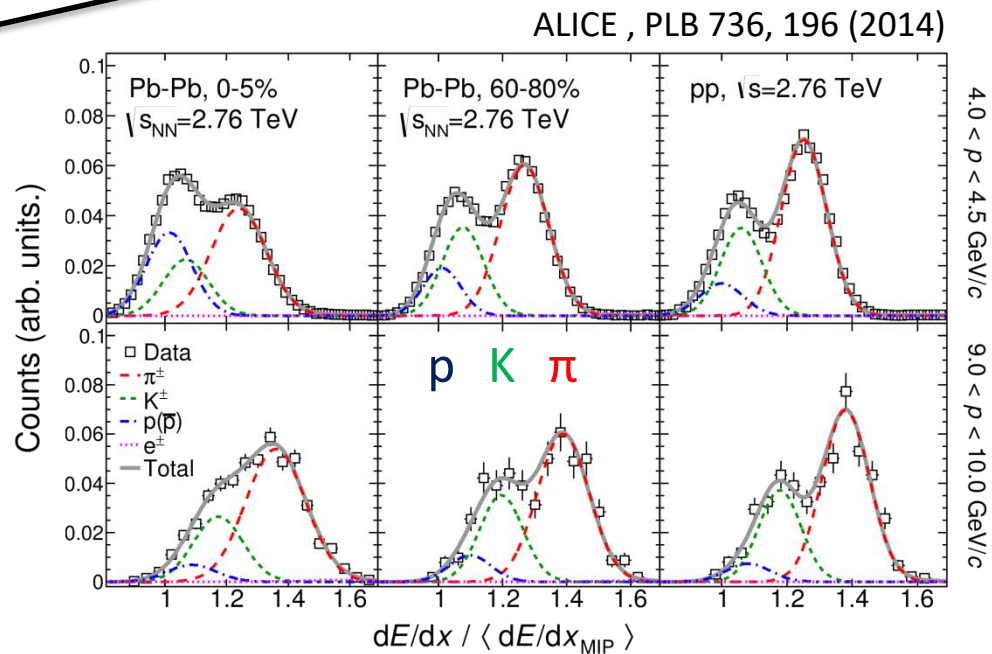
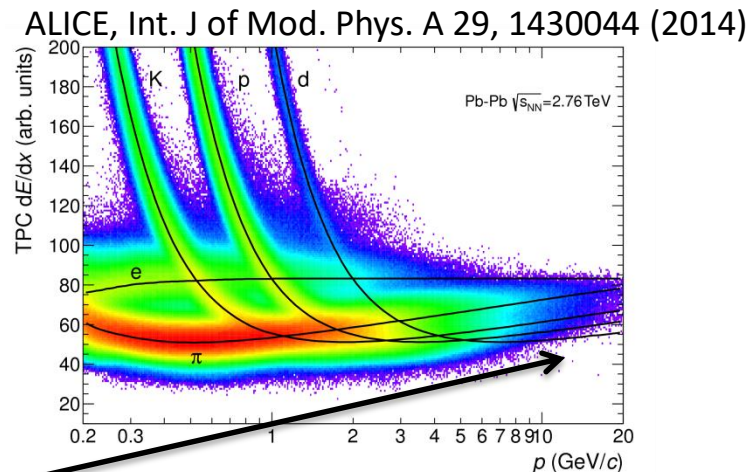
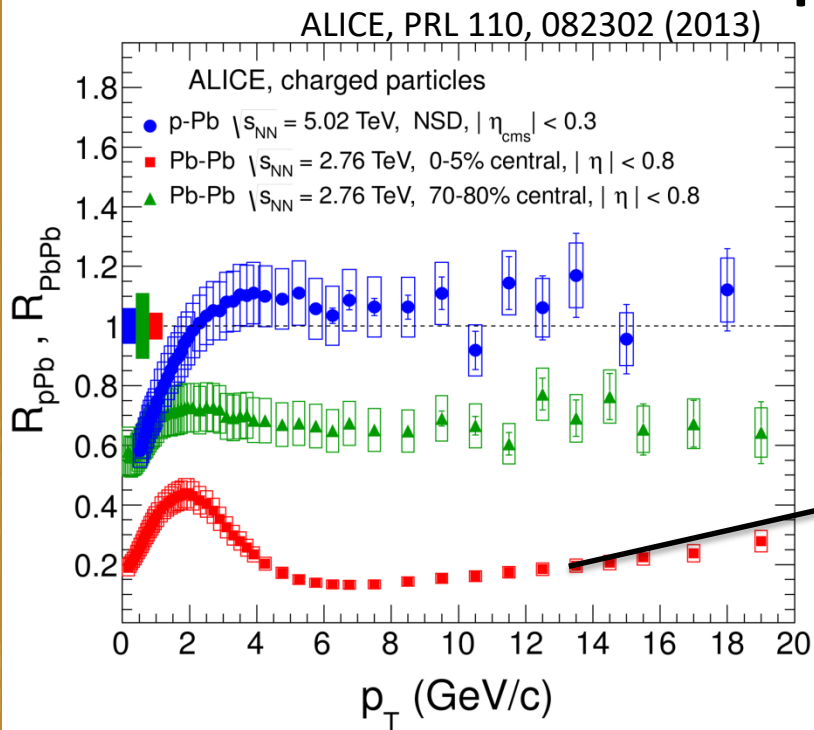
CMS-HIN-14-016



CMS have shown that the energy loss can be recovered at low  $p_T$  at large  $\Delta\eta$  and large  $\Delta\phi$



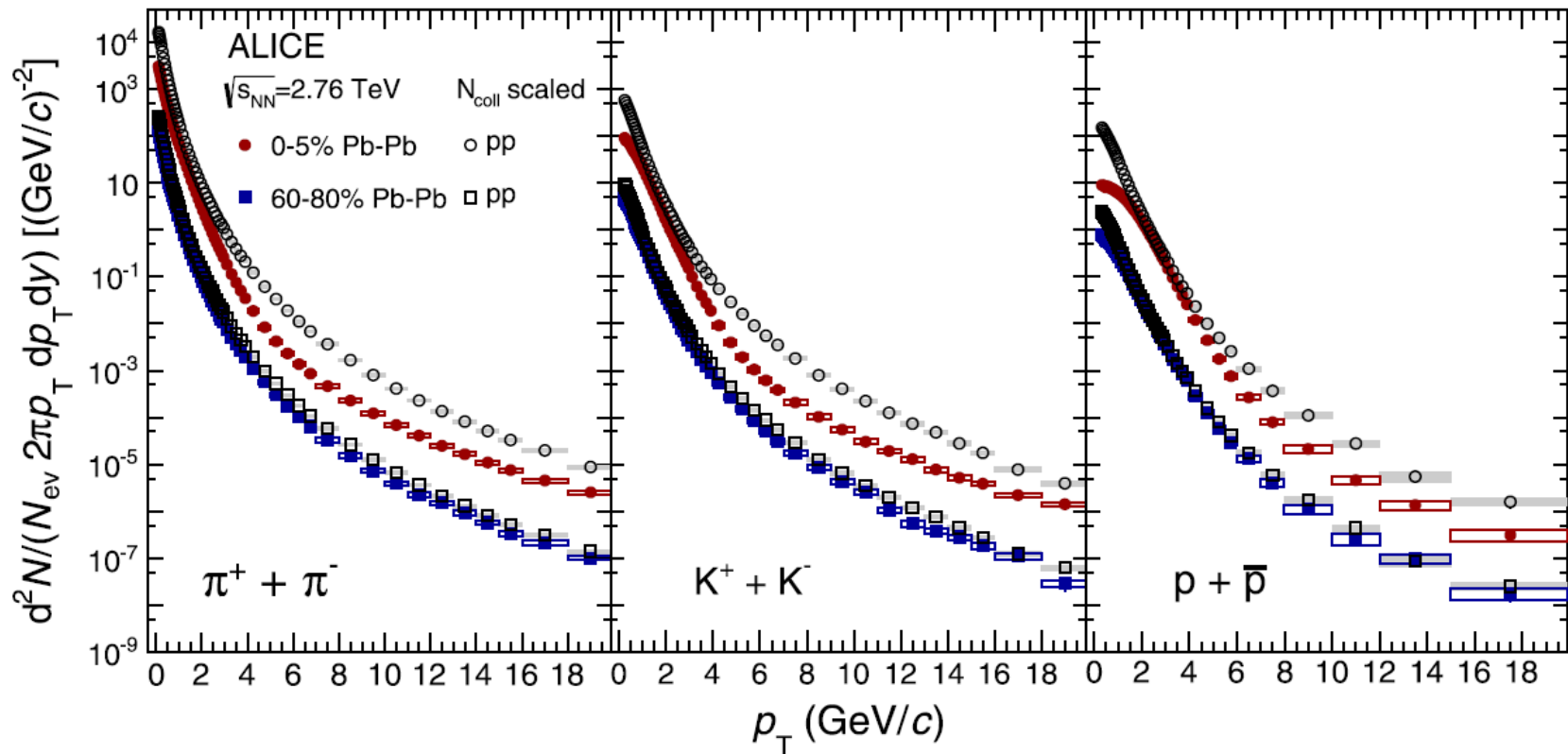
# Extending the $R_{AA}$ to identified particles



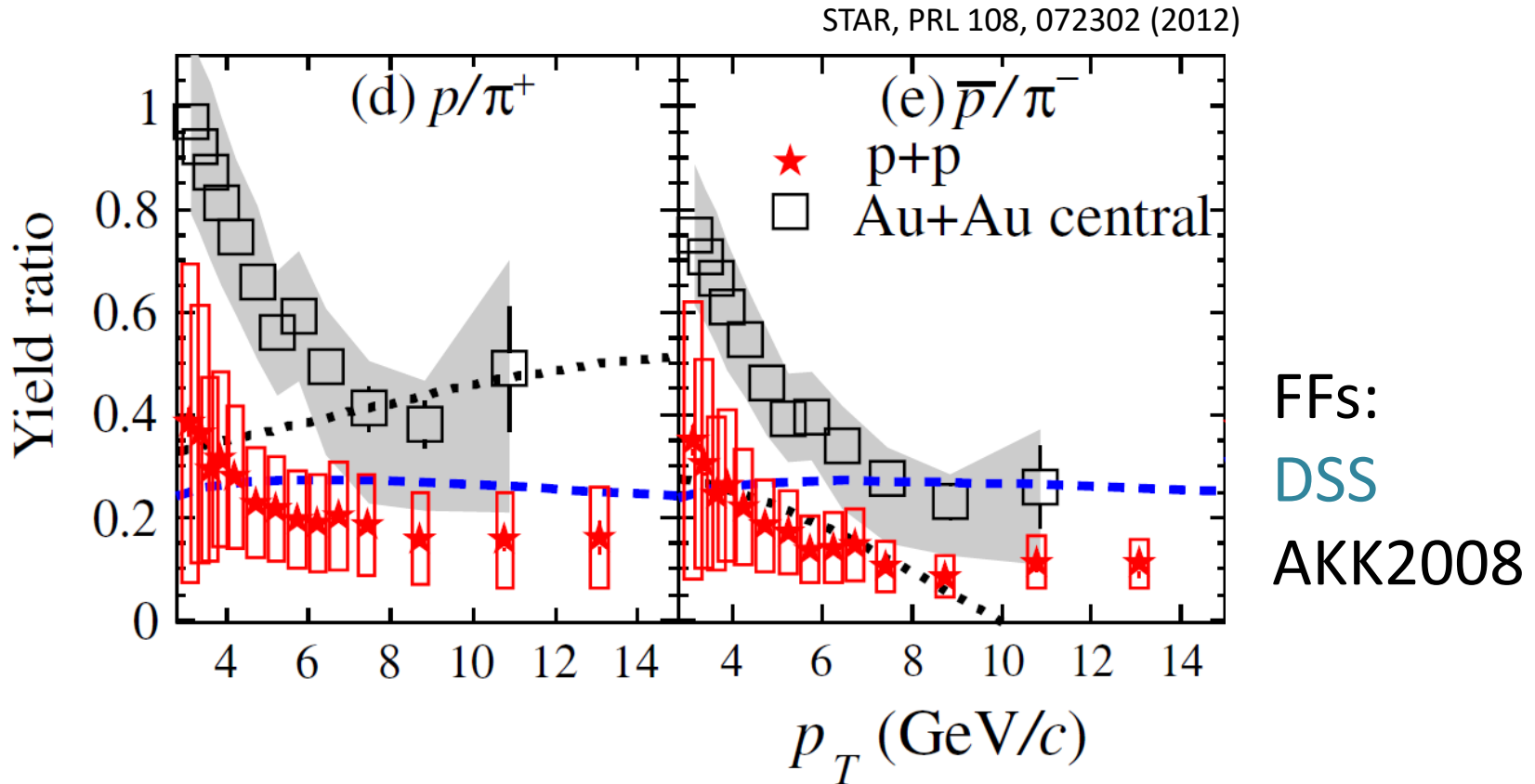
Each TPC track also has an associated  $dE/dx$  that can be used for ID on the relativistic rise

# Charged $\pi$ , $K$ , and $p$ spectra in pp and Pb-Pb collisions

ALICE, PLB 736, 196 (2014)



# High $p_T$ $p/\pi$ ratio in pp and central Au-Au collisions at $\sqrt{s_{NN}}=200$ GeV

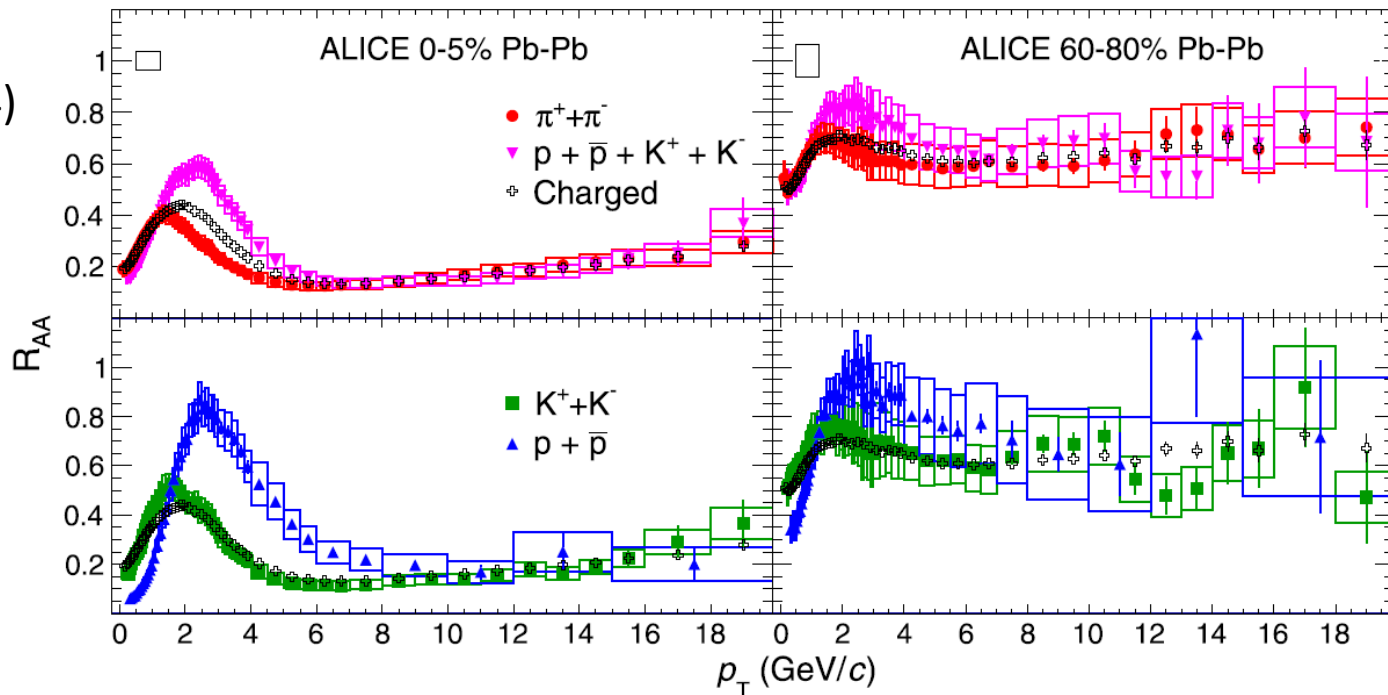


NB! At this low energy there is significant difference between baryon and anti-baryon production.



