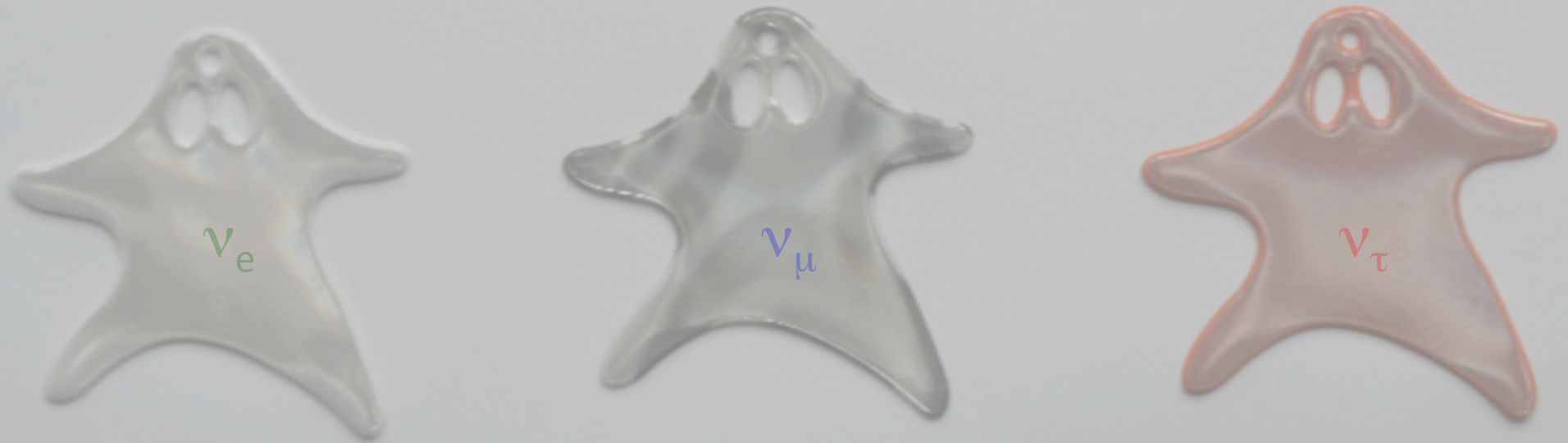


Introduction to Neutrino Physics

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Lecture 4

Neutrinoless double beta decay Part 2

Recap lecture 3 neutrino mass

- Neutrino oscillations don't tell us anything about absolute neutrino masses

- From neutrino oscillations and cosmological observations:

$$50 \text{ meV} < m_\nu < \sim 1 \text{ eV}$$

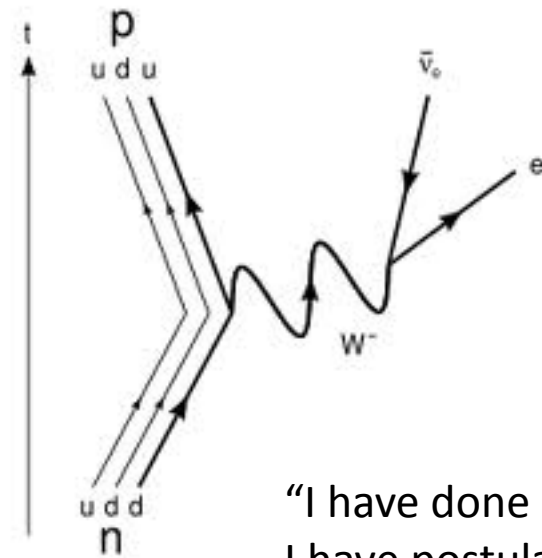
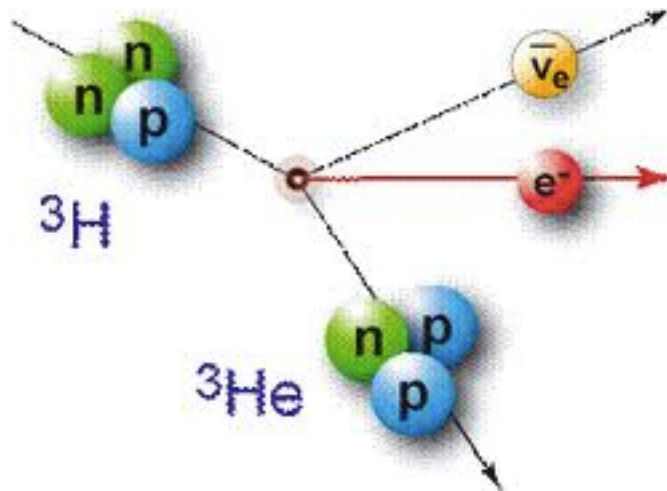
- Straightforward to extend the SM with Dirac masses for neutrinos
 - But this introduces an unnecessary constraint on lepton-number conservation: the neutrino has no electric charge to conserve!
- Can construct a Majorana mass term out of the right-handed field and its charge-conjugate
 - Mixes ν and anti- ν , so no conservation of lepton number
- See-saw mechanism combines Dirac and Majorana to give an extremely light \sim left-handed neutrino and an extremely heavy \sim right-handed neutrino

Outline lecture 4

- Beta decay: single and double, with and without neutrinos
- Experimental considerations for a search for neutrinoless double beta decay
- Tour of a few neutrinoless double beta decay experiments

Single beta decay

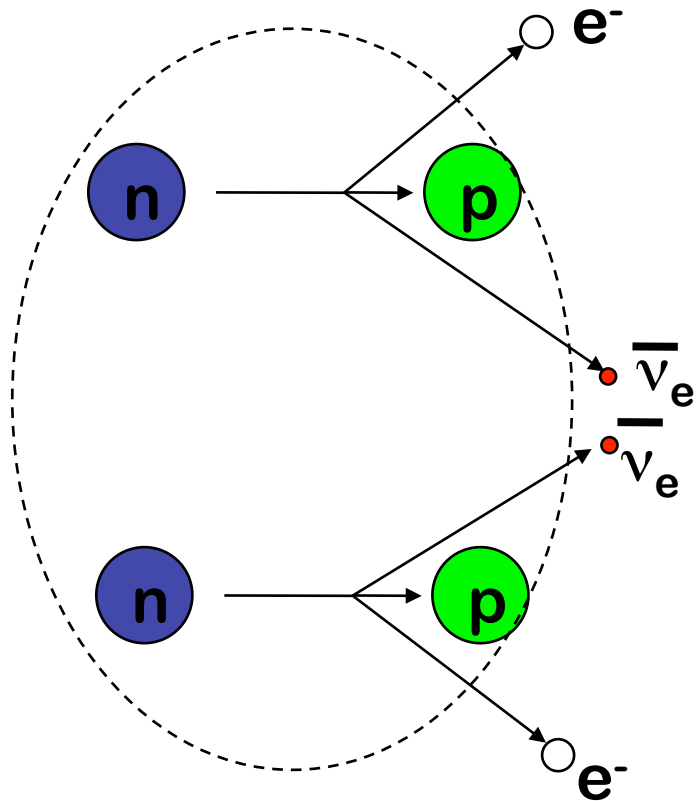
That's where it started...



