

P and CP violation of the weak decays



The Nobel Prize in Physics 1957

Chen Ning Yang, Tsung-Dao Lee

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Chen Ning Yang

Tsung-Dao Lee



Chen Ning Yang



Tsung-Dao (T.D.) Lee

The Nobel Prize in Physics 1957 was awarded jointly to Chen Ning Yang and Tsung-Dao (T.D.) Lee *"for their penetrating investigation of the so-called parity laws which has led to important discoveries regarding the elementary particles"*

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Reminder about parity

The experiment of Wu



So: $n \rightarrow p + e^- + \bar{\nu}_e$

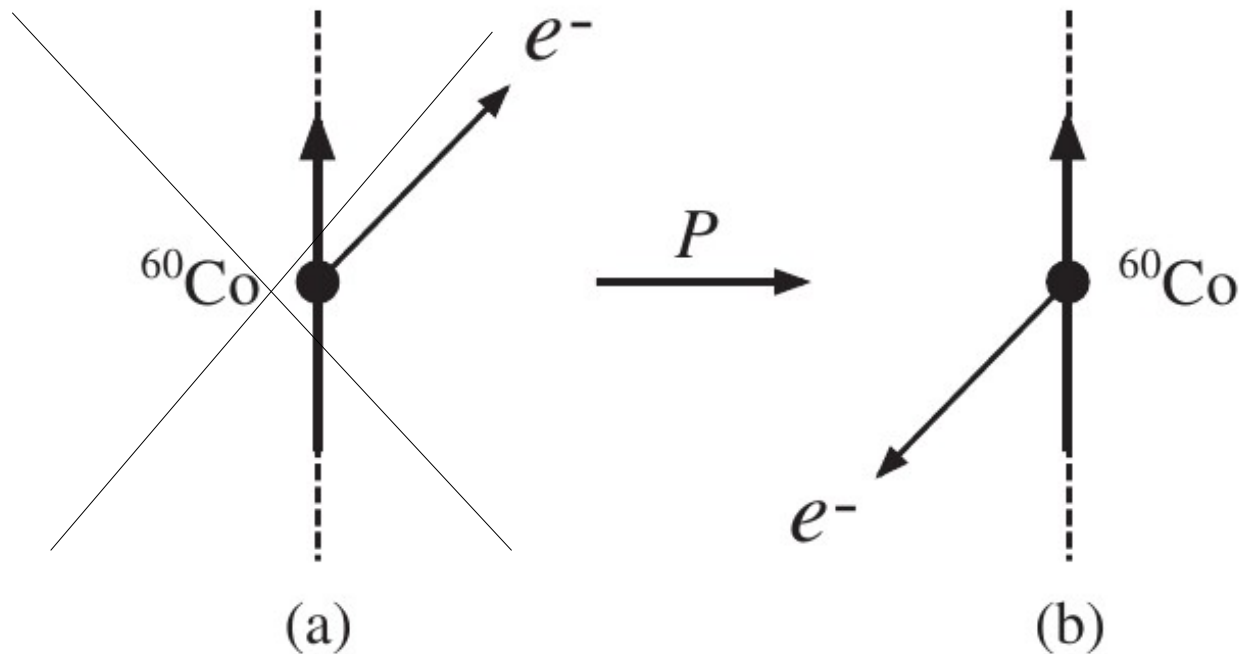


Figure 10.1 Effect of a parity transformation on ${}^{60}\text{Co}$ decay (10.1). The short thick arrows indicate the direction of the spin of the ${}^{60}\text{Co}$ nucleus, while the long arrows show the direction of the electron's momentum.