# The nuclear modification factor $R_{AA}$ for charged $\pi$ , $K$ , and $p$

PLB 736,  
196 (2014)



The nuclear modification factors are consistent for  $\pi$ ,  $K$ ,  $p$  for  $p_T > 10$  GeV/ $c$ .  
For 0-5% Pb-Pb collisions:  $K$  and  $\pi$  to within  $\sim 10\%$  (syst),  $p$  and  $\pi$  to within  $\sim 20\%$  (syst)

The results rule out popular ideas of jet quenching where the large energy loss is coupled to large leading order particle species dependent effects.





# Conclusions on jets

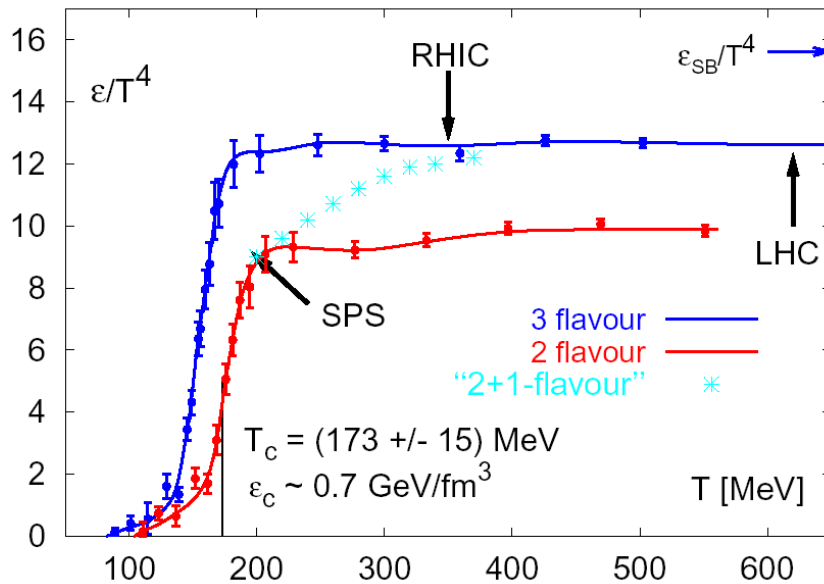
- High  $p_T$  particles (and jets) were one of the hottest topics at RHIC where they were difficult to measure
  - And it seemed at RHIC that jets were modified in a spectacular way giving rise to exotic effects
- At LHC the picture we have of jet quenching has turned out to be surprisingly simple
  - Experimentally the jets seem to first lose energy in the QGP and then afterwards to fragment as vacuum jets

The energy loss is observed as soft particles at large angles



# SOFT MEDIUM PROPERTIES

# Lattice QCD calculation of the energy density



$$\epsilon_{\text{Quark-Gluon gas}} = \frac{\pi^2}{30} \left( 2 \times 8 + \frac{7}{8} 2 \times 2 \times 3 \times 3 \right) T^4$$

↑  
**Gluon spin and color**

↑  
**(Anti+)quark spin, color and flavor**

Because of the similarity with Stefan-Boltzmann energy density for a quark-gluon gas:

QGP should be weakly coupled

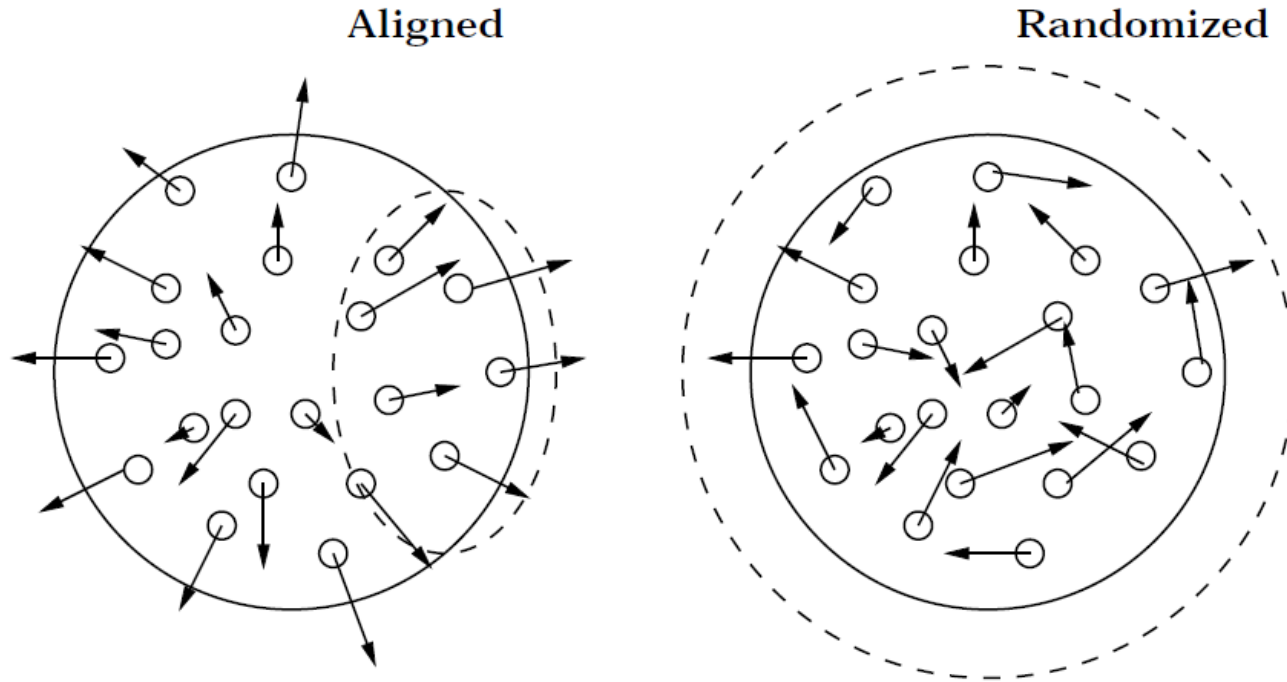
At a deeper level this is also what we expect from asymptotic freedom







# Radial flow

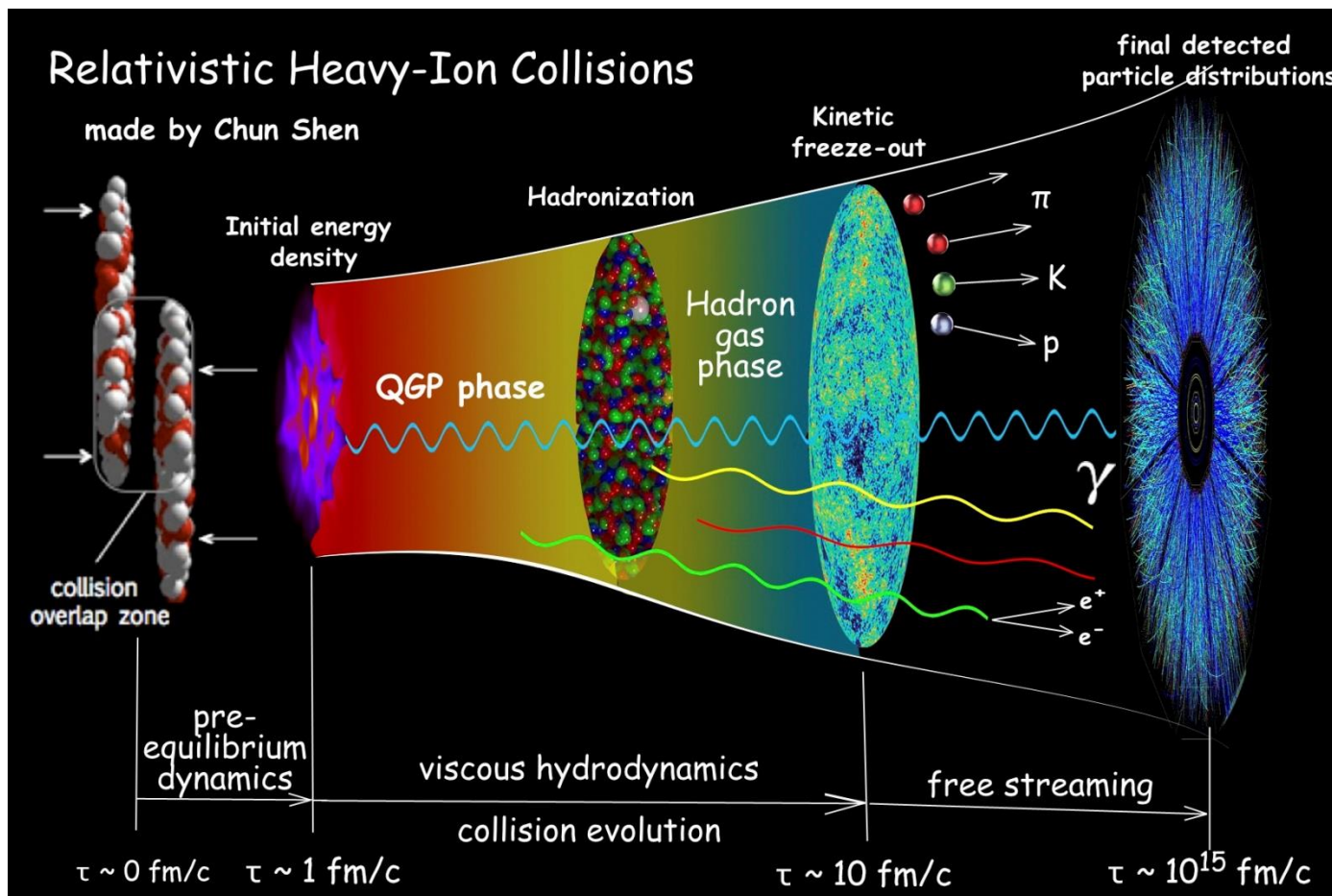


Flow in general plays a very important role in heavy ion collisions. Flow means that there are strong correlation between position and momentum in the source.

We believe that flow in the partonic phase is imprinted on the final state hadrons at freeze out.

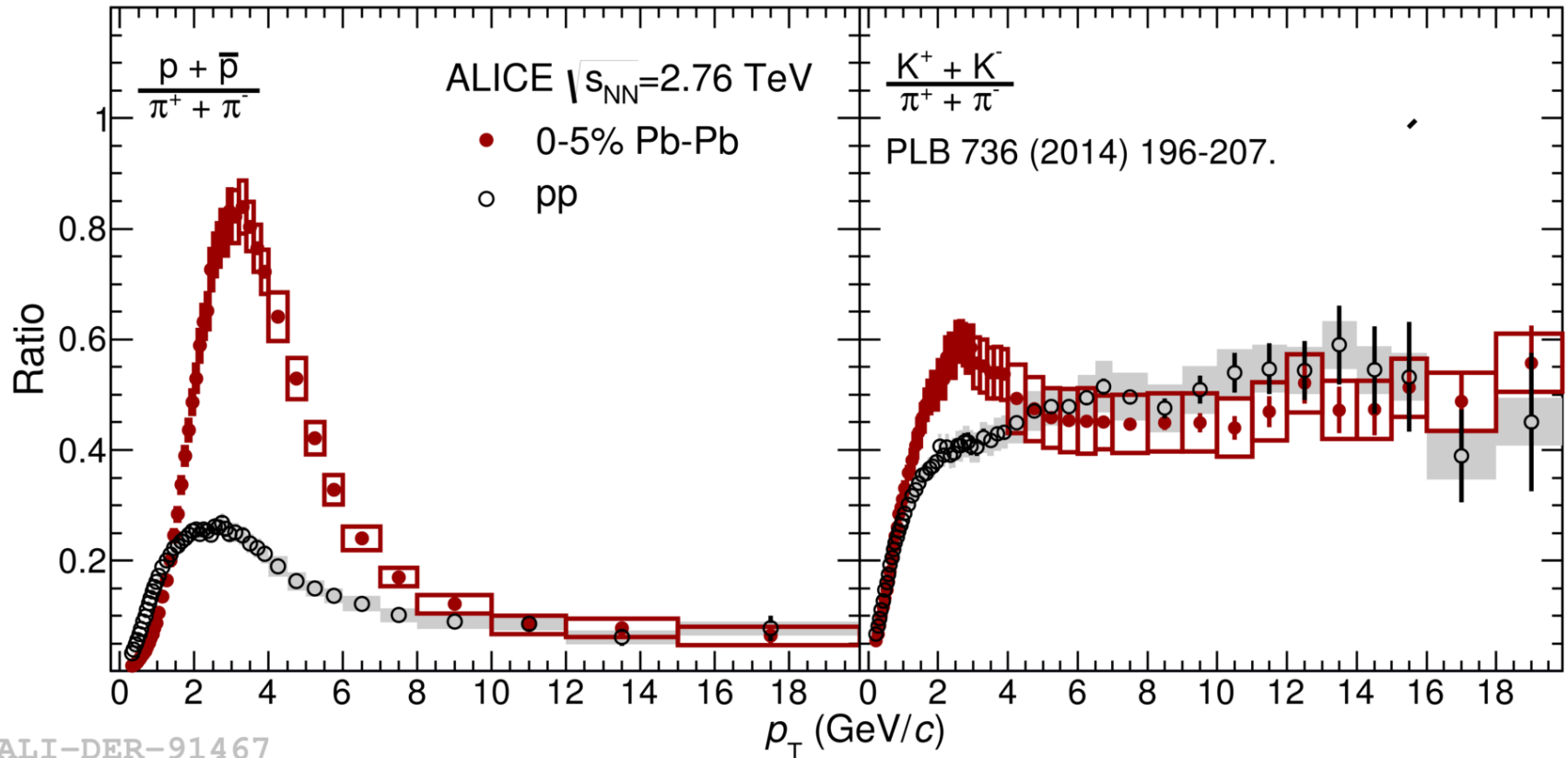


# Evolution of the heavy-ion collision



Mainly this period where flow build up

# Radial flow pushes heavier particles out to larger $p_T$

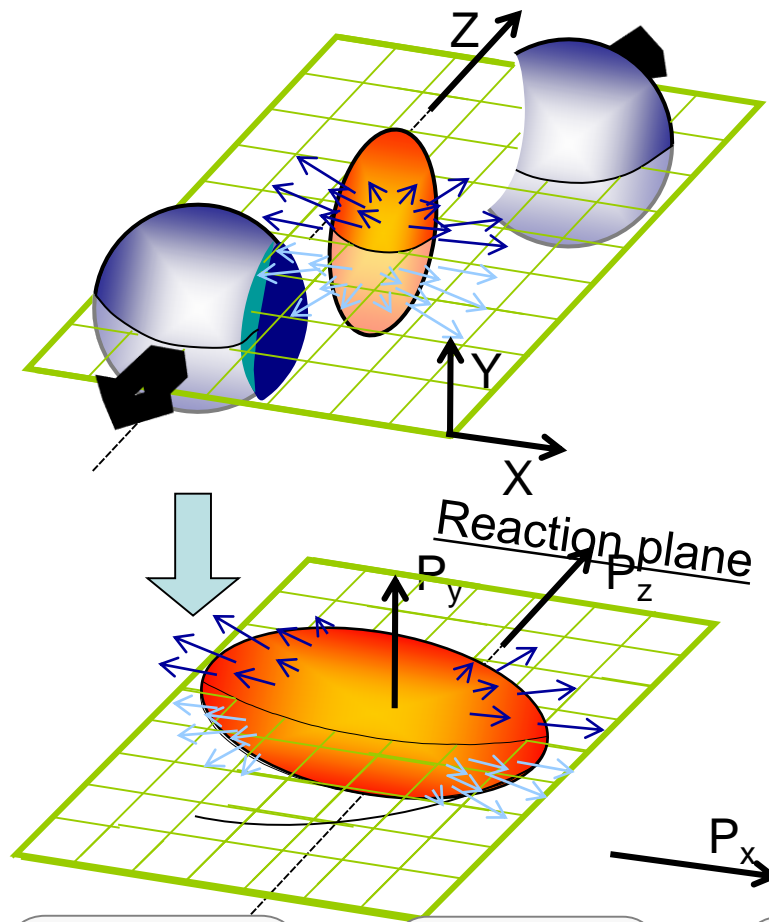
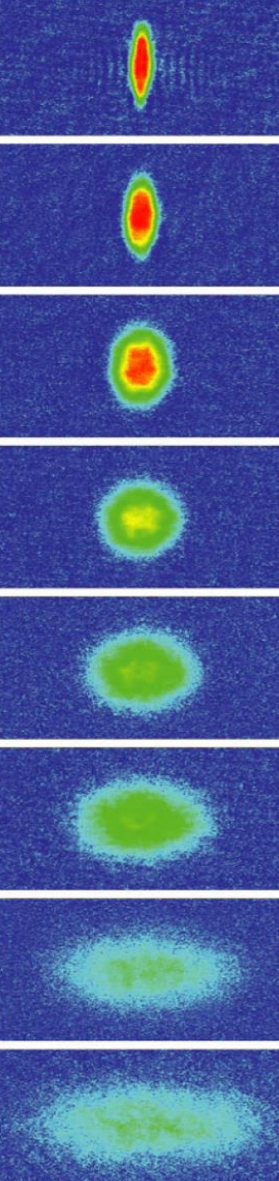


If the source has a radial flow velocity  $\beta_r$  then particles at rest in the flowing frame are boosted in the lab as:  $p_T \sim \gamma \beta_r m$   
 So they get a mass dependent boost!