“I have done a terrible thing,
I have postulated a particle
that can not be detected”
(W. Pauli)

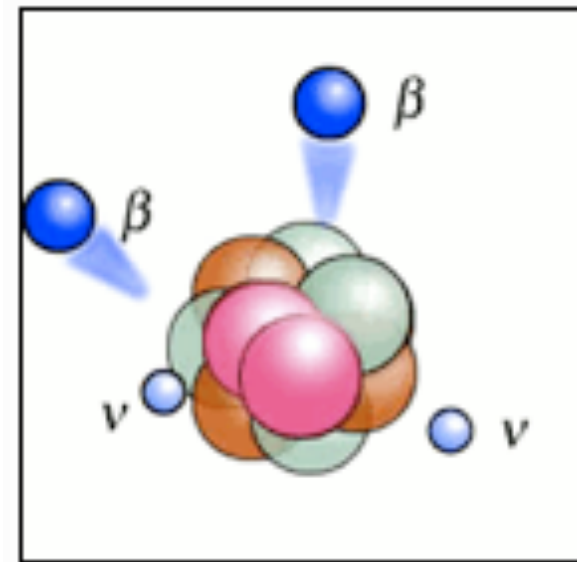
First-order Standard-Model process

Double beta decay



$$\mathcal{N}(A,Z) \rightarrow \mathcal{N}(A,Z+2) + 2e^- + 2\bar{\nu}_e$$

$2\beta^-$ process

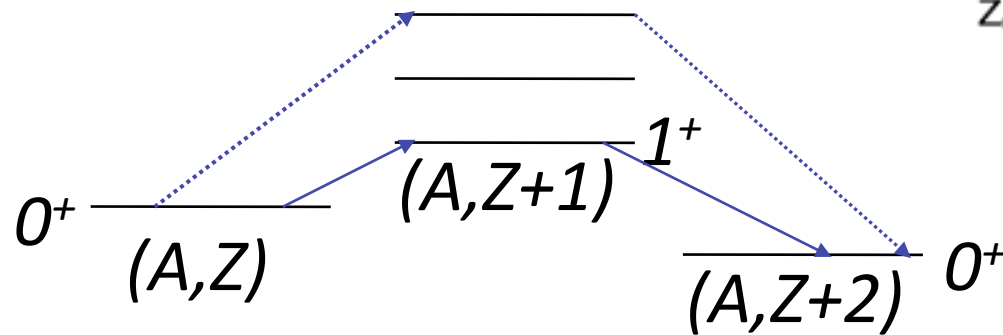
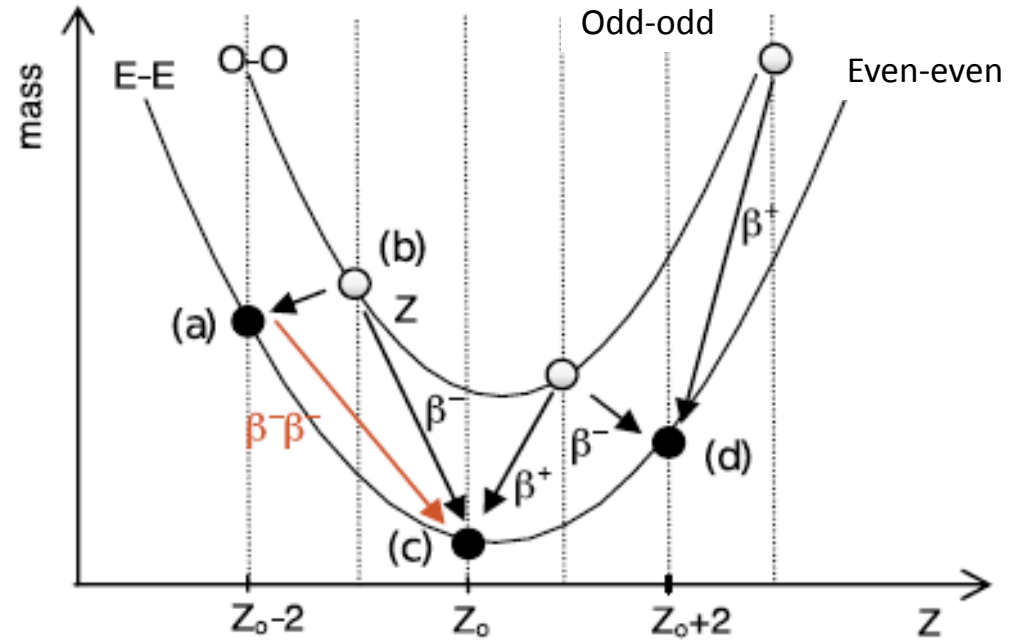


Second-order SM process

Two neutrinos emitted: $2\nu\beta\beta$

Double beta decay

Requires even-even nuclei



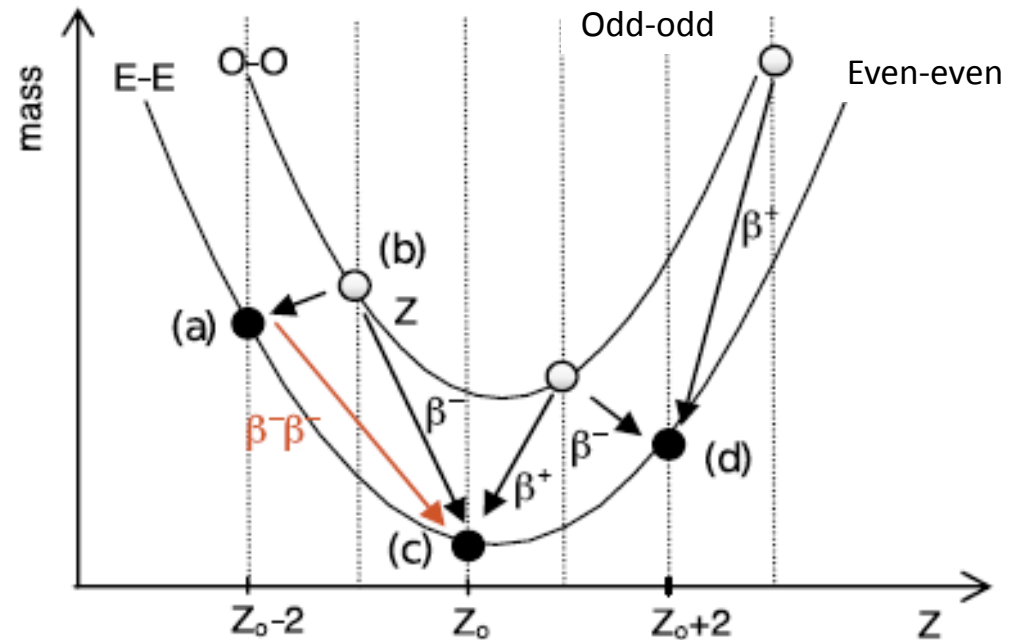
Double beta decay

Rarest kind of radioactive decay:

Only 35 naturally occurring isotopes can decay via $2\beta^-$

Half-lives of $O(10^{16}$ years) and upwards
C.f. age of the universe $O(10^{10}$ years)

Double beta decay

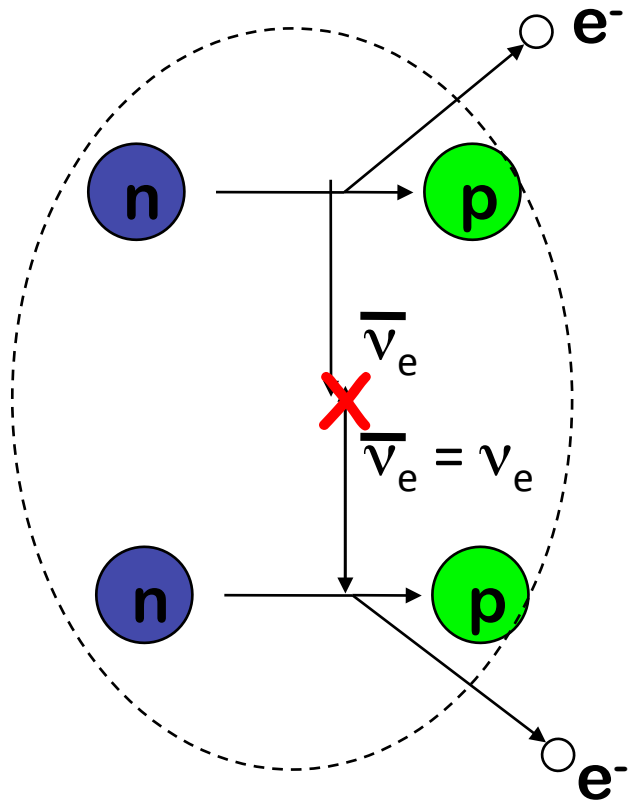


Rarer still:

Can occur also via the $2\beta^+$ process: $\mathcal{N}(A,Z) \rightarrow \mathcal{N}(A,Z-2) + 2e^+ + 2\nu_e$
 and via electron capture (EC): $e^- + \mathcal{N}(A,Z) \rightarrow \mathcal{N}(A,Z-2) + e^+ + 2\nu_e$
 $2e^- + \mathcal{N}(A,Z) \rightarrow \mathcal{N}(A,Z-2) + 2\nu$

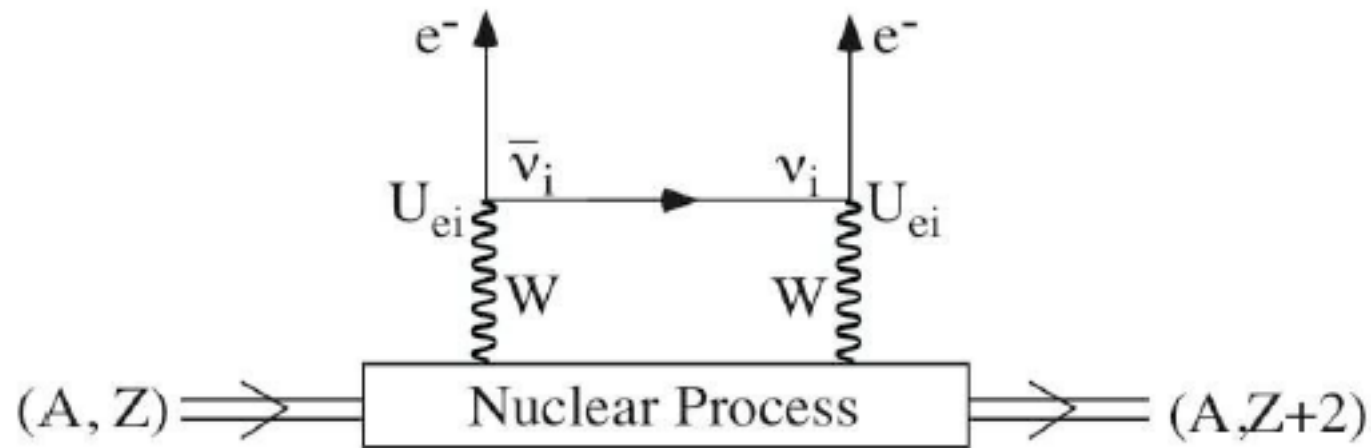
Neutrino-less double beta decay ($0\nu\beta\beta$)

$$\mathcal{N}(A,Z) \rightarrow \mathcal{N}(A,Z+2) + 2e^-$$



- Violates lepton number: $\Delta L = 2$
- Forbidden in the Standard Model

$0\nu\beta\beta$ and neutrino mass



Amplitude for $0\nu\beta\beta$ is proportional to the **effective Majorana mass**:

$$m_{\beta\beta} = \sum_{k=1}^3 U_{ek}^2 m_k = |U_{e1}|^2 m_1 + e^{i\alpha} |U_{e2}|^2 m_2 + e^{i\beta} |U_{e3}|^2 m_3$$

$0\nu\beta\beta$ half-life

$$(T_{1/2})^{-1} = G_{0\nu} |M_{0\nu}|^2 |m_{\beta\beta}|^2 / m_e^2$$

Phase-space factor

Calculable with small uncertainties

Depends on $Q_{\beta\beta}$ and Z

Nuclear Matrix Element

Incorporates knowledge of all nuclear structural effects

Requires accurate nuclear model

Models differ by up to x3

Experimental input could help

The Q value



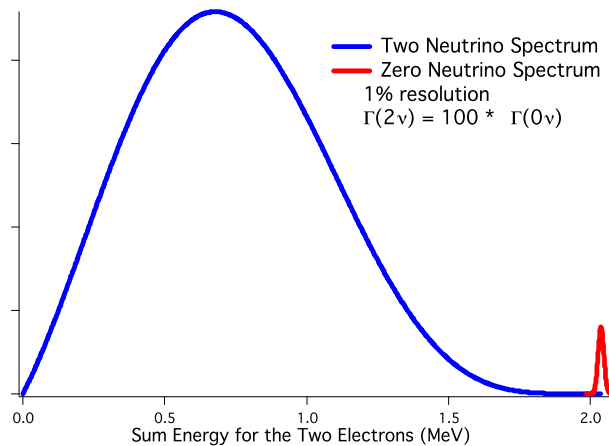
The sum of the kinetic energies of the electrons
is always the same, namely
the mass difference $m_{(A,Z)} - m_{(A,Z+2)}$,
i.e. $Q_{\beta\beta}$ (the “Q value”)

The Q value



The sum of the kinetic energies of the electrons is always the same, namely the mass difference $m_{(A,Z)} - m_{(A,Z+2)}$, i.e. $Q_{\beta\beta}$ (the “Q value”)

$2\nu\beta\beta$ has a continuous spectrum



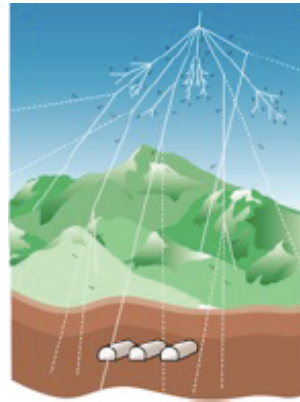
$0\nu\beta\beta$ has a discrete spectrum
At the upper end of the $2\nu\beta\beta$ spectrum

Considerations for a $0\nu\beta\beta$ experiment

- Which isotope to choose?
 - 35 to choose from...
 - Ex ^{76}Ge , ^{130}Te , ^{82}Se , ^{136}Xe , ^{150}Nd
- Dealing with backgrounds
- Detector techniques
- A few example experiments