Handedness for massless particles

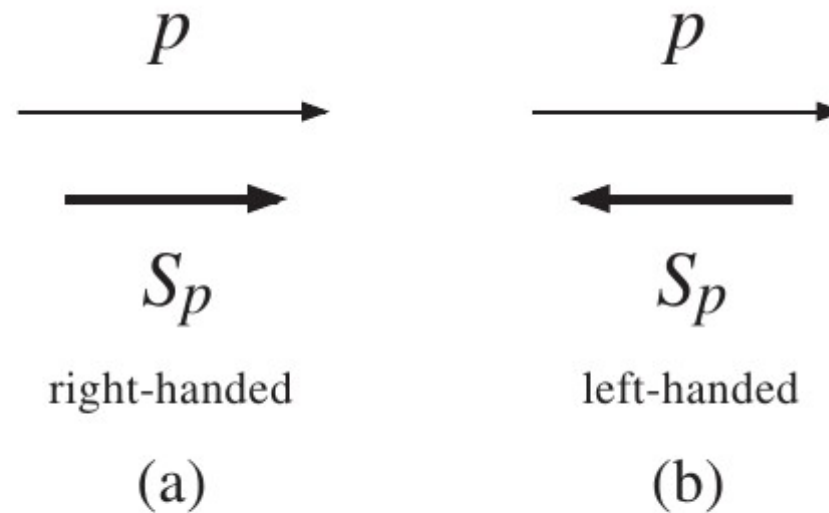


Figure 10.3 Helicity states of a spin- $\frac{1}{2}$ particle. The long thin arrows represent the momenta of the particles and the short thick arrows their spins.

Why only the left is observed

The parent nucleus has the z-component of the spin $S_z = +5$, the (excited) daughter nucleus $S_z = +4$.

So spin of neutrino and electron must point up!

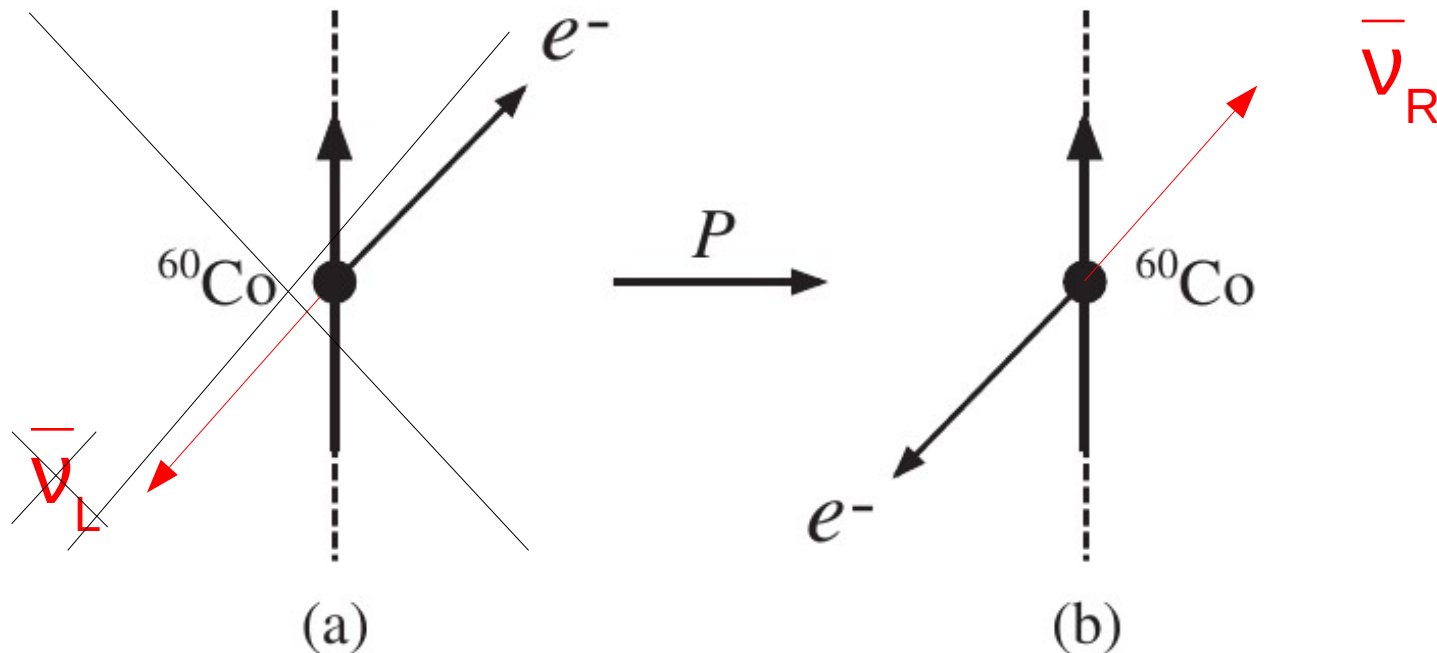


Figure 10.1 Effect of a parity transformation on ^{60}Co decay (10.1). The short thick arrows indicate the direction of the spin of the ^{60}Co nucleus, while the long arrows show the direction of the electron's momentum.