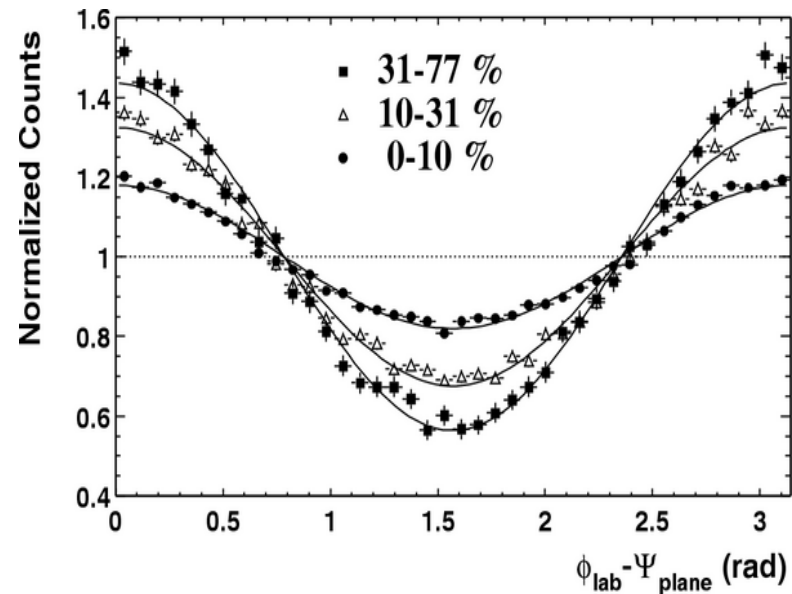




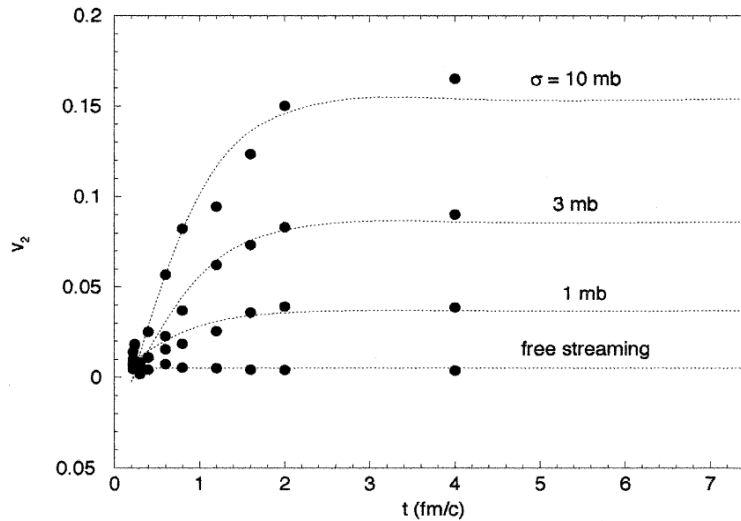
# Elliptic flow ( $v_2$ )



Fourier decomposition:  
 $dN/d\Delta\phi = 1 + 2 v_2 \cos(2 \Delta\phi)$



# Elliptic flow requires early strong interactions to form

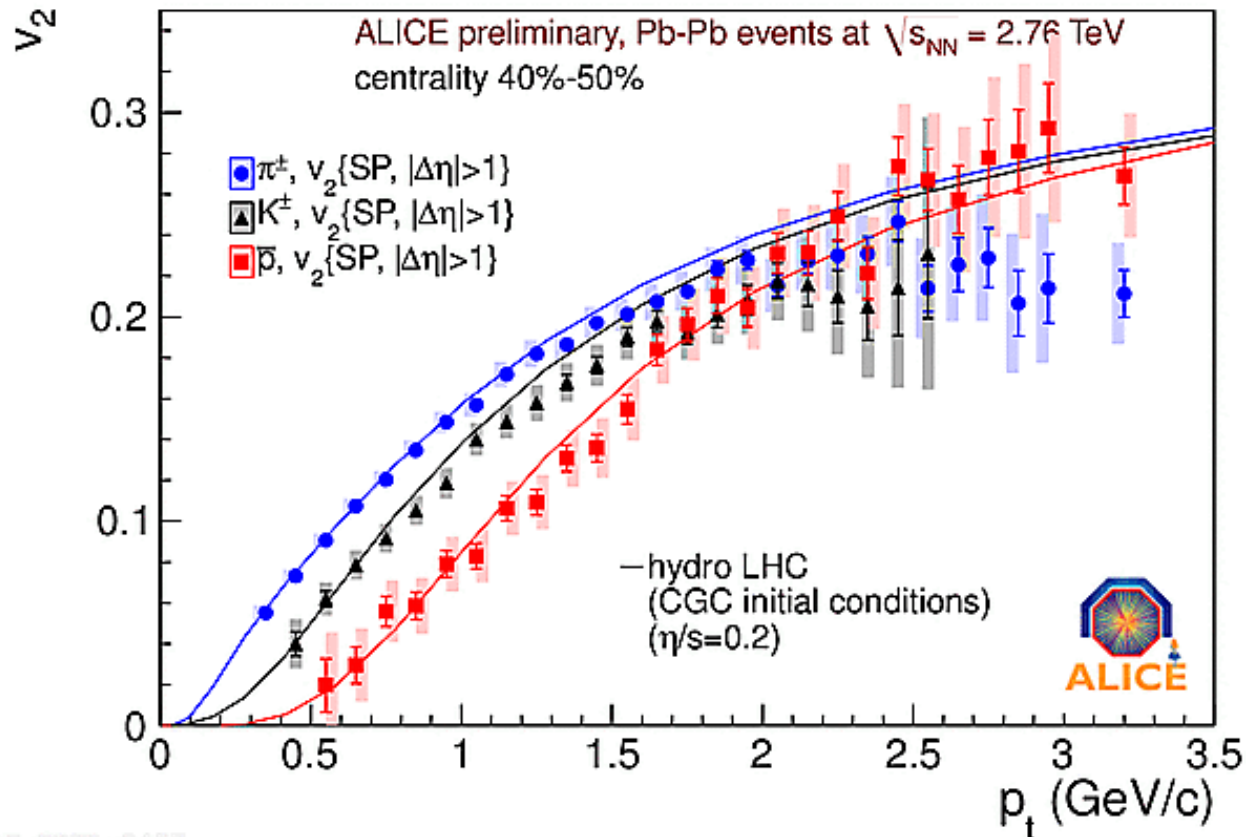


Zhang, Gyulassy, Ko,  
Phys. Lett. B455 (1999) 45

- Each nucleon-nucleon interaction produces on average a spherical symmetric distribution.  
Only by interacting elliptic flow is generated



# Mass ordering of the elliptic flow



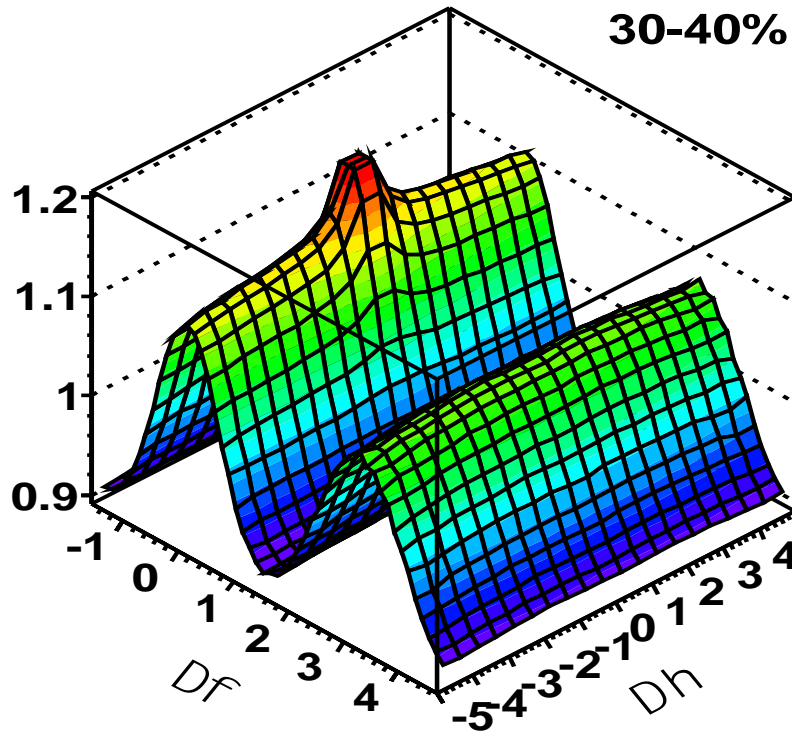
ALICE-PREL-2457

The mass ordering is characteristic of flow since heavier particles are pushed out to higher  $p_T$   
Surprisingly this is well described by nearly ideal hydrodynamics





# Ridge (ATLAS)



One can study flow using two particle correlations. Here  $2 < p_{T1}, p_{T2} < 3$  GeV/c

The long range correlations are imprinted by the flow because **the initial elliptic overlap is a global property of a collision!**

For completeness:

$$\text{Singles: } \frac{dN}{df} \propto 1 + \sum_n \hat{a}_n 2v_n \cos n(f - Y_n) \quad \text{EP method}$$

$$\text{Pairs: } \frac{dN}{dDf} \propto 1 + \sum_n 2v_n^a v_n^b \cos(nDf) \quad \text{2PC method}$$



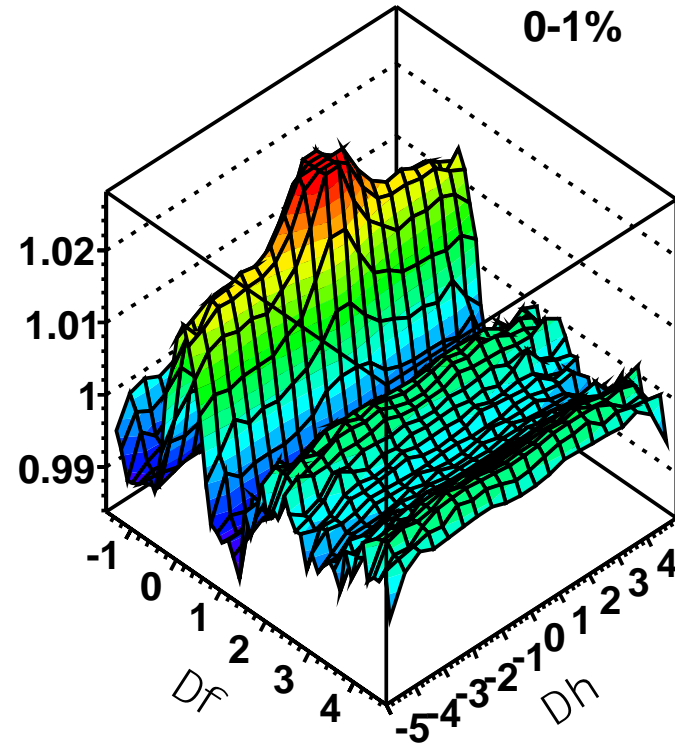
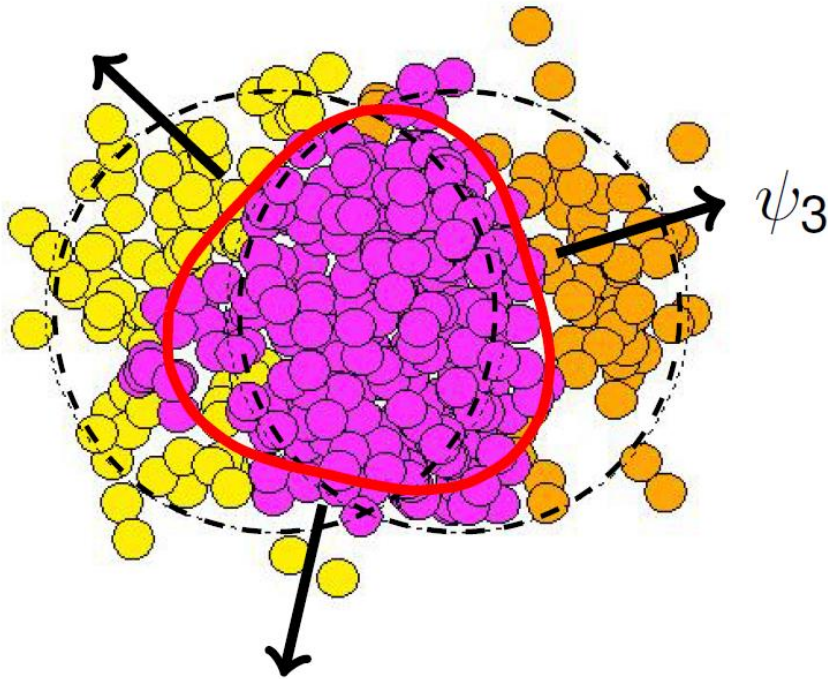
# Some simulation movies of hydrodynamics

- <https://www.youtube.com/watch?v=G18pyV0mSRw>





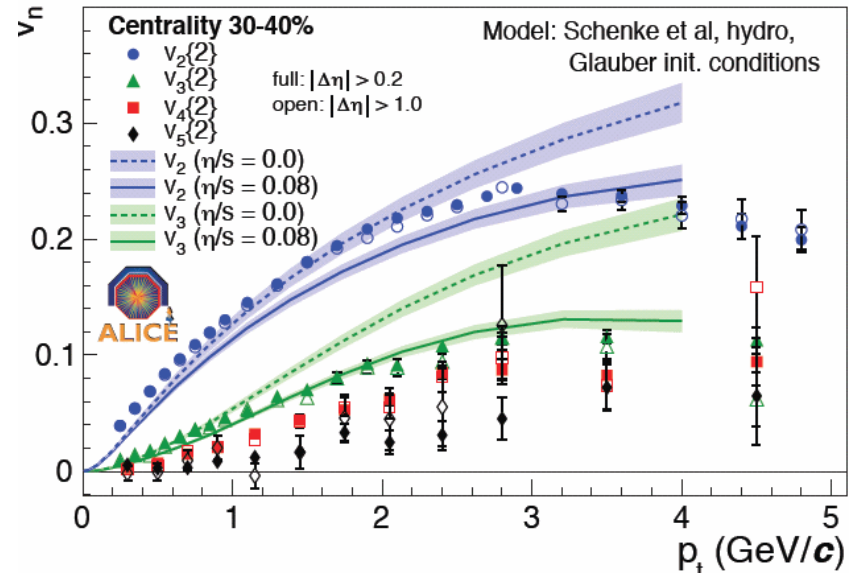
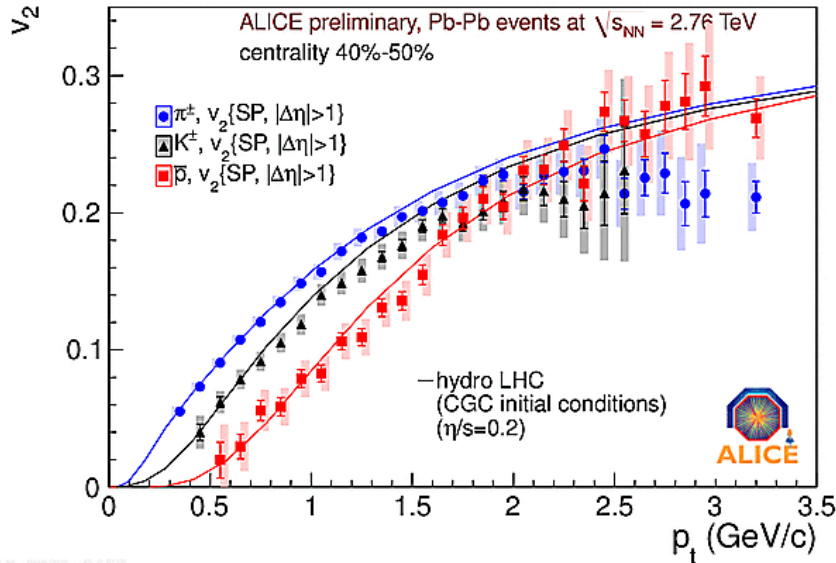
# Flow fluctuations



Because the nuclei are not homogenous the initial state can be quite asymmetric giving rise to e.g. triangular flow!