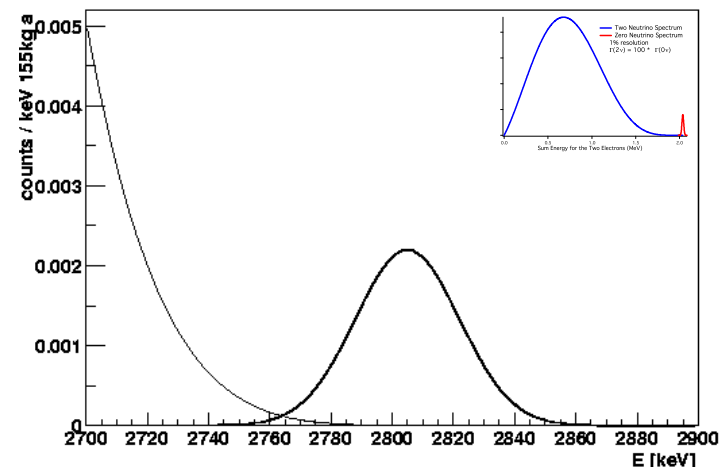
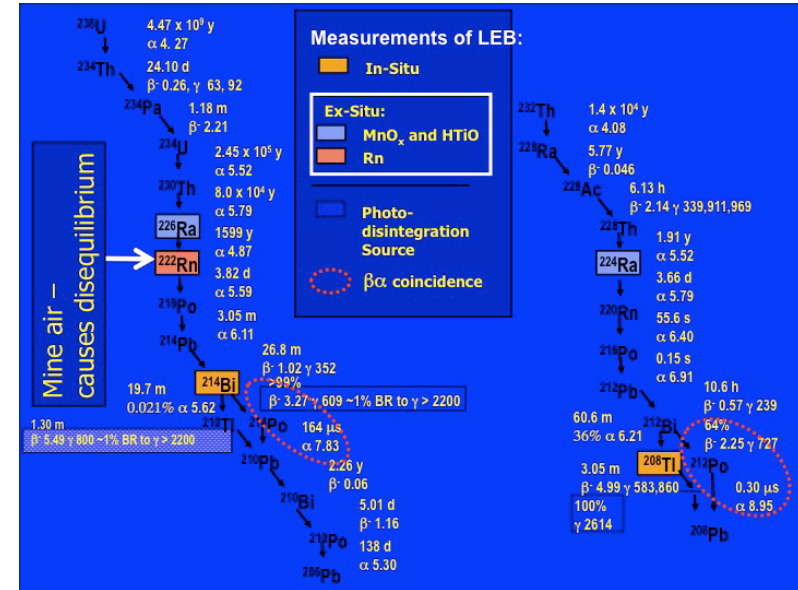
Backgrounds

- Cosmic rays



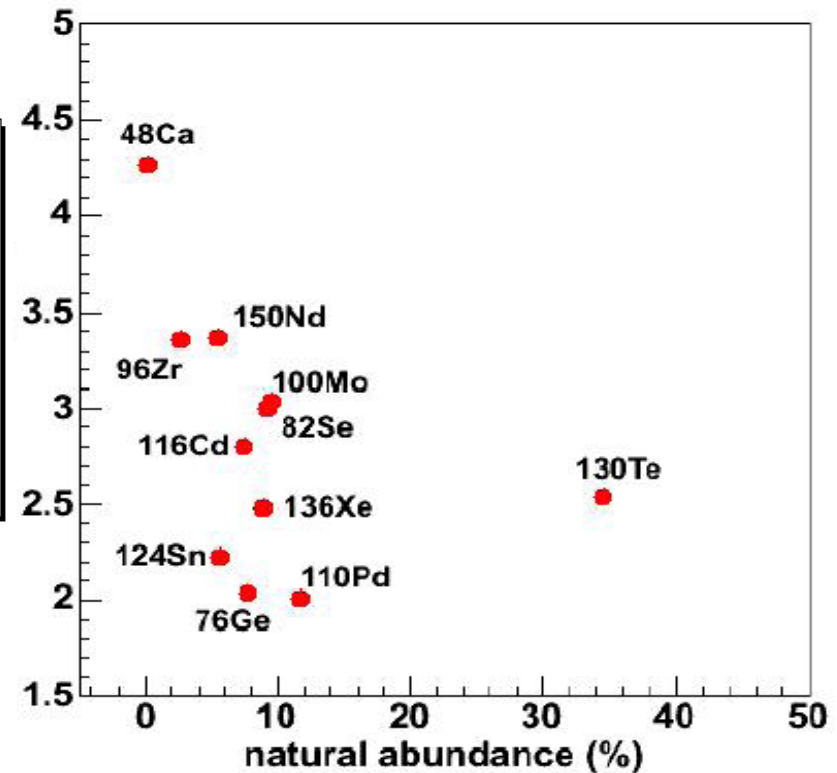
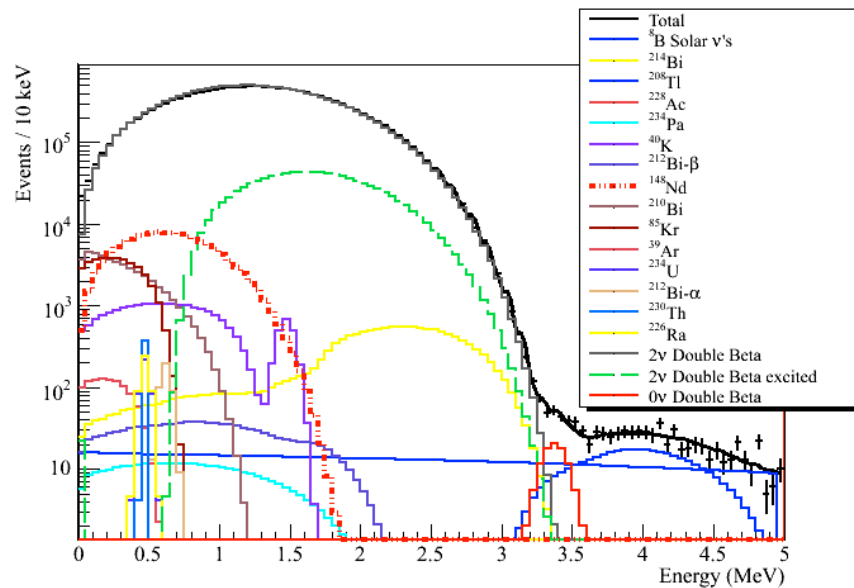
- Natural radioactivity
 - $T_{1/2}(^{238}\text{U}, ^{232}\text{Th}) \sim 10^{10}$ years
 - $T_{1/2}(0\nu\beta\beta) \sim 10^{25}$ years

- Standard-Model $2\nu\beta\beta$
 - Irreducible background



Dealing with backgrounds I

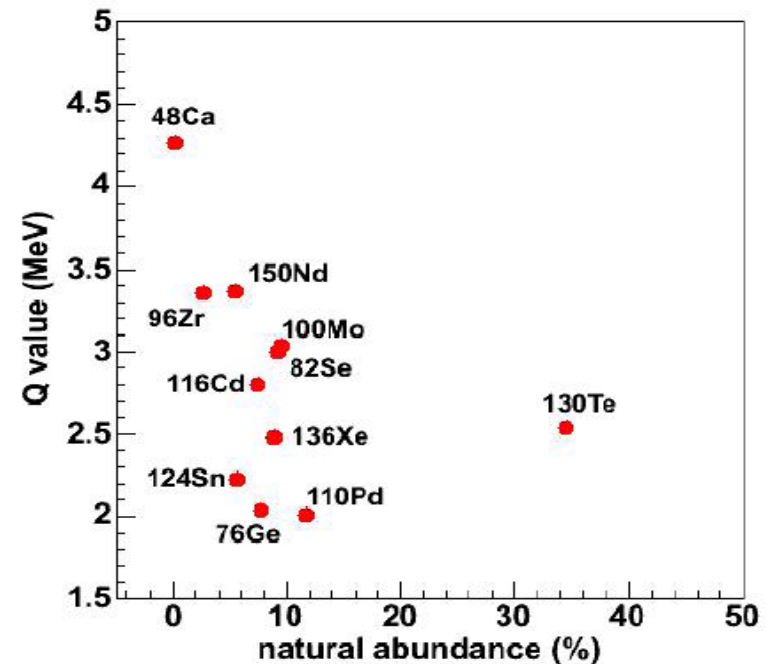
- Choose isotope with a high Q value
 - Energy above as much of the background as possible
- Also long $2\nu\beta\beta$ half-life