For neutrinos

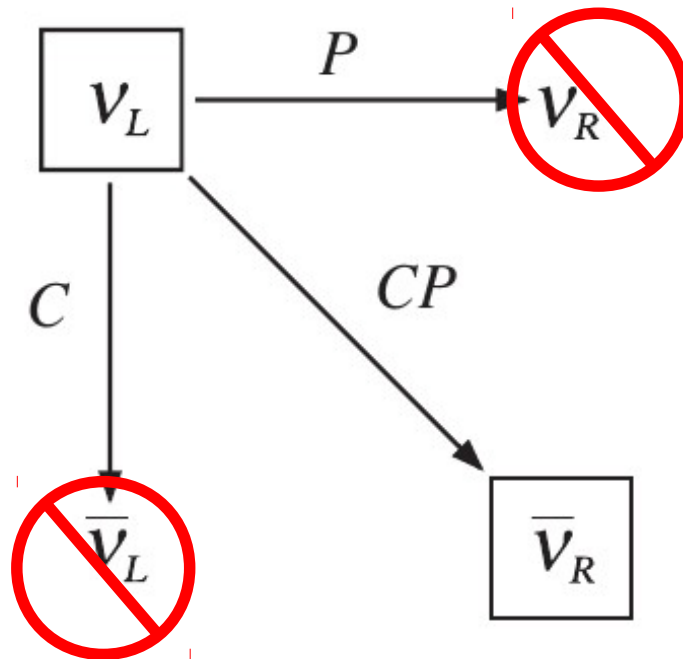


Figure 10.4 Effect of C , P and CP transformations on a left-handed neutrino ν_L . Only the states shown in boxes have been observed in nature.

Left-handed particles & Right-handed anti-particles

Three generations of matter (fermions)

	I	II	III		
mass →	2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²	0	? GeV/c ²
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin →	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
name →	u up	c charm	t top	γ photon	H Higgs boson
Quarks	4.8 MeV/c ²	104 MeV/c ²	4.2 GeV/c ²	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	d down	s strange	b bottom	g gluon	
Leptons	<2.2 eV/c ²	<0.17 MeV/c ²	<15.5 MeV/c ²	91.2 GeV/c ²	
	0	0	0	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z⁰ Z boson	
	0.511 MeV/c ²	105.7 MeV/c ²	1.777 GeV/c ²	80.4 GeV/c ²	
	-1	-1	-1	±1	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	e electron	μ muon	τ tau	W[±] W boson	

Gauge bosons

Right-handed particles & Left-handed anti-particles

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	-1	-1	-1	0	
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	e electron	μ muon	τ tau	Z⁰ Z boson	

Gauge bosons

The Z⁰ couples different to right handed particles. See e.g. Leif's notes.
(This is because the Z⁰ is a mix between B which couples and W⁰ which does not couple to right handed particles)