Famous paper: B. Alver, G. Roland, Phys.Rev. C81 (2010) 054905  
(420 citations)

# Elliptic flow and triangular flow is almost ideal!

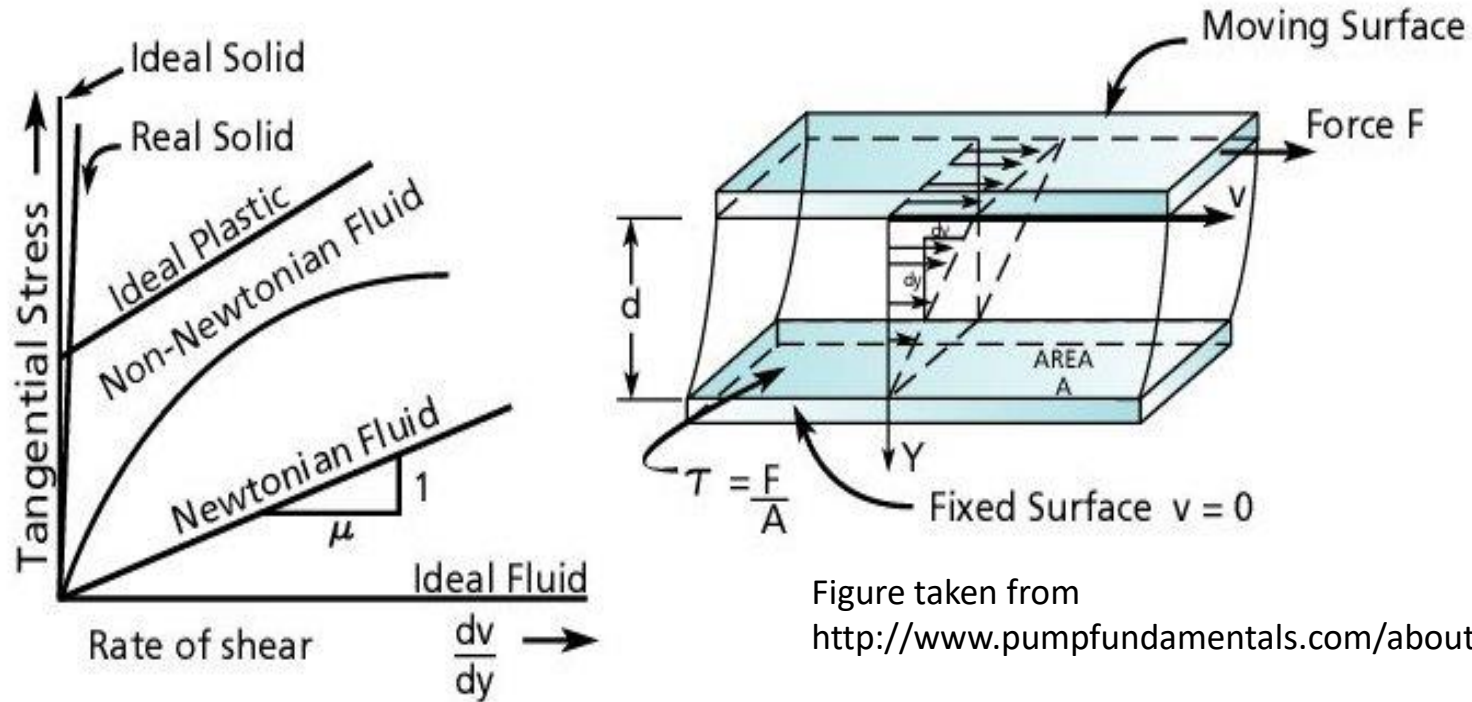


- Huge flow at intermediate  $p_T$ :  
2 times more particles in plane than out  
Nearly ideal fluid
- Significant higher order flow caused by fluctuations – also described by nearly ideal hydro + initial state





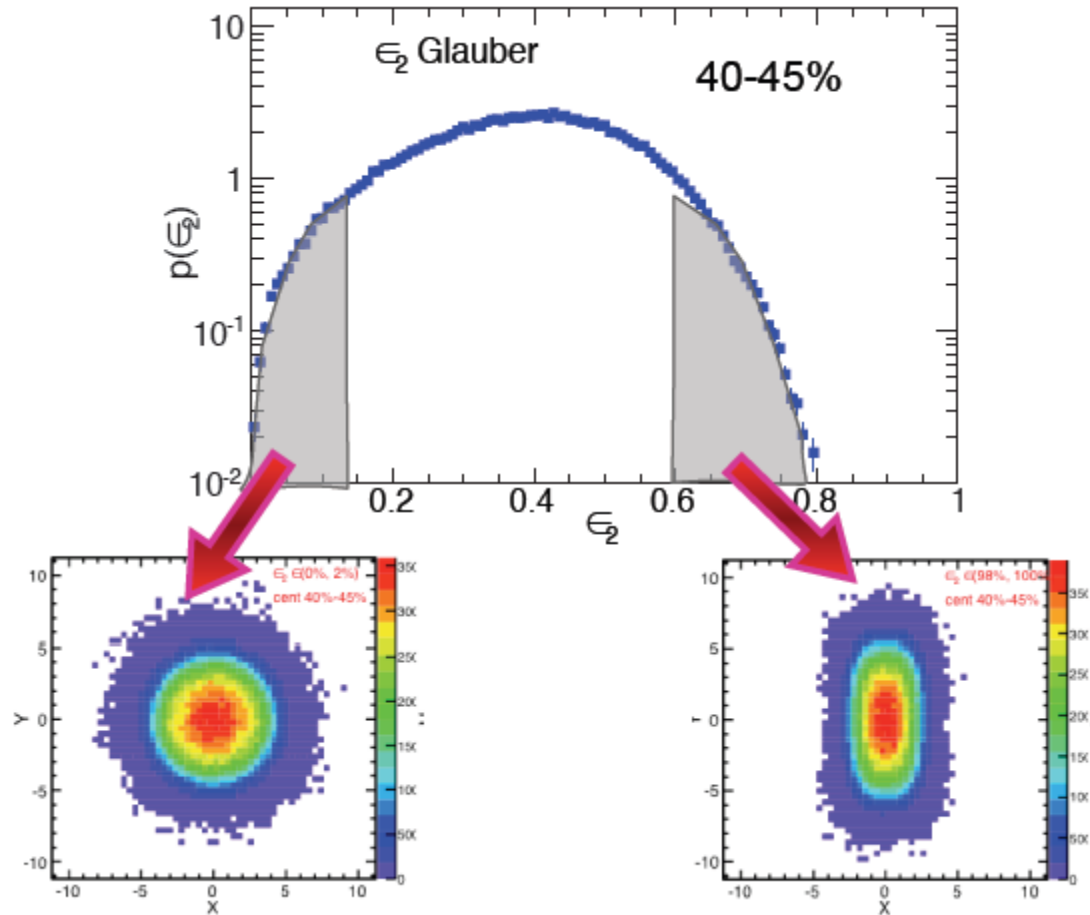
# Shear viscosity



The shear force is given as  $F = \eta A v / d$

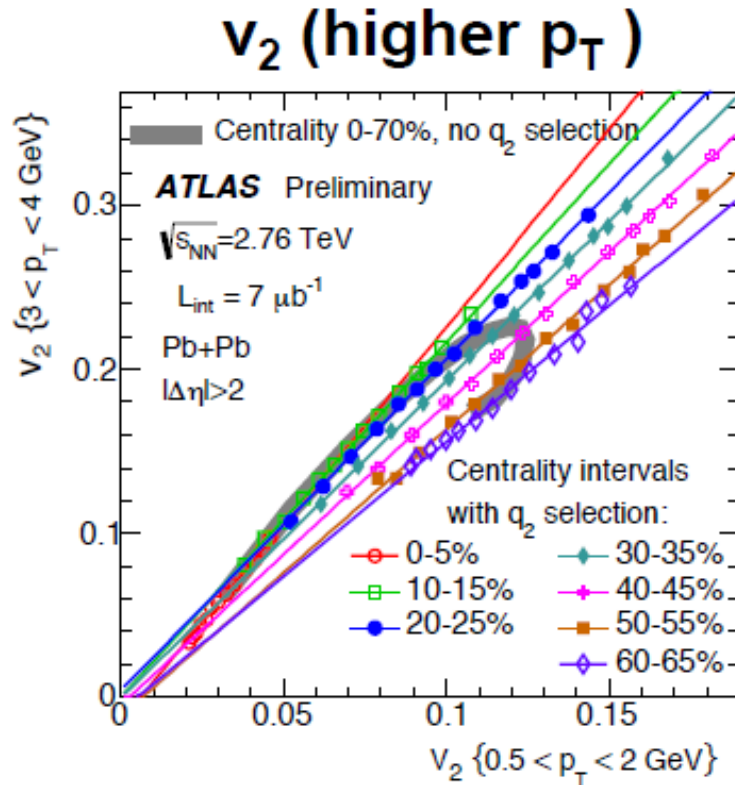
The shear viscosity-to-entropy density ratio,  $\eta/s$ , is a unitless quantity for characterizing fluids

# How ideal is it? (1/2)

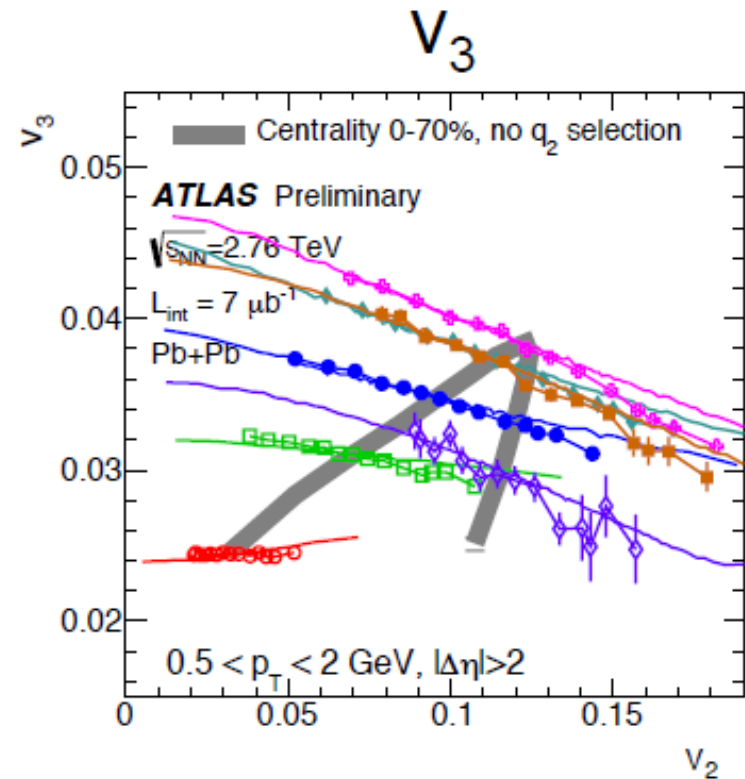


Ideal = reversible  $\Rightarrow$  linear response to initial state ( $v_2 \sim k^* \epsilon_2$ )  
 $\Rightarrow$  fluctuations should be equal to geometrical

# How ideal is it? (2/2)



Left:  $v_2$  response is linear at low and high  $p_T$  (=global property)



Right:  $v_2$  and  $v_3$  correlations are given by pure Glauber model! (solid curves)



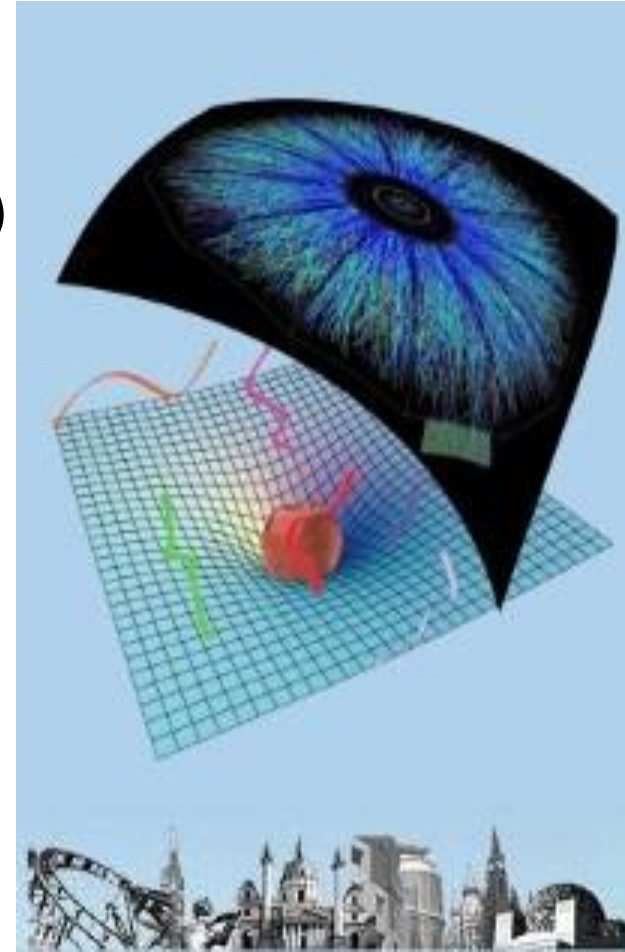
# The QGP is less like a crowd and more like a synchro team



- Big theoretical challenges:
  - Why is the QGP behaving like a liquid? (next slide)
  - how to go from initial random collisions to organized state in a VERY short time ( $<1\text{fm}/c \sim 10^{-23}\text{s}$ ). This remains to be understood

# How to understand this? AdS-CFT

- How to reconcile nearly ideal fluid with energy density like a relativistic gas?
- AdS-CFT correspondence (conjecture)  
J.M. Maldacena,  
Adv.Theor.Math.Phys.2:231-252, 1998,  
~10,000 citations on inspire=most cited
- Duality between **weakly** coupled gravity like theory (AdS) and **strongly** coupled QCD like theory (CFT)
- QCD like theory, but  
conformal (no confinement, no running coupling) (like QGP!)  
infinite Ncolors  
SUSY