Maximising the decay rate

$$(T_{1/2})^{-1} = G_{0\nu} |M_{0\nu}|^2 |m_{\beta\beta}|^2 / m_e^2$$

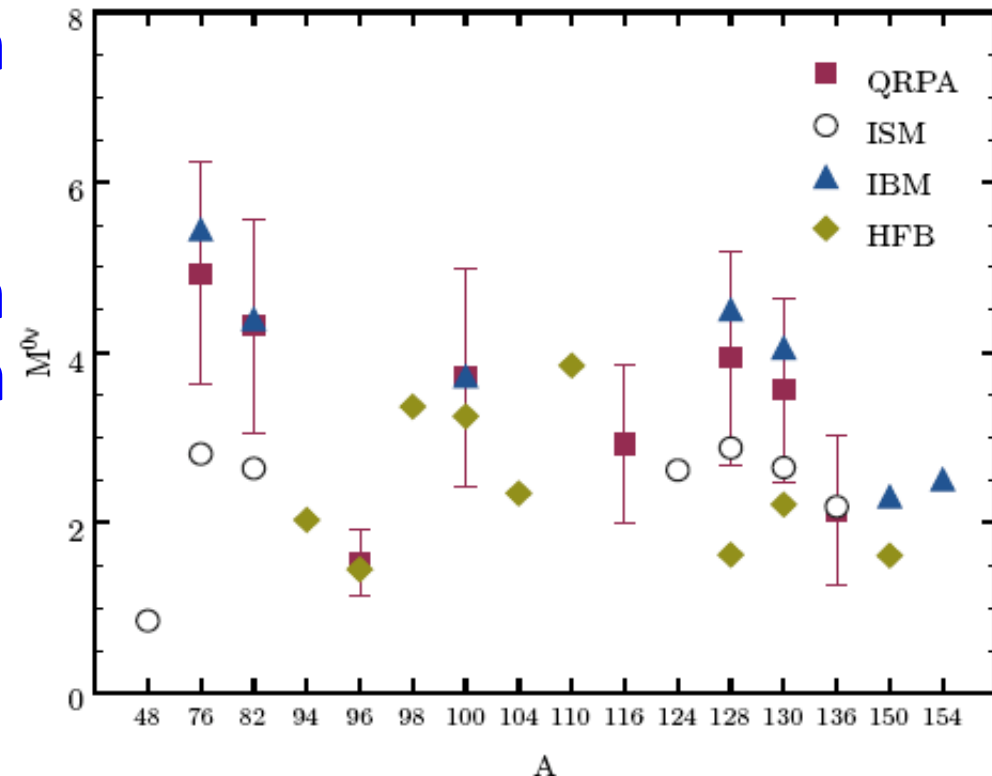
- The more isotope mass you can put in your experiment, the more decays there will be
 - Choose isotope with high natural abundance
 - Choose isotope that is “cheap” per unit mass
 - Consider isotope enrichment



Maximising the decay rate

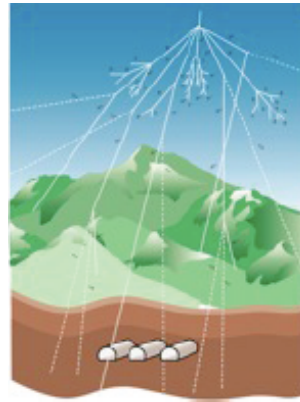
$$(T_{1/2})^{-1} = G_{0\nu} |M_{0\nu}|^2 |m_{\beta\beta}|^2 / m_e^2$$

- Choose isotope with large phase space
- Choose isotope with small uncertainty on nuclear matrix element



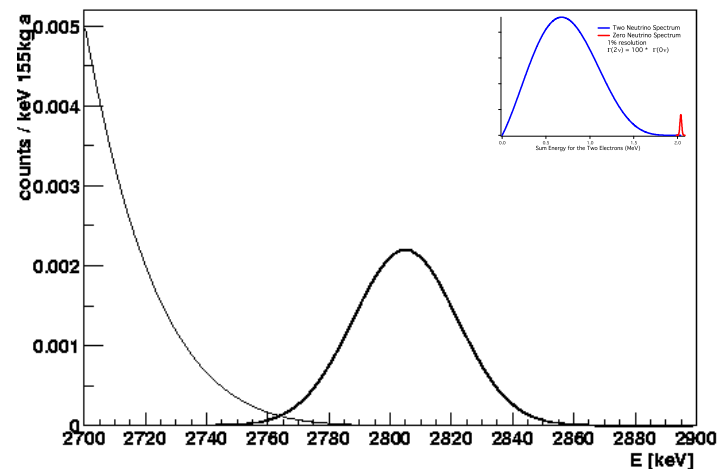
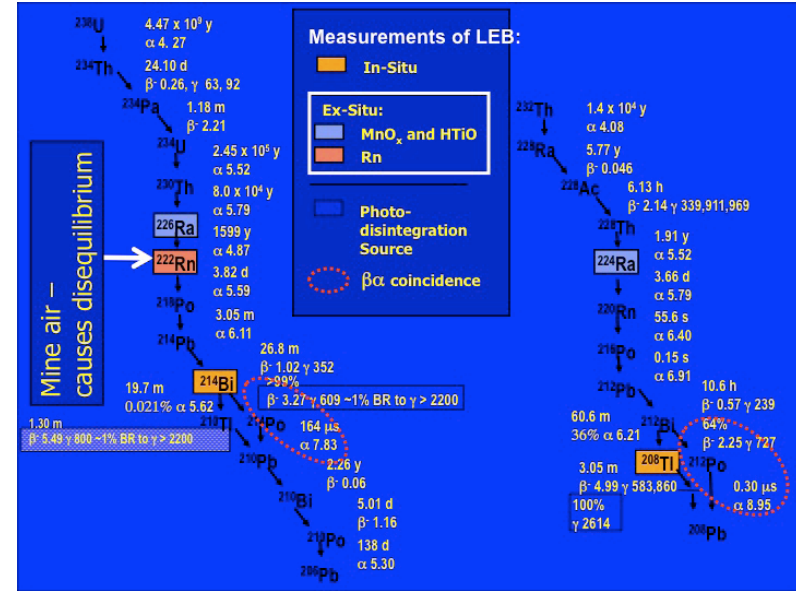
Dealing with backgrounds II

- Cosmic rays
 - Underground lab

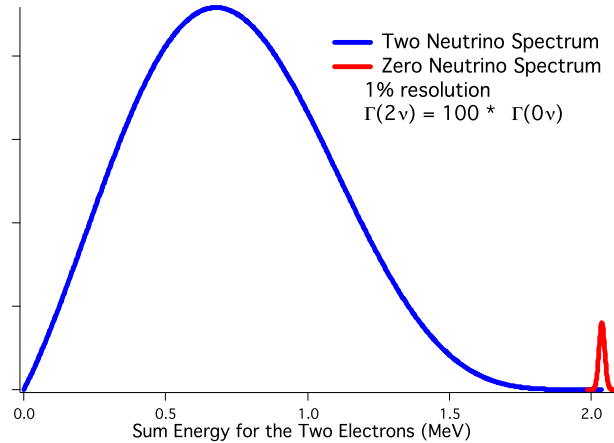


- Natural radioactivity
 - Radiopurity
 - Choice of detector materials
 - Cleanliness
 - Background identification

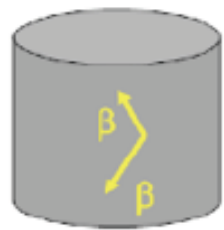
- Standard-Model $2\nu\beta\beta$
 - Irreducible background
 - Good energy resolution