Handedness for massive particles

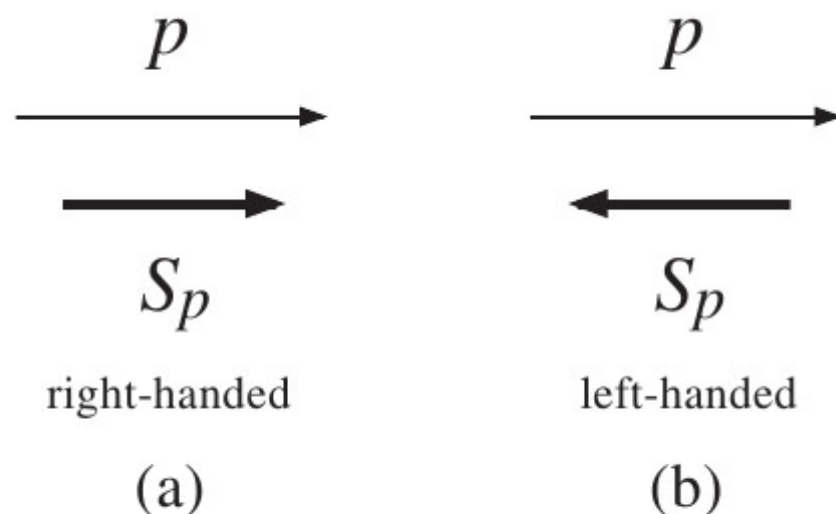


Figure 10.3 Helicity states of a spin- $\frac{1}{2}$ particle. The long thin arrows represent the momenta of the particles and the short thick arrows their spins.

Handedness is a well defined property for massless particles (coincides with helicity).

For massive particles the handedness can be flipped by a boost!

For a massive particle the forbidden helicity states are only suppressed by:

$$\left(1 - \frac{v}{c}\right) \approx \frac{m^2}{2E^2}, \quad (10.13)$$

This is also why the pion prefers to decay to a muon!

- Consider:
 - Decay 1: $\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$
 - Decay 2: $\pi^- \rightarrow e^- + \bar{\nu}_e$
- As pion has spin 0, the e^-/μ^- has to have opposite spin of the anti-neutrino.
- In the CM:

Momentum: $\langle \text{----- } e^- \quad \pi^- \quad \bar{\nu}_e \text{ -----} \rangle$

Spin (only $\bar{\nu}_e$): $\langle \text{----- } e^- \quad \pi^- \quad \bar{\nu}_e \text{ -----} \rangle$
- But the W^- coupling to e^-_R is “forbidden”/suppressed!

$$\frac{\Gamma(\pi^+ \rightarrow e^+ \nu_e)}{\Gamma(\pi^+ \rightarrow \mu^+ \nu_\mu)} = (1.230 \pm 0.004) \times 10^{-4} \quad (10.15)$$