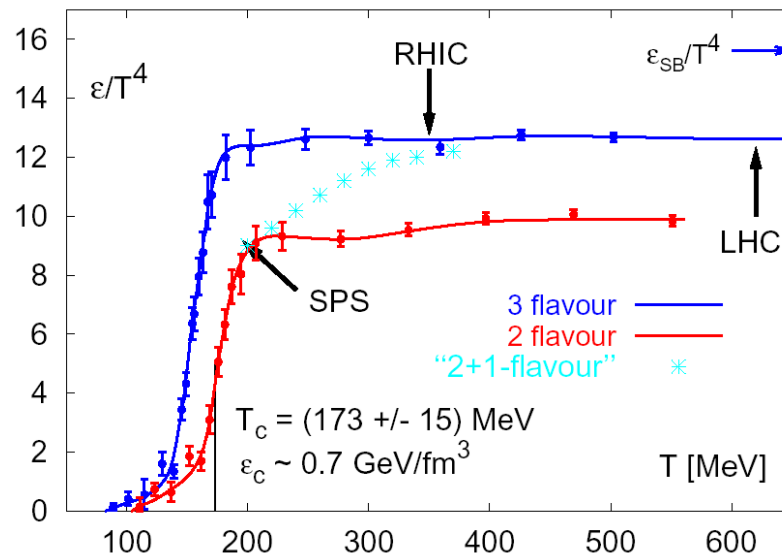
# AdS-CFT

- Two very important results:

Conjectured bound on shear viscosity:  $\eta/s \geq 1/(4\pi) \sim 0.08$

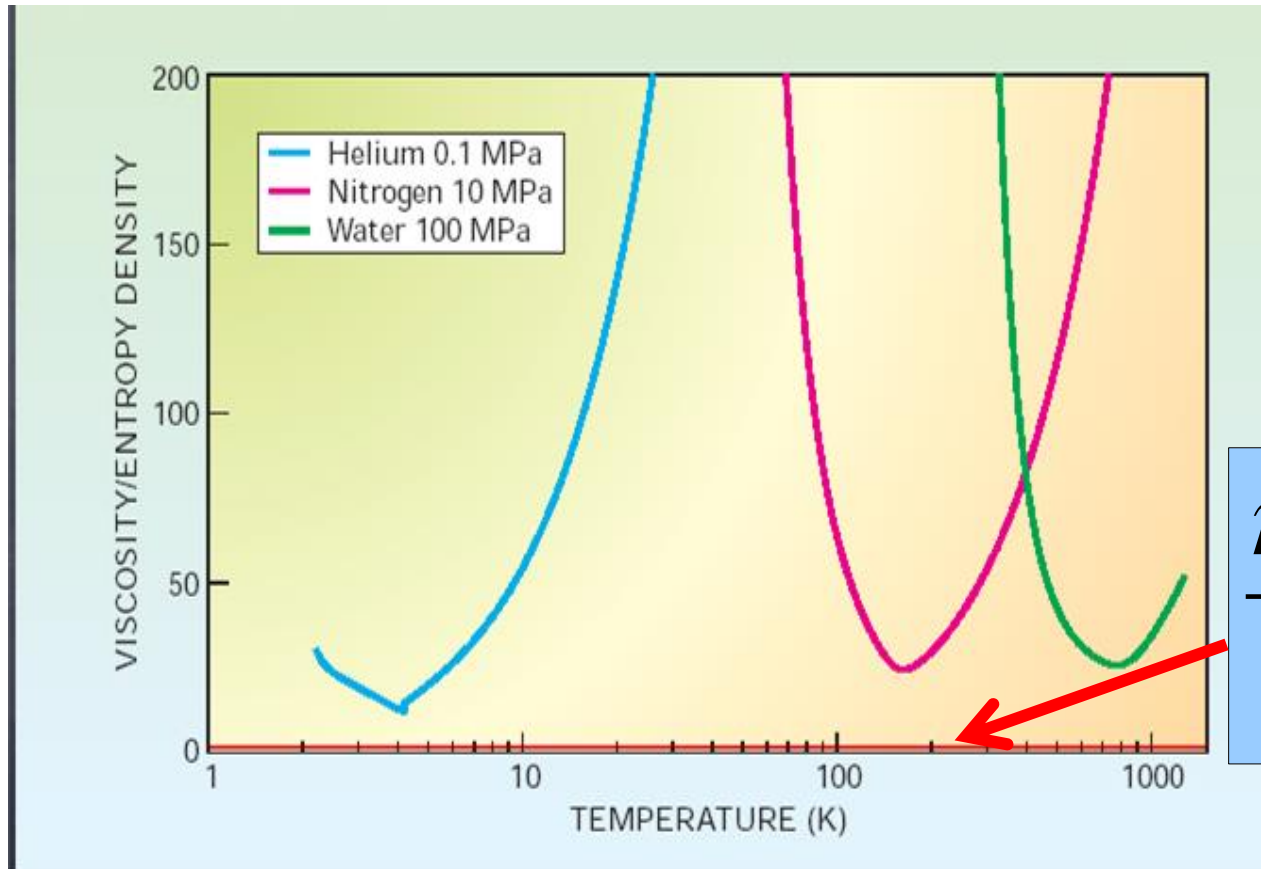
- Viscosity in strongly interacting quantum field theories from black hole physics, P. Kovtun, D.T. Son, A.O. Starinets, Phys.Rev.Lett. 94 (2005) 111601. (~1400 citations on inspire.)

Possibility of infinitely strong coupling at energy density of  $\frac{3}{4}$  SB gas





# The QGP fluid compared to other fluids

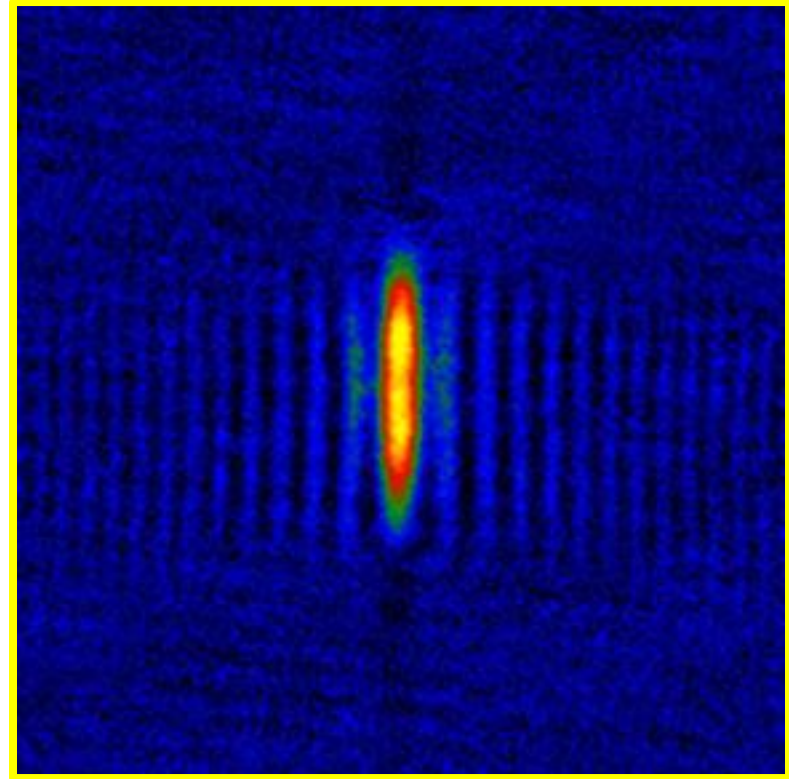
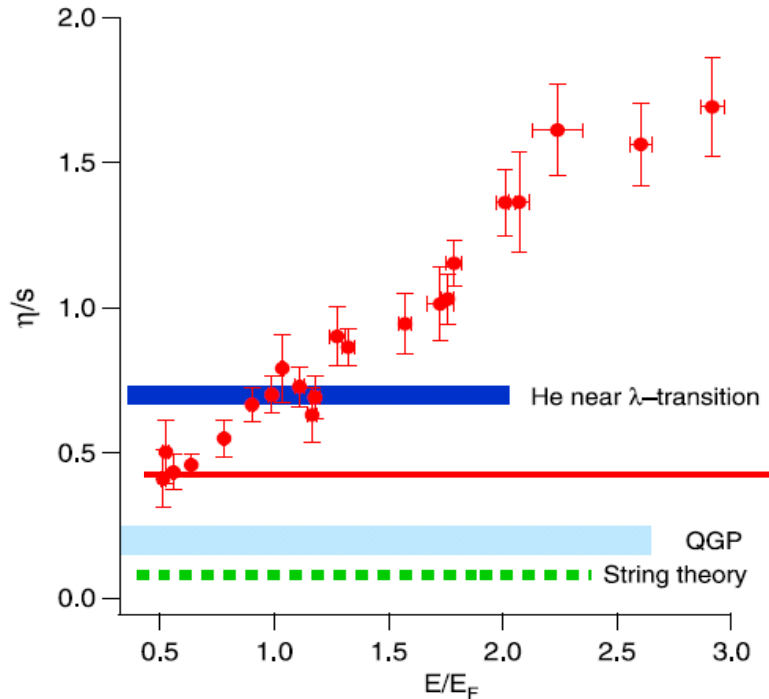


$$\frac{\eta}{s} \geq \frac{\hbar}{4\pi} \frac{1}{k_B}$$



# The QGP fluid compared to other fluids

Strongly interacting ultra-cold Li atoms released from a trap




$$\eta/s \sim 7 \times 1/4\pi$$

<http://www.physics.ncsu.edu/jet/index.html>



# The QGP fluid compared to other fluids

PRL 103, 025301 (2009)

 Selected for a Viewpoint in *Physics*  
PHYSICAL REVIEW LETTERS

week ending  
10 JULY 2009

## Graphene: A Nearly Perfect Fluid

Markus Müller,<sup>1</sup> Jörg Schmalian,<sup>2</sup> and Lars Fritz<sup>3</sup>

<sup>1</sup>*The Abdus Salam International Center for Theoretical Physics, Strada Costiera 11, 34014 Trieste, Italy*

<sup>2</sup>*Ames Laboratory and Department of Physics and Astronomy, Iowa State University, Ames, Iowa 50011, USA*

<sup>3</sup>*Department of Physics, Harvard University, Cambridge, Massachusetts 02138, USA*

(Received 24 March 2009; revised manuscript received 18 May 2009; published 6 July 2009)

Hydrodynamics and collision-dominated transport are crucial to understand the slow dynamics of many correlated quantum liquids. The ratio  $\eta/s$  of the shear viscosity  $\eta$  to the entropy density  $s$  is uniquely suited to determine how strongly the excitations in a quantum fluid interact. We determine  $\eta/s$  in clean undoped graphene using a quantum kinetic theory. As a result of the quantum criticality of this system the ratio is smaller than in many other correlated quantum liquids and, interestingly, comes close to a lower bound conjectured in the context of the quark gluon plasma. We discuss possible consequences of the low viscosity, including preturbulent current flow.

DOI: 10.1103/PhysRevLett.103.025301

PACS numbers: 05.60.Gg, 71.10-w, 73.23.-b, 81.05.Uw





# Conclusions about soft physics

- The medium produced in heavy-ion collisions behaves like a nearly perfect liquid. In fact like the most perfect liquid we know!
- This was completely unexpected based on lattice QCD results
- We can get some insight into the liquid nature from AdS-CFT but so far this does not give a full picture e.g. it does not yet describe jet quenching

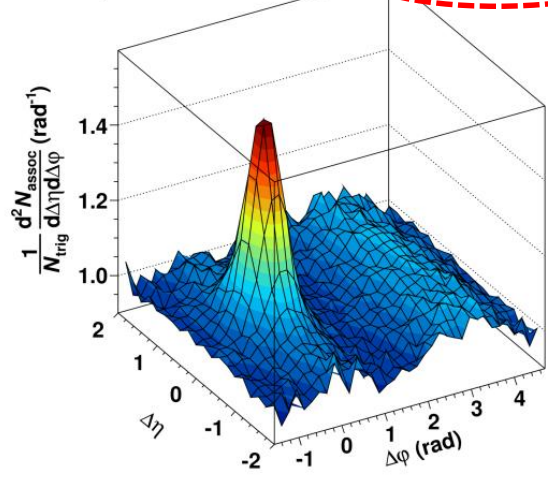


# COLLECTIVITY IN SMALL SYSTEMS

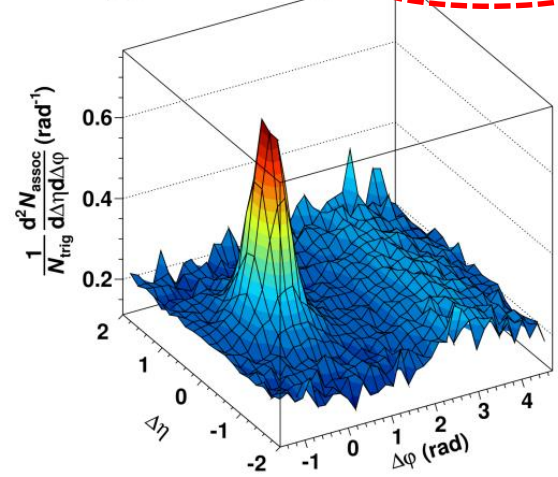


# Collectivity in p-Pb collisions? The rise of the double ridge

$2 < p_{T, \text{trig}} < 4 \text{ GeV}/c$   
 $1 < p_{T, \text{assoc}} < 2 \text{ GeV}/c$   
 p-Pb |  $s_{NN} = 5.02 \text{ TeV}$   
 0-20%



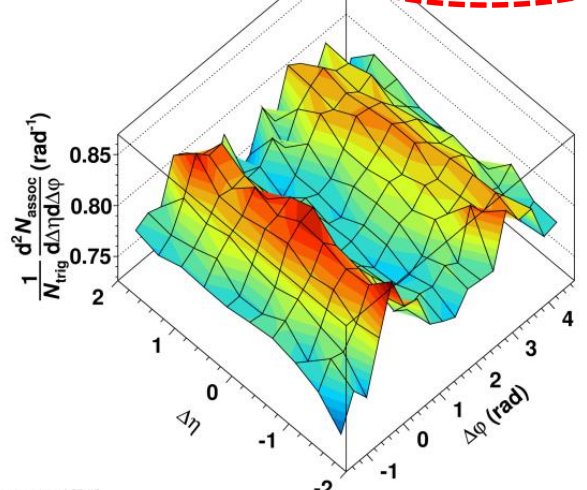
$2 < p_{T, \text{trig}} < 4 \text{ GeV}/c$   
 $1 < p_{T, \text{assoc}} < 2 \text{ GeV}/c$   
 p-Pb |  $s_{NN} = 5.02 \text{ TeV}$   
 60-100%



—

==

ALI-PUB-46228  
 $2 < p_{T, \text{trig}} < 4 \text{ GeV}/c$   
 $1 < p_{T, \text{assoc}} < 2 \text{ GeV}/c$   
 p-Pb |  $s_{NN} = 5.02 \text{ TeV}$   
 (0-20%) - (60-100%)



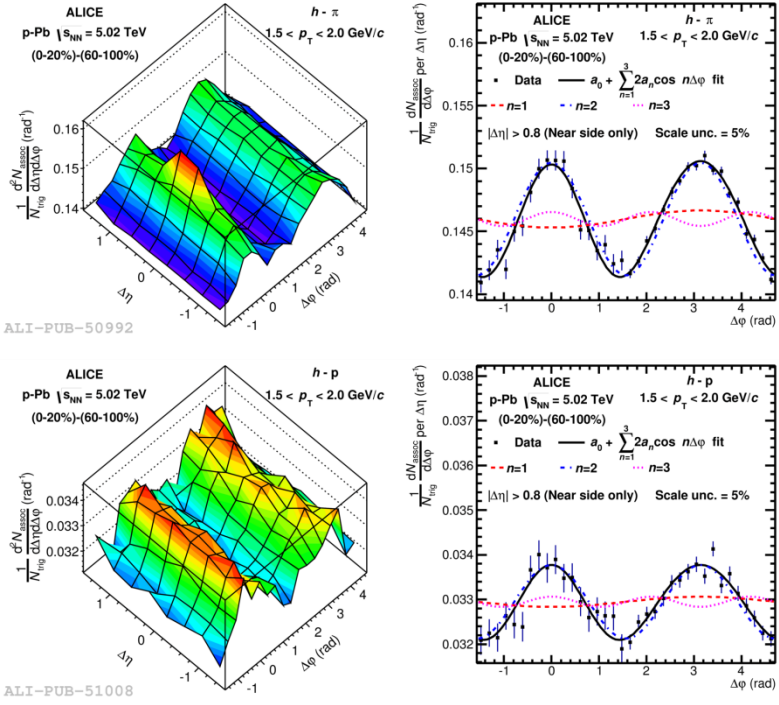
ALICE: Physics Letters B 719 (2013)