Detector techniques



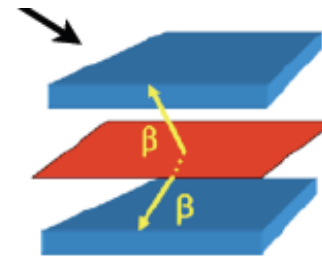
1. Two coincident electrons from the same vertex
2. $E_{e1} + E_{e2} = Q_{\beta\beta}$
3. Angular distribution



Source = detector
(calorimeter)

Good E resolution

Two approaches



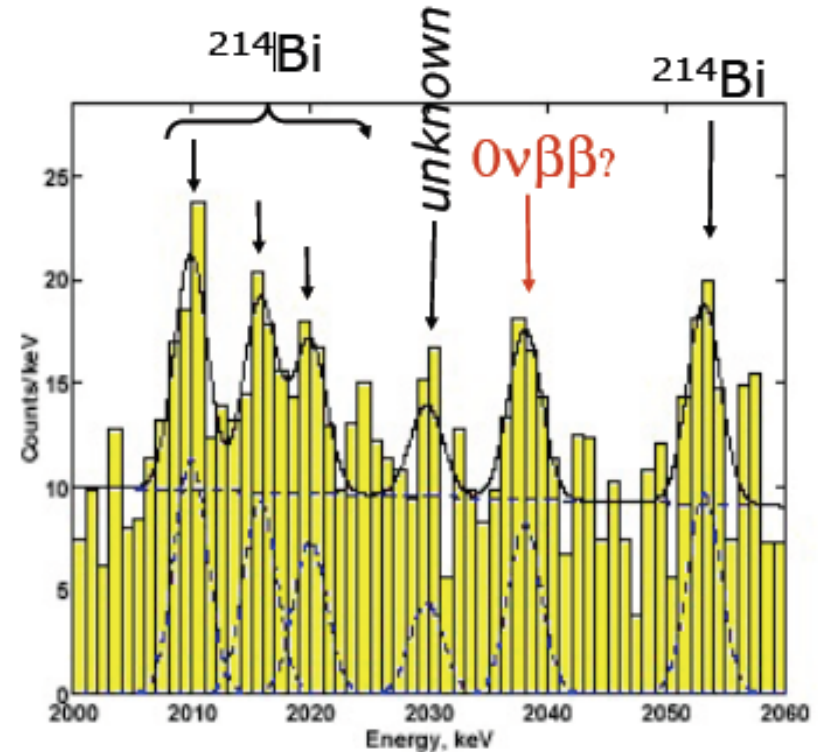
Source ≠ detector
(foil + tracking + calorimetry)

Background
by topology

- Scintillator calorimeters
- Semiconductors
- Bolometers (energy release via heating)
- Liquid/gas Xe TPC

A controversial claim

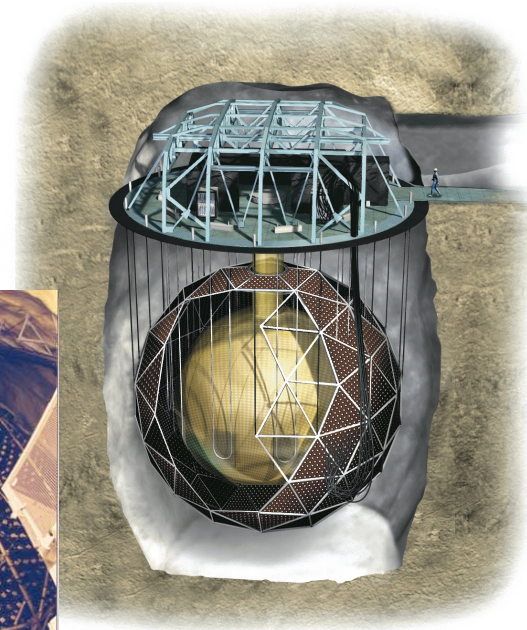
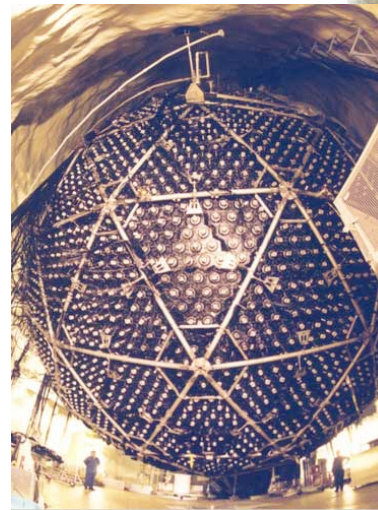
- Klapdor-Kleingrothaus et al., 2004
- High-Purity Ge detector
- Found unknown line at 2038 keV with 4.2σ significance
- $T_{1/2} \sim 10^{25}$ y
- Cannot be dismissed out of hand, but
 - Background underestimated
 - Unknown line in the same region
 - Problem with relative intensities of ^{214}Bi lines



$\langle m_\nu \rangle = 0.1 - 0.9 \text{ eV}$ (KKDC, 2004)

SNO+

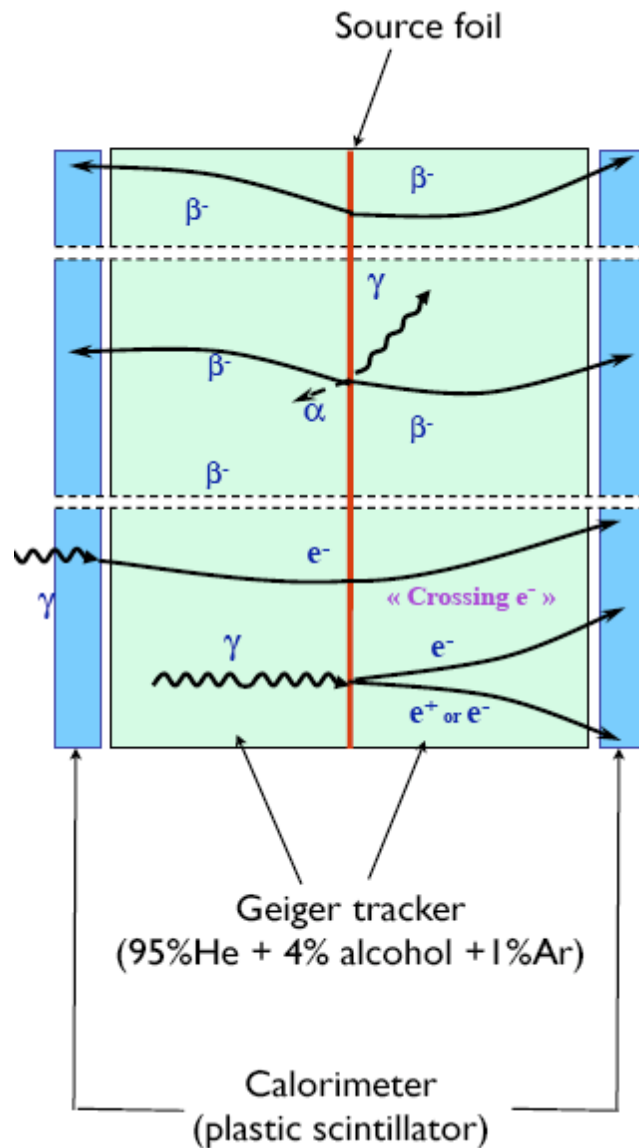
- Re-use old SNO (Sudbury Neutrino Observatory) vessel in Creighton nickel mine, Canada
- Candidate isotope ^{150}Nd dissolved in liquid scintillator
 - ^{96}Zr also interesting
- Also other ν sources, e.g.:
 - Low-E solar neutrinos: neutrino-matter interactions; solar composition
 - Reactors: neutrino oscillations
 - Supernova neutrinos
- Data collection to start in 2013