The Nobel Prize in Physics 1980

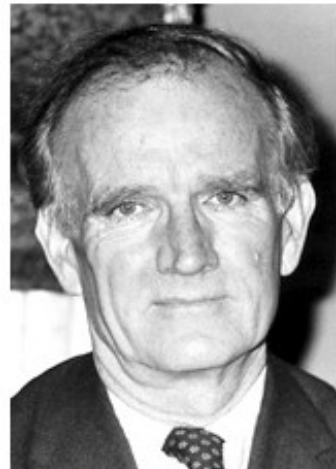
James Cronin, Val Fitch

The Nobel Prize in Physics 1980

Nobel Prize Award Ceremony

James Cronin

Val Fitch



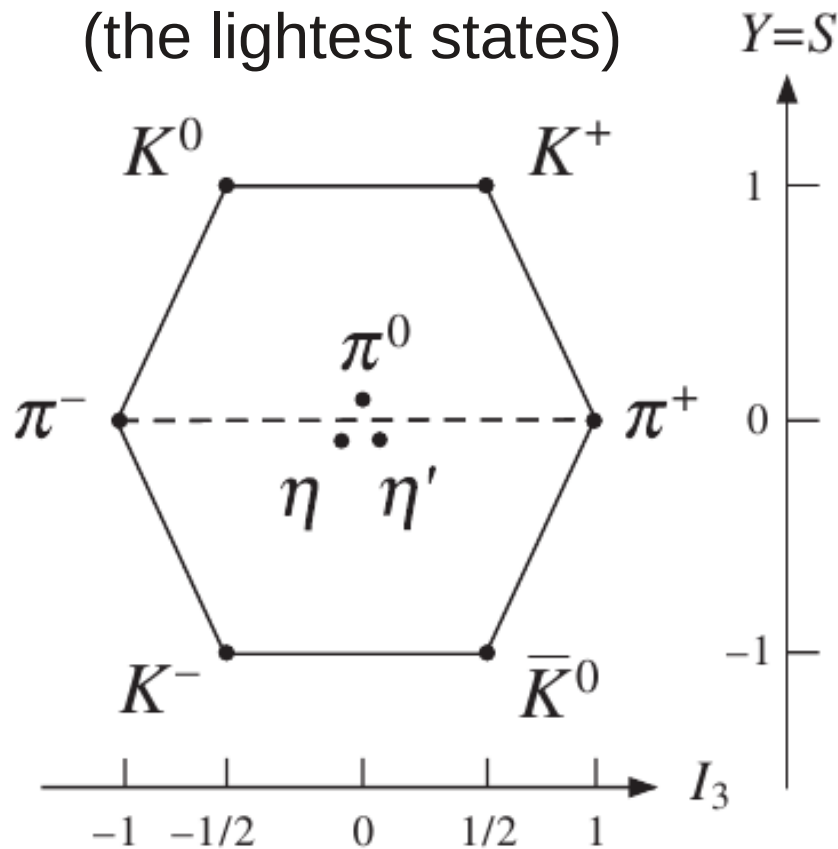
James Watson Cronin Val Logsdon Fitch

The Nobel Prize in Physics 1980 was awarded jointly to James Watson Cronin and Val Logsdon Fitch *"for the discovery of violations of fundamental symmetry principles in the decay of neutral K-mesons"*

Photos: Copyright © The Nobel Foundation

C and P for the light mesons

The $L=0$, $S=0$, $J=0$ states
(the lightest states)



- Rule for P

$$P_M = P_a P_{\bar{b}} (-1)^L = (-1)^{L+1}.$$

(5.39a)

So these states all have $P = -1$

- There is no general rule for C
- $C = +1$ for pions
- $C = -1$ for K^0 and \bar{K}^0

$K^0 \leftrightarrow \bar{K}^0$ oscillations

$$K^0(498) = d\bar{s} \quad \text{and} \quad \bar{K}^0(498) = s\bar{d}, \quad (10.17)$$

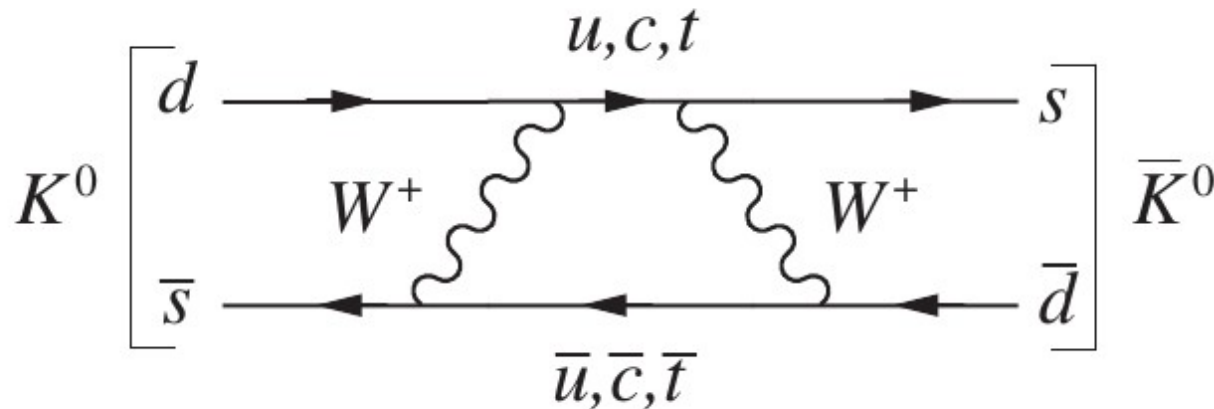


Figure 10.8 Example of a process that can convert a $K^0 \equiv \bar{s}d$ state into a $\bar{K}^0 \equiv s\bar{d}$ state. The labels on the internal lines identify the quark or antiquark going from left to right, i.e. not necessarily in the direction of the arrow.