- Double ridge structure reminiscent of azimuthal flow in Pb-Pb collisions

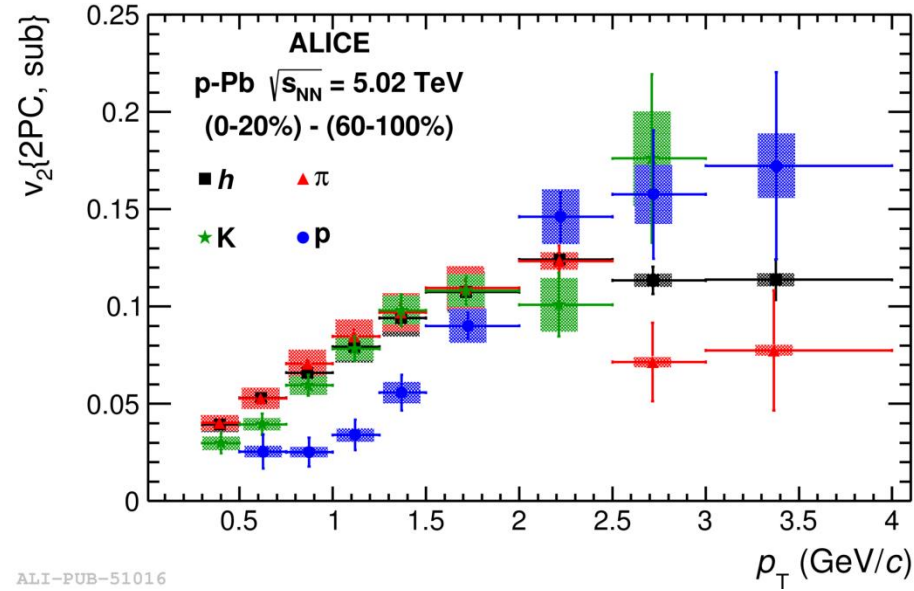


# Using Particle IDentification to study the double ridge

Heavy-ion physics and the QGP (P. Christiansen, Lund)



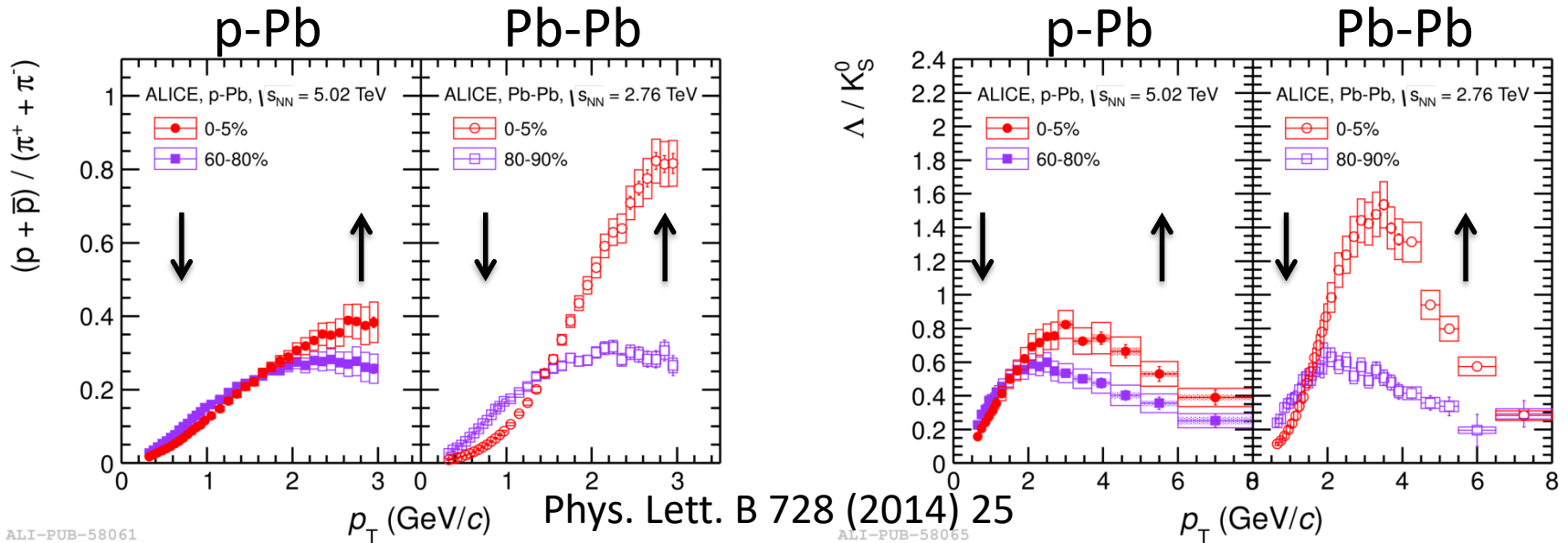
## Fourier coefficients:



ALI-PUB-51016  
Phys. Lett. B 726 (2013) 164–177

- Clear mass ordering suggests flow

# Particle ratios in p-Pb and Pb-Pb show similar features



- Characteristic evolution of  $p/\pi$  and  $\Lambda / K_S^0$  with multiplicity is reminiscent of Pb-Pb where it is believed to be due to radial flow
- NB! The solid boxes for p-Pb ratios indicate the uncorrelated systematic error  
 $\Rightarrow$  the relative trend can be measured rather precisely

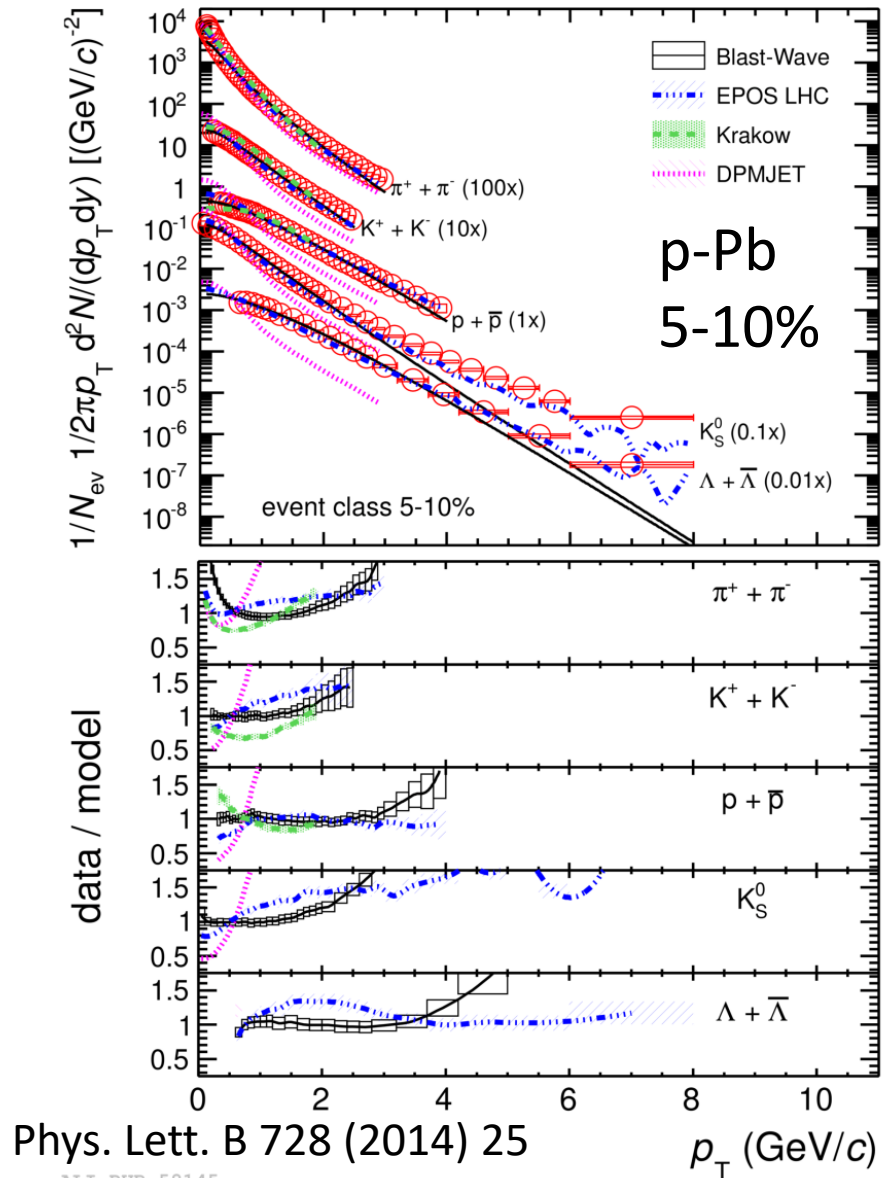






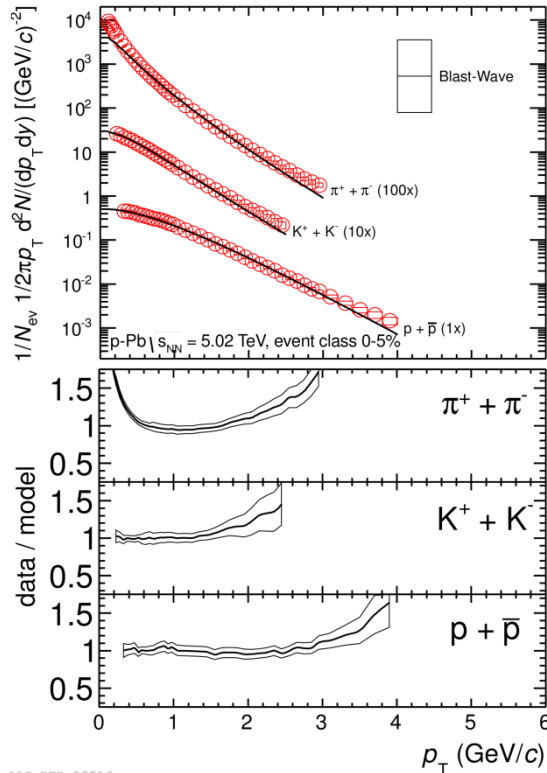
# Which models can capture these features?

- Models
  - Blast-wave (next slides)
  - EPOS LHC (full event generator including hydro)
  - Krakow (hydro calculation focused on low  $p_T$ )
  - DPMJET (PHOJET pp + nuclei via Glauber-Gribov theory)
- Only models which employ hydrodynamics can describe the  $p_T$  spectra



# A blast wave study of the data

p-Pb 0-5%



- Simultaneous fits
  - $\pi$ :  $0.5 < p_T < 1.0$  GeV/c
  - K:  $0.2 < p_T < 1.5$  GeV/c
  - pp:  $0.3 < p_T < 1.5$  GeV/c
  - p:  $0.3 < p_T < 3.0$  GeV/c
  - pp:  $0.5 < p_T < 2.5$  GeV/c

- Adding  $K_S^0$  ( $0.0 < p_T < 1.5$  GeV/c) and  $\Lambda$  ( $0.6 < p_T < 2.0$  GeV/c) does not significantly change extracted parameters

Schnedermann et al, Phys. Rev. C 48, 2462 (1993)

$$\frac{dN}{p_{\perp} dp_{\perp}} \propto \int_0^R r dr m_{\perp} I_0 \left( \frac{p_{\perp} \sinh \rho}{T_{\text{kin}}} \right) K_1 \left( \frac{m_{\perp} \cosh \rho}{T_{\text{kin}}} \right)$$

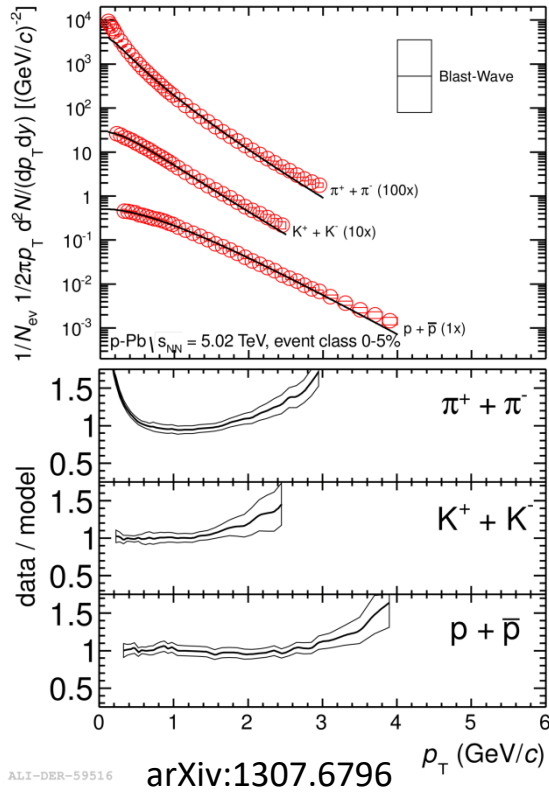
$$\begin{aligned} \rho &= \tanh^{-1} \beta_T \\ \beta_T &= \beta_S (r/R)^n \\ \langle \beta_T \rangle &= \frac{2}{2+n} \beta_S \end{aligned}$$



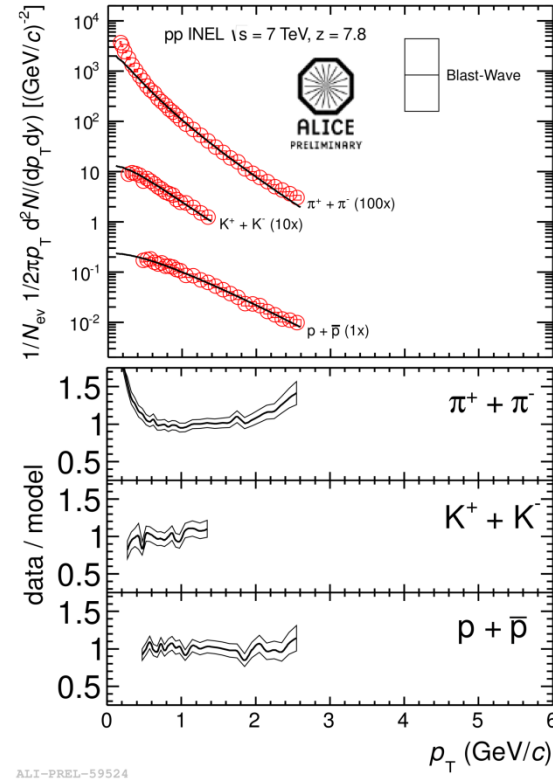


# A blast wave study of the data

p-Pb 0-5%



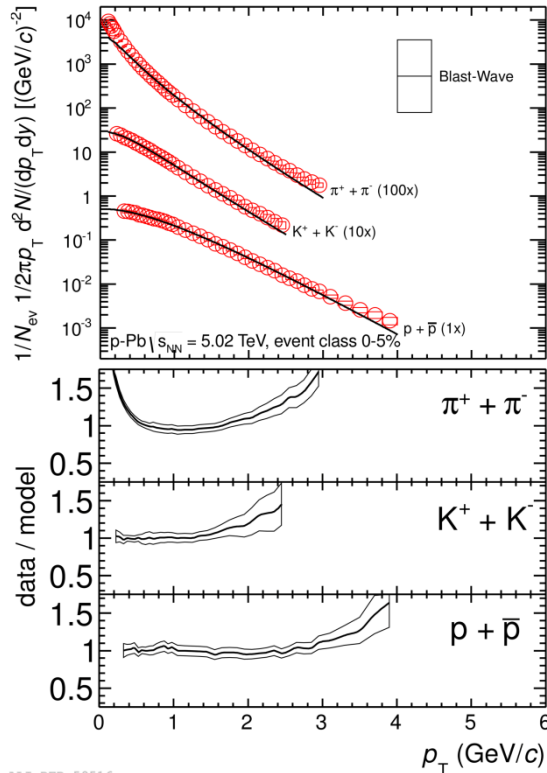
pp highest mult vs = 7 TeV



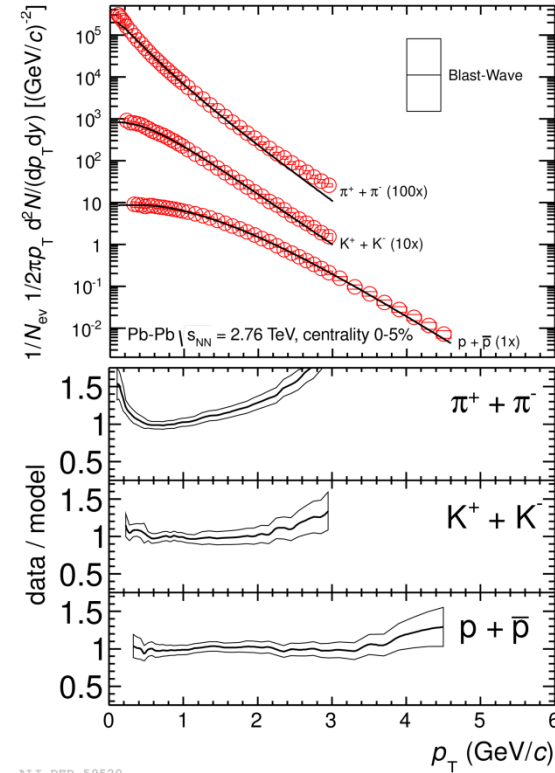
The description of p-Pb and pp data by the blast-wave fit is reasonable without being excellent