Art by Don Foley
Copyright (C) 2006, National Geographic Society
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NEMO3: topological signature



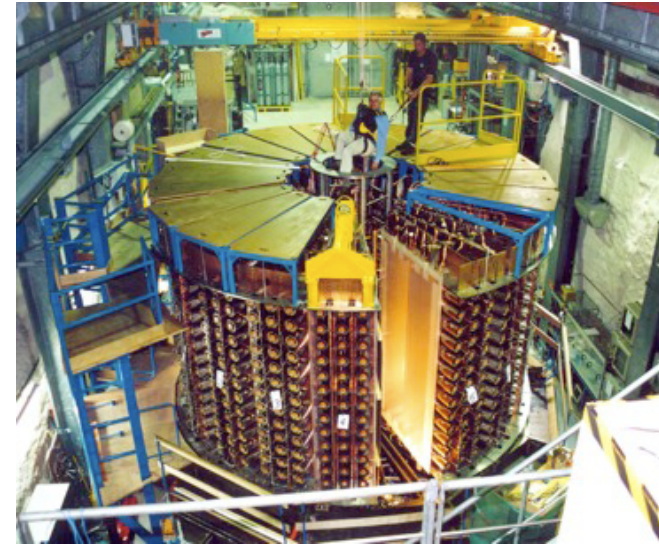
$\Delta t \sim 0$ ns Signal

$\Delta t \sim 0$ ns
Internal (e.g. ^{214}Bi)

$\Delta t \geq 3$ ns

External (e.g. ^{208}Tl , neutrons) Nemo3 (Modane lab)

$\Delta t \sim 0$ ns

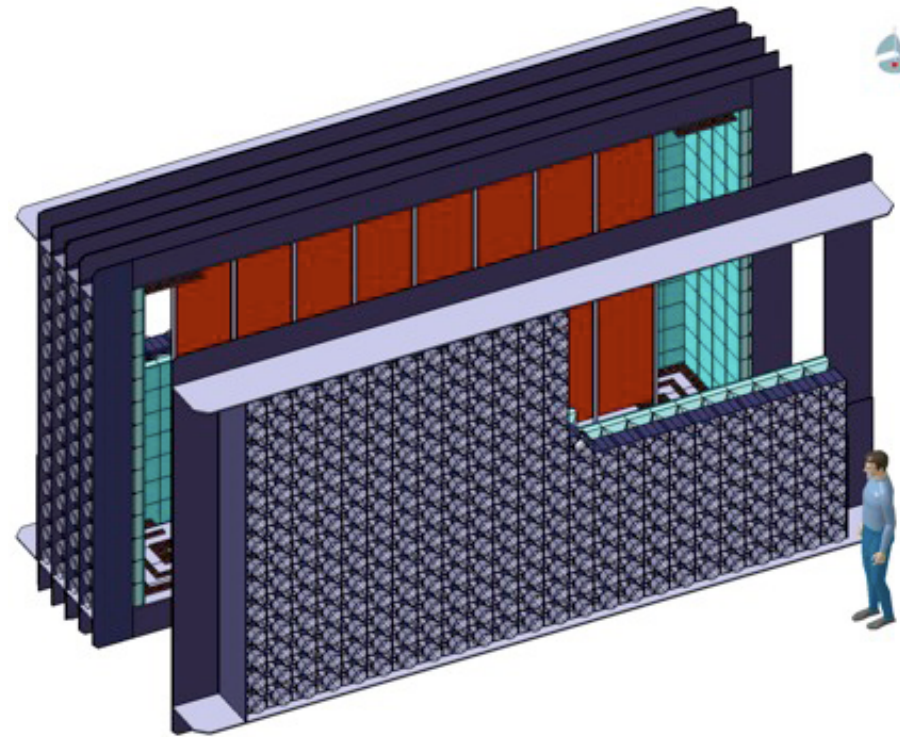


- Measure final-state observables:
 - Individual electron energies ($E_1 + E_2 = Q_{\beta\beta}$)
 - Electron trajectories and vertices
 - Time of flight
 - Angular distribution between electrons
- Background rejection through particle ID: e^- , e^+ , α , γ
- Has made many (world-leading) measurements of $2\nu\beta\beta$ half-lives
 - ^{100}Mo , ^{130}Te , ^{82}Se , ^{116}Cd , ^{150}Nd , ^{48}Ca , ^{96}Zr



SuperNEMO

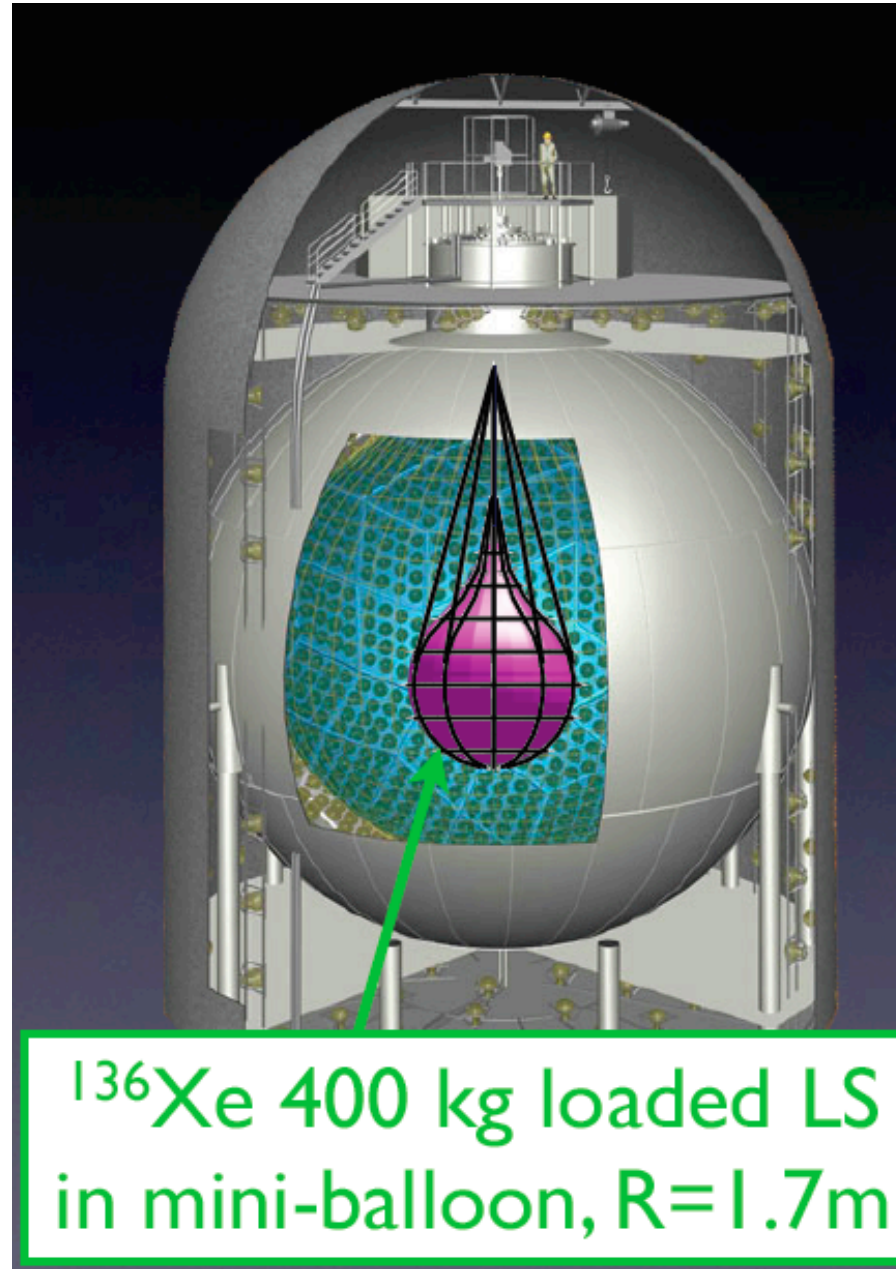
- Successor to NEMO3
 - Again, foils + tracking + calorimetry
- Planar and modular:
20 modules x 5 kg of enriched isotopes
 - 10x NEMO3
- Tracking:
Drift cells in Geiger mode
- Calorimeter:
Solid scintillators + PMTs
- Surrounded by water shielding
- Modane lab in France
- Next step: build one “demonstrator” module