# A blast wave study of the data

p-Pb 0-5%



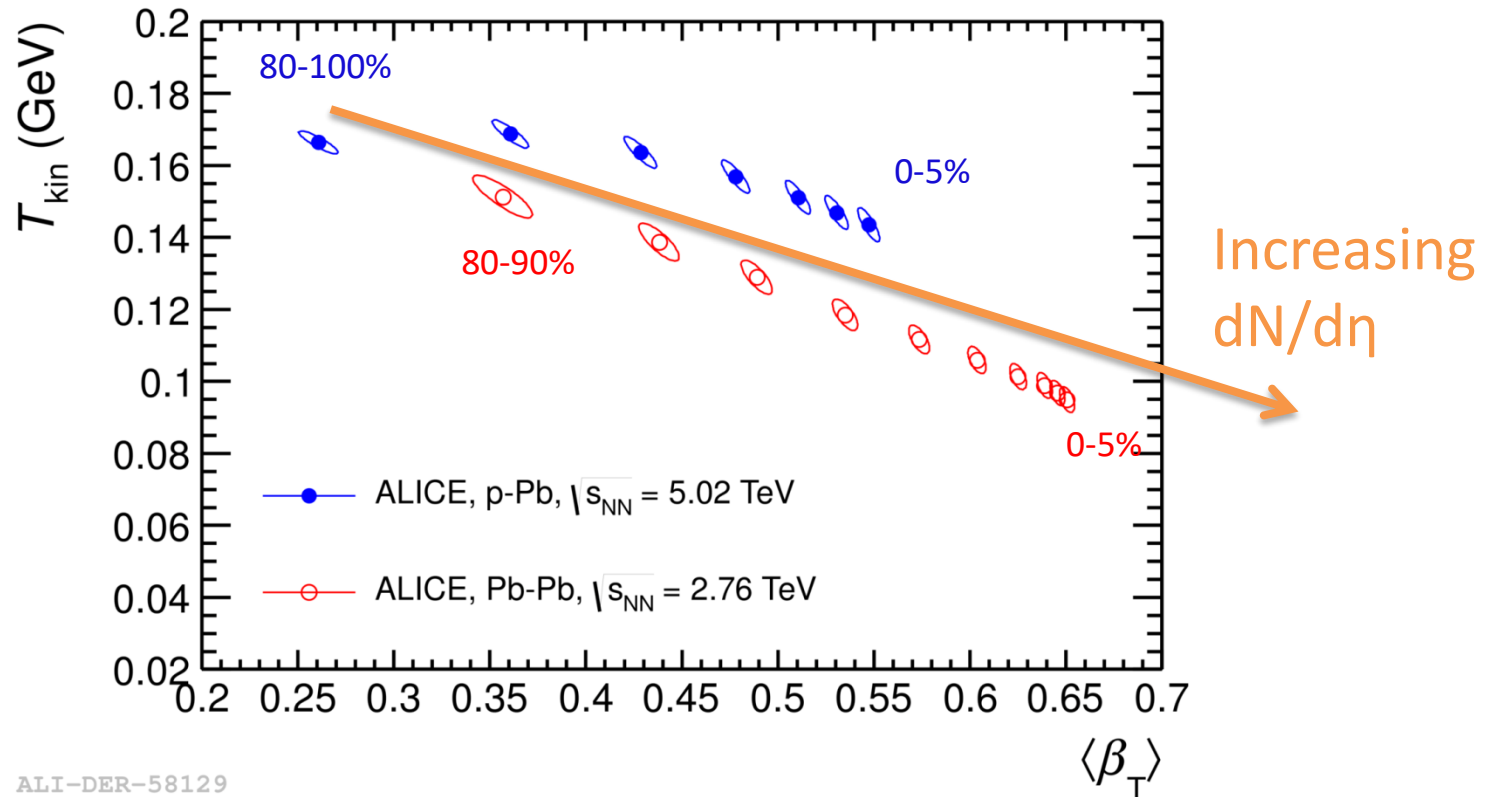
Pb-Pb 0-5%  $\sqrt{s_{NN}} = 2.76$  TeV



The  $p_T$  region where the blast-wave fit describes the data is in general broader for Pb-Pb



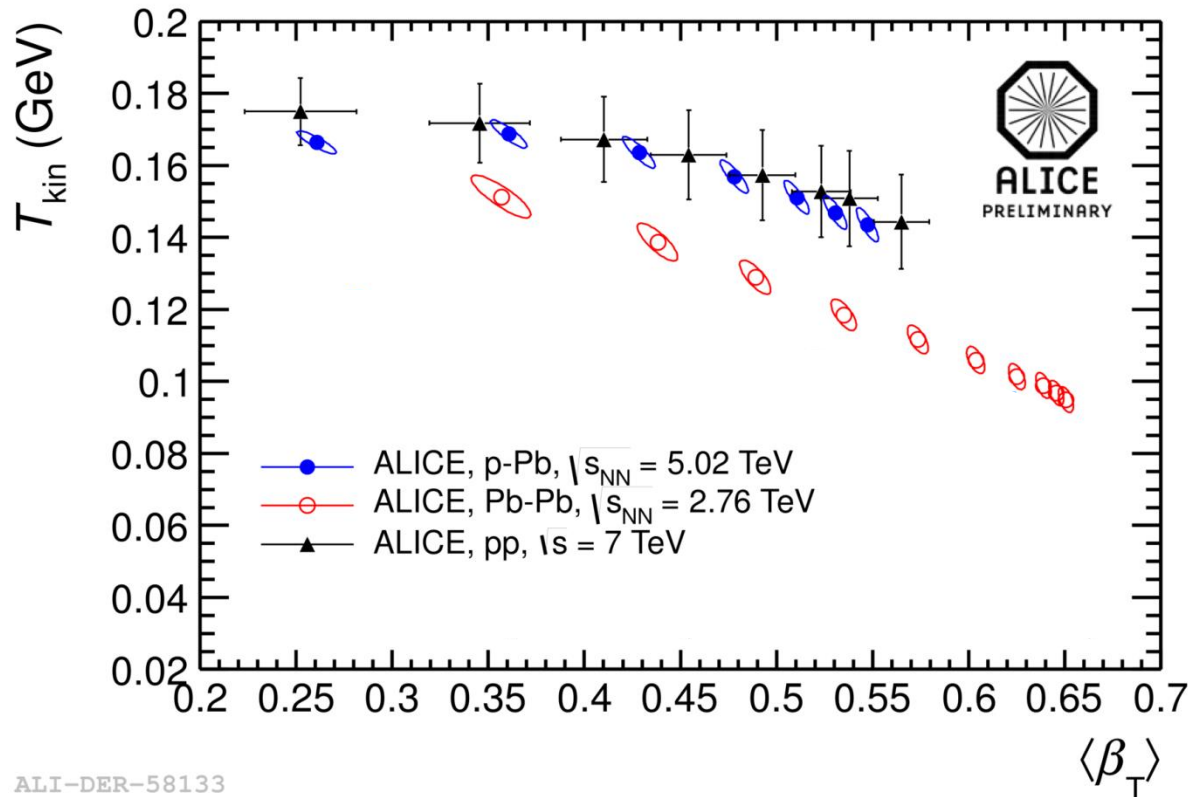
# Summarizing the results from the blast-wave studies



There is a strong common trend between the parameters extracted from Pb-Pb and p-Pb



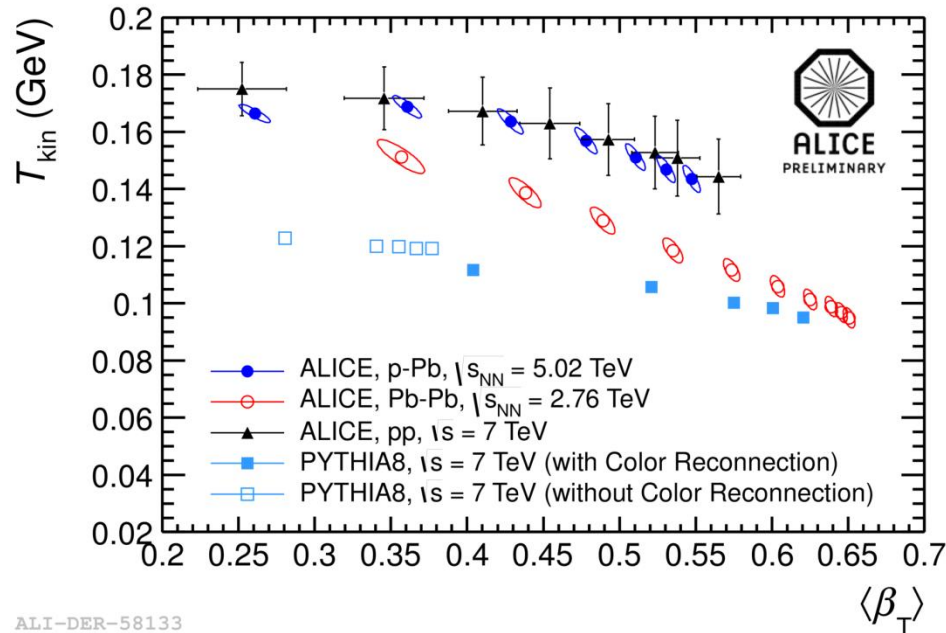
# Summarizing the results from the blast-wave studies



Even the pp data seems to follow the same trend!  
 It seems that if we ascribe the change in spectral shape to radial flow in p-Pb then the same can be done in pp



# Summarizing the results from the blast-wave studies

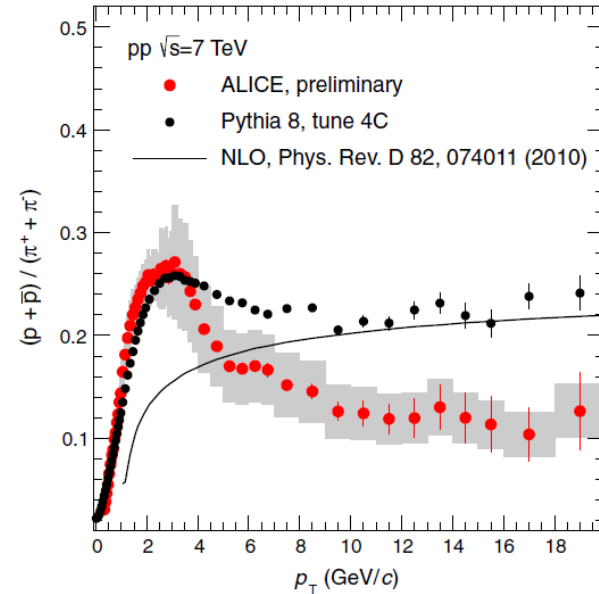
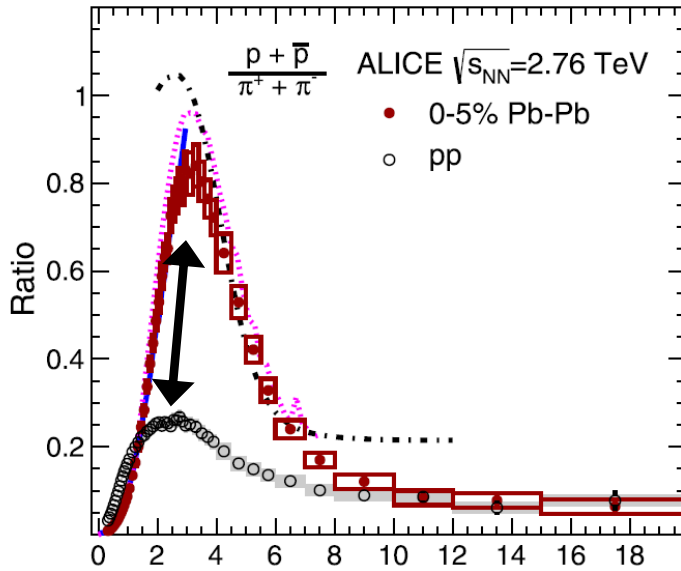


BUT also simulated PYTHIA8 pp events follow a qualitatively similar trend when Color Reconnection (CR) is enabled  
 CR has been shown to mimic radial flow but without requiring the formation of a medium

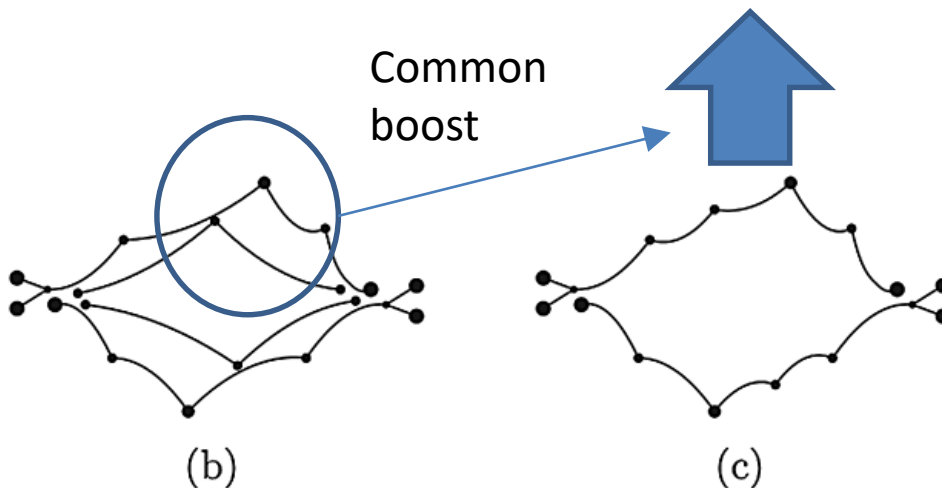


# The “flow peak” in pp

ALICE , PLB 736,  
196 (2014)



Realized that Color Reconnection in PYTHIA gives rise to flow like boosts, A. O. Velasquez, P. Christiansen, et al, PRL 111, 042001 (2013). For details, see T. Sjöstrand, arXiv:1310.8073.



Can be interpreted different ways: CR as microscopic model of flow or one needs (hydro) flow in pp.







# Conclusions

- The reason we believe in a QGP
  - We clearly form a medium that behaves as a nearly ideal fluid with a temperature in the range given by LQCD
  - It dissolves quarkonia as expected indicating that there are strong screening effects = high density of colored objects
  - It quenches jets and gives rise to energy losses of order GeV/fm which is as expected from color fields
- The challenge to this paradigm is that we now observe medium like effects in small systems
  - We need to understand this better!



# What I did not cover

- Charged multiplicity and hadronization
  - Color Glass Condensate and Statistical Model
- Quark scaling of elliptic flow
  - Recombination of quark like degrees of freedom (picture is more complicated at LHC)
- Quenching and flow of heavy quarks
  - Suggests that there are also large collisional energy loss



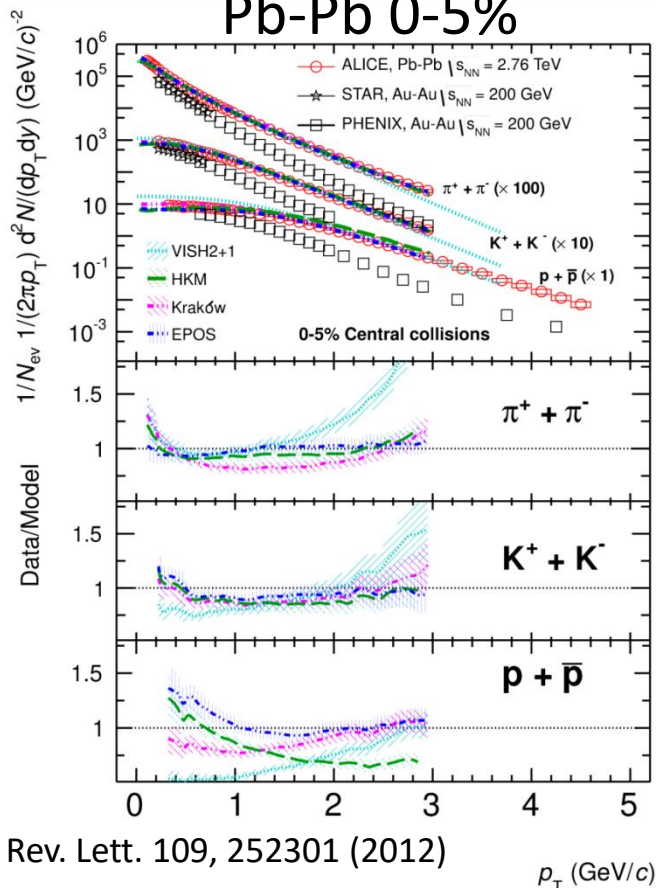
# Why are the small systems so critical to understand

- Before we were happy to have
  - Small systems = QCD
    - Phenomenology: PYTHIA
  - Large systems = QGP
    - Phenomenology: Glauber + hydro + stat. model + some hadronic rescattering



# Geometry: is the Pb-Pb centrality evolution described by hydro?

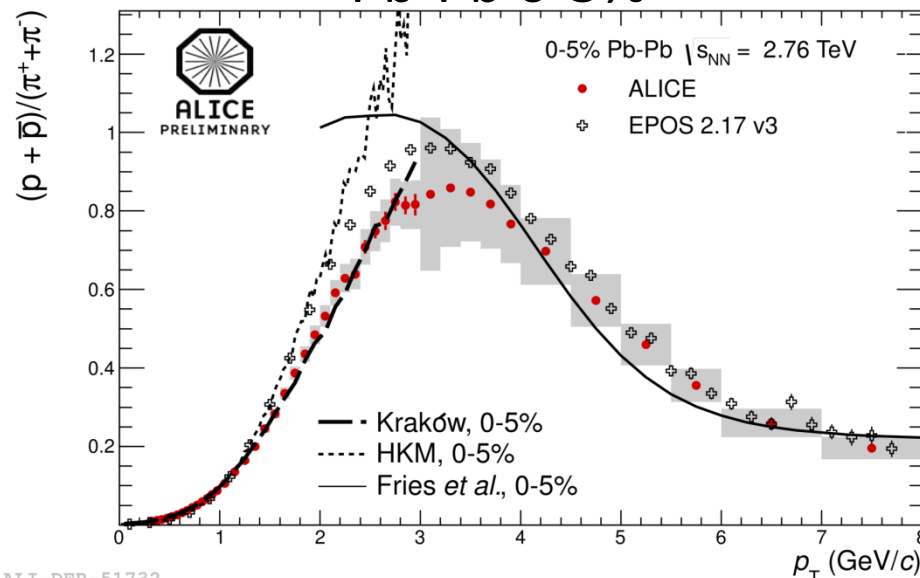
Pb-Pb 0-5%



Phys. Rev. Lett. 109, 252301 (2012)

- Near ideal hydrodynamics with some implementation of the hadronic phase describes well  $p_T$  spectra in central collisions

Pb-Pb 0-5%



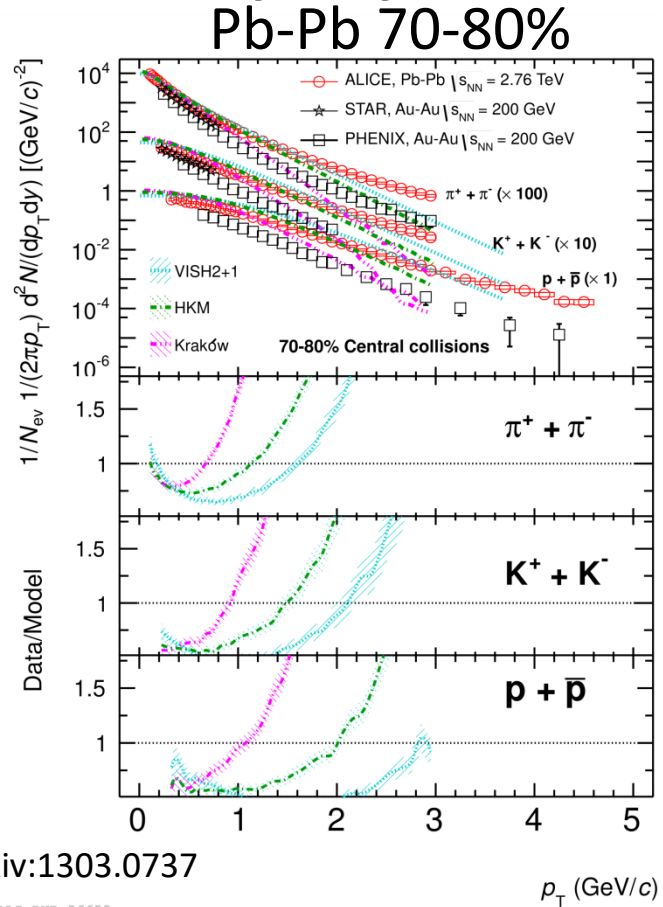
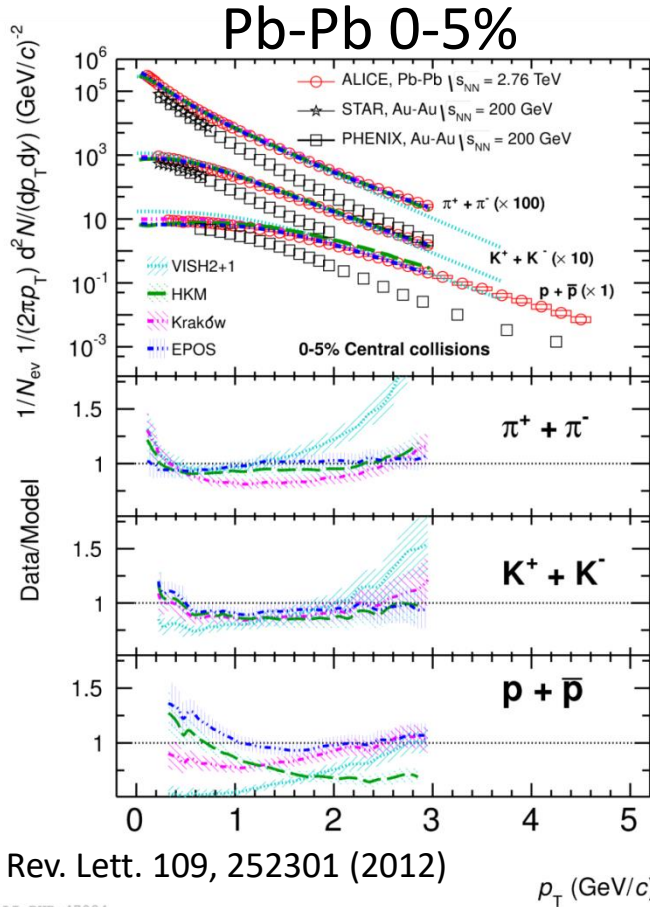
ALI-DER-51732

- Krakow: PRC85, 064915 (2012)
- HKM: PRC87, 024914 (2013)
- Fries: PRL90, 202303 (2003) and private communication
- EPOS: PRL109, 102301 (2012) and private communications

Heavy-ion physics and the QGP (P. Christiansen, Lund)

ALI-PUB-47084

# Geometry: is the Pb-Pb centrality evolution described by hydro?



Phys. Rev. Lett. 109, 252301 (2012)

arXiv:1303.0737

ALI-PUB-47084

ALI-PUB-56658

- The same models fail to describe the  $p_T$  spectra in peripheral collisions
- Typically hydro has not been expected to work in peripheral collisions but if it is at work in p-Pb and pp collisions should it not work there?





# Why is this a problem

- Our description of the medium is not only critical for the low  $p_T$  observables, the jet quenching also relies on our medium properties
- If we cannot resolve this issue then we are a bit stuck
- Need to understand if a medium is also created in small systems!
  - My own view: need to look for energy loss in small systems

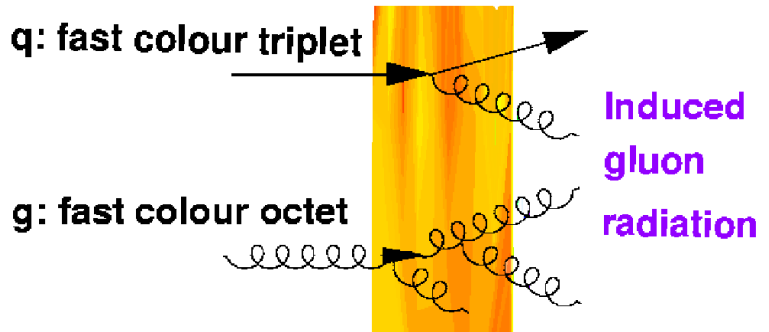
# Backup slides



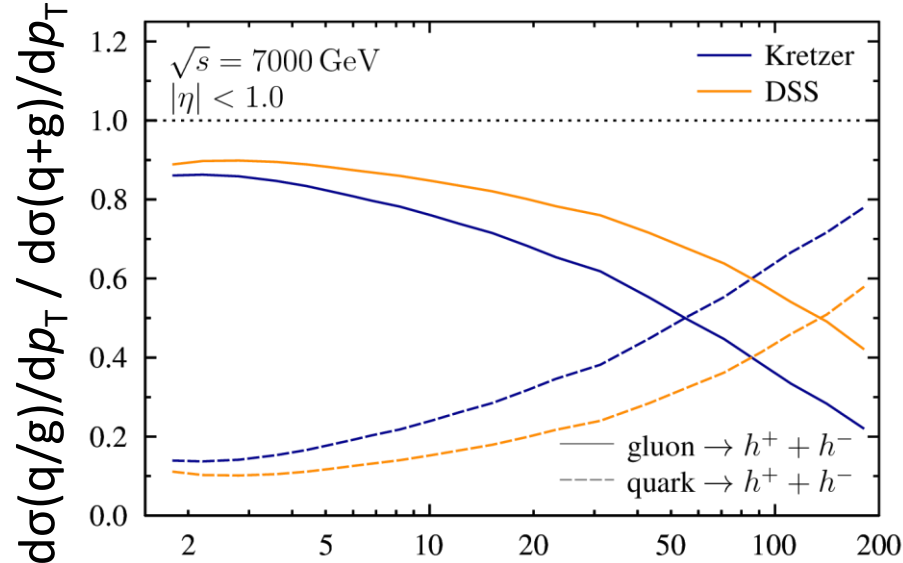


# Motivation: searching for "footprints" of the energy loss

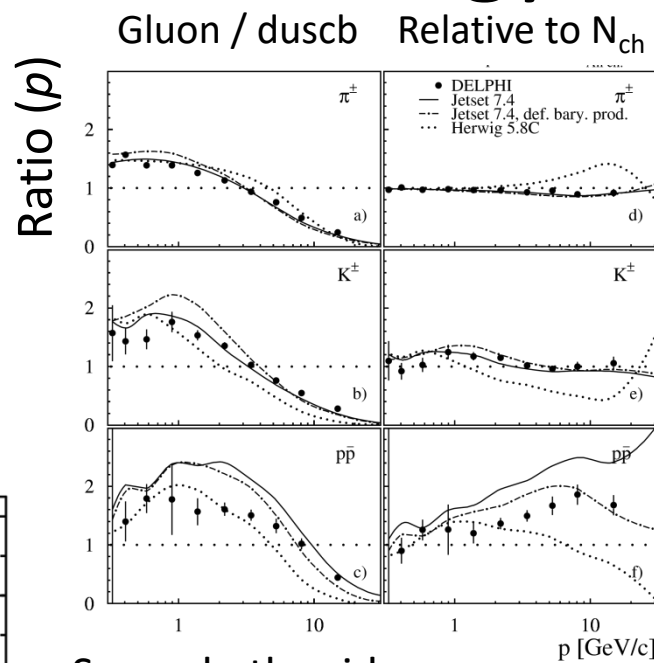
Heavy-ion physics and the QGP (P. Christiansen, Lund)



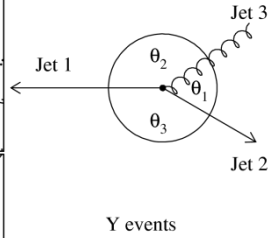
Gluons lose 2 times (color factor) more energy than quarks in the medium



d'Enterria et al, *p\_T* [GeV/c]  
Nucl.Phys. B883, 615 (2014)



DELPHI Collaboration  
EPJC17, 207 (2000)



Several other ideas:

Color flow (Sapeta, Wiedemann, EPJC55, 293, 2008)

Color structure (Aurenche, Zakharov, EPJC71, 1829, 2011)

In medium formation time effects (Bellwied, Markert, PLB691, 208, 2010)

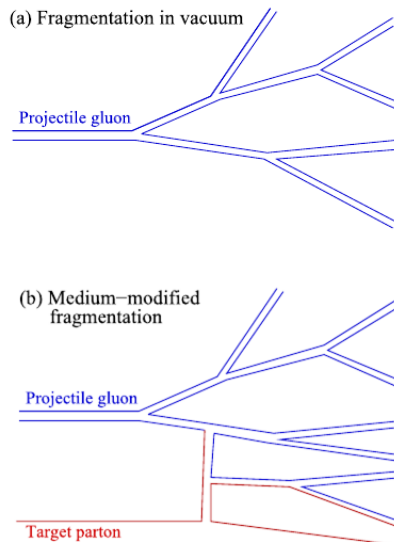
Magnitude of the effects are large (50+%) since it is linked to the large energy loss





# Why expect particle species dependent $R_{AA}$ at high $p_T$ ?

- Large effects at intermediate  $p_T$  – does this effect just disappear?
- The low value of  $R_{AA}$  suggests that most hard partons interacts strongly with the medium



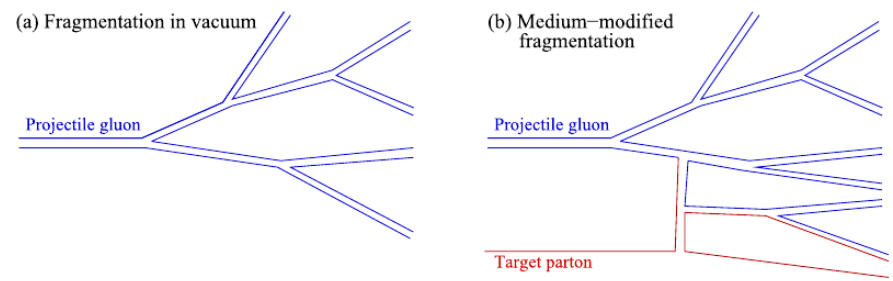
S. Sapeta and U.A. Wiedemann, Eur.Phys.J. C55 (2008) 293:

- Indirect
  - “in all models of radiative parton energy loss, the interaction of a parent parton with the QCD medium transfers color between partonic projectile and target. This changes the color flow in the parton shower and is thus likely to affect hadronization.”
- Direct
  - “In addition, flavor or baryon number could be exchanged between medium and projectile.”



# A general model with particle species dependent modifications

S. Sapeta and U.A. Wiedemann, Eur.Phys.J. C55 (2008) 293



- Effect inside jet
- But for  $p_T \gg 8 \text{ GeV}/c$  we expect all hadrons to belong to jets