KamLAND-Zen

Reuse old KamLAND detector, with ^{136}Xe as candidate isotope

Balloon to optimise light yield/E resolution vs. amount of ^{136}Xe

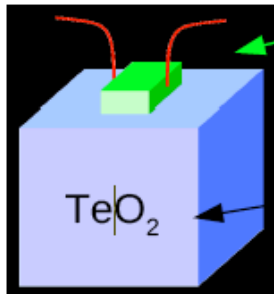


CUORE/CUORICINO: the bolometric way

Temperature sensor: $\Delta T \rightarrow \Delta V$

NTD thermistor

$R = R_0 \exp(T/T_0)^\gamma \rightarrow$ high sensitivity



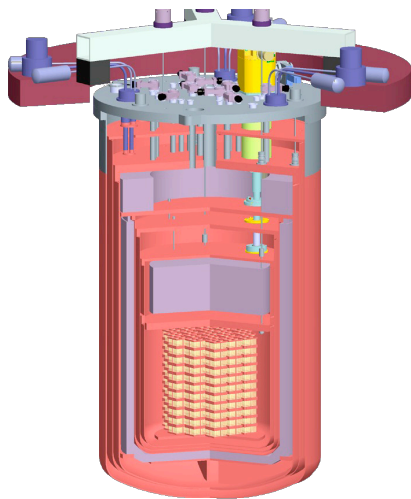
Absorber: $E \rightarrow \Delta T \sim E/C(T)$

TeO₂ crystals

Low heat capacity

High radio-purity

Large size crystals available

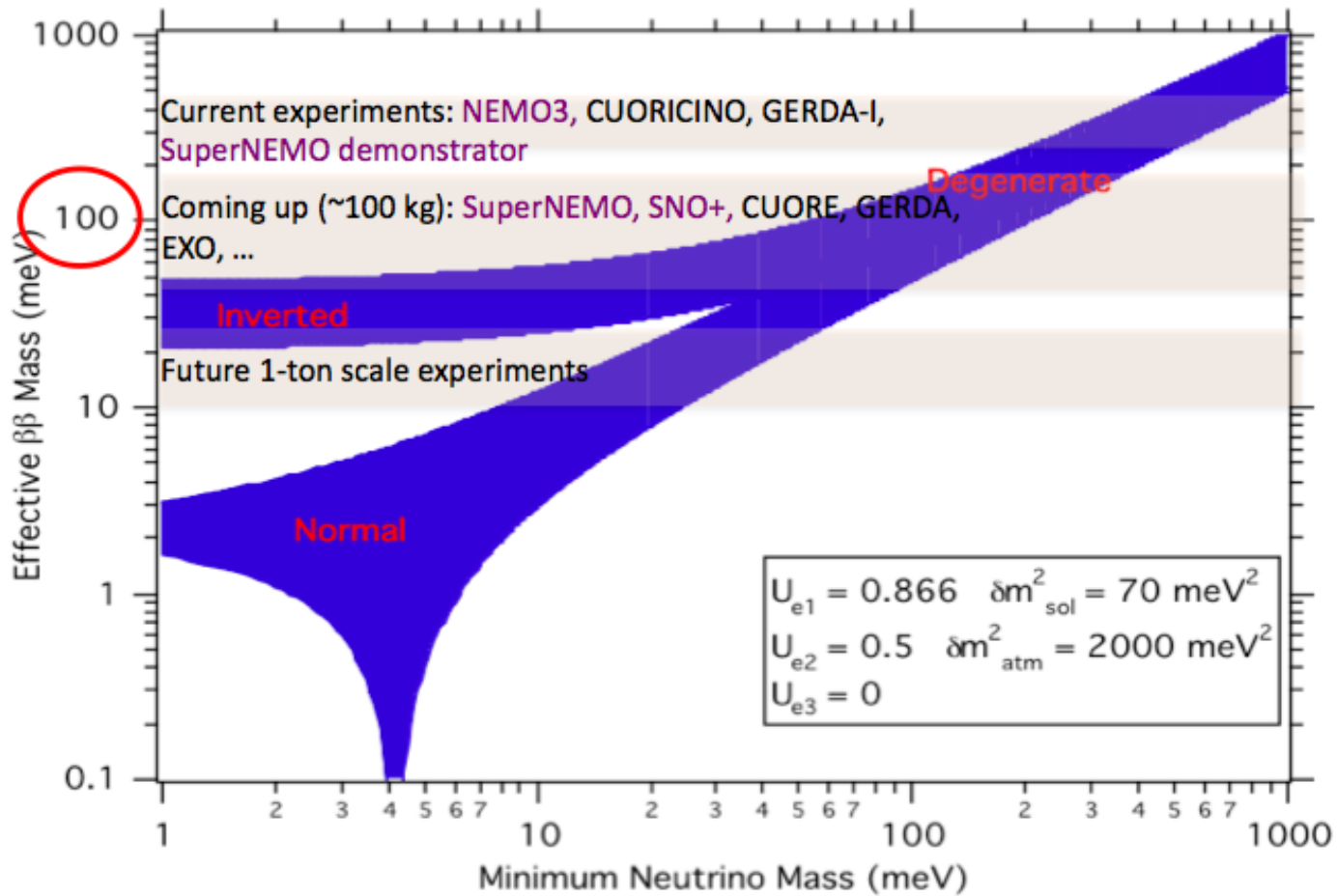


Tower of TeO₂ crystals in cryostat

Plan for ~ 200 kg of ¹³⁰Te

Many players lined up

Probing $m_{\beta\beta}$
down to
 $O(100 \text{ meV})$



Finally: other limits on neutrino mass

Measurement	Origin of Limit	Limit
Tritium Decay	$m_{\nu_e}^2$	$< 2.0 \text{ eV}$
π^+ Decay	$m_{\nu_\mu}^2$	$< 0.17 \text{ MeV}$
τ Decay	$m_{\nu_\tau}^2$	$< 18.2 \text{ MeV}$
SN1987A Time-of-Flight	$m_{(\nu_e \rightarrow \nu)}^2$	$< 5.7 \text{ eV}$
Terrestrial Time-of-Flight	$m_{(\nu_e \rightarrow \nu)}^2$	$< 50 \text{ MeV}$
Cosmology	$\sum_i m_{\nu_i}$	$< 0.66 \text{ eV}$

From PhD thesis by B. Still, 2009

Recap

- Neutrinoless double beta decay, if it exists, would violate lepton-number conservation by 2
- Half-lives upwards of $\sim 10^{18}$ years
- Neutrinoless double beta decay is the only experimental technique known to be able to reveal whether the neutrino is Majorana
- Many experiments in the pipeline

Back-ups