Question?

$$|K_1^0, \mathbf{p} = \mathbf{0}\rangle \equiv \frac{1}{\sqrt{2}} [|K^0, \mathbf{p} = \mathbf{0}\rangle + |\bar{K}^0, \mathbf{p} = \mathbf{0}\rangle] \quad (10.21a)$$

and

$$|K_2^0, \mathbf{p} = \mathbf{0}\rangle \equiv \frac{1}{\sqrt{2}} [|K^0, \mathbf{p} = \mathbf{0}\rangle - |\bar{K}^0, \mathbf{p} = \mathbf{0}\rangle], \quad (10.21b)$$

- What is the CP of each state?
- Which state can decay to 2 pions and which to 3 pions?

$$CP \left| K_1^0, \mathbf{p} = \mathbf{0} \right\rangle = \left| K_1^0, \mathbf{p} = \mathbf{0} \right\rangle \quad (10.22a)$$

$$CP \left| K_2^0, \mathbf{p} = \mathbf{0} \right\rangle = - \left| K_2^0, \mathbf{p} = \mathbf{0} \right\rangle, \quad (10.22b)$$

$$K_1^0 \rightarrow \pi^+ \pi^-, \pi^0 \pi^0 \quad (10.23a)$$

$$K_2^0 \rightarrow \pi^+ \pi^- \pi^0, \pi^0 \pi^0 \pi^0 \quad (10.23b)$$

Argument

- Initially $J = 0$ and as pions have $J=0$ then $L=0$
- $P=P_{\pi}^2 (-1)^L = 1$ (-1 for 3 pions)
- $C=(+1)^2 (-1)^L = 1$ (also 1 for 3 pions)

The real states

- Experimentally one observes

$$K_S^0 \rightarrow \pi^0 + \pi^0, (B=0.31) \quad \text{and} \quad K_S^0 \rightarrow \pi^+ + \pi^-, (B=0.69), \quad (10.28)$$

with lifetime $\tau \sim 9 \times 10^{-11} \text{s}$, and

$$K_L^0 \rightarrow \pi^0 + \pi^0 + \pi^0 (B=0.20), \quad K_L^0 \rightarrow \pi^+ + \pi^- + \pi^0 (B=0.13) \quad (10.29)$$

with lifetime $\tau \sim 5 \times 10^{-8} \text{s}$, and so we guess

$$K_S^0 = K_1^0 \quad \text{and} \quad K_L^0 = K_2^0. \quad (10.31)$$

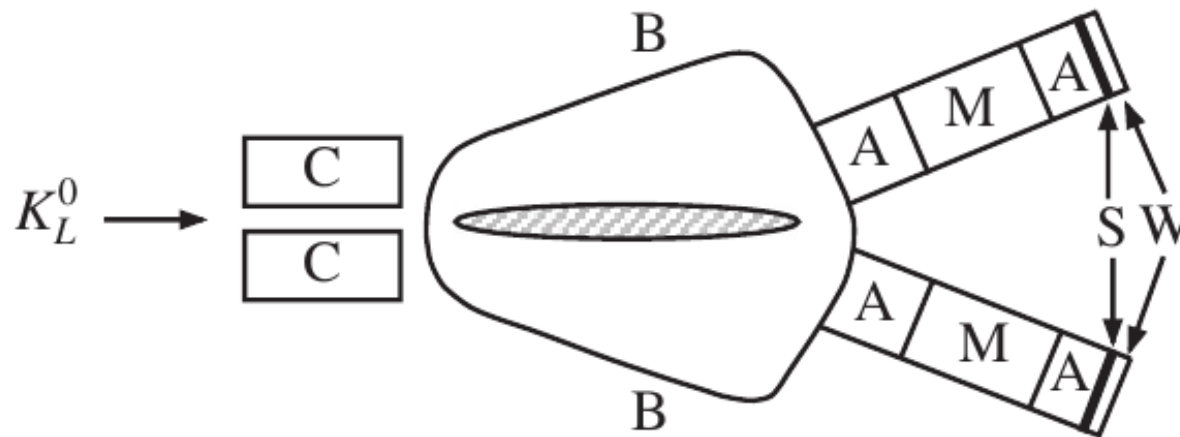


Figure 10.10 Schematic diagram of the apparatus used in the discovery of CP violation. The K_L^0 beam entered a helium-filled bag B through a lead collimator C, and those CP -violating decays that occurred in the shaded region were detected by the symmetrically spaced spectrometers. These each contained a pair of spark chambers A separated by a magnet M, followed by a scintillation counter S and a water Čerenkov counter W. (Reprinted Figure 1 with permission from J. H. Christenson *et al.*, *Phys. Rev. Lett.*, **13**, 138. Copyright 1964 American Physical Society.)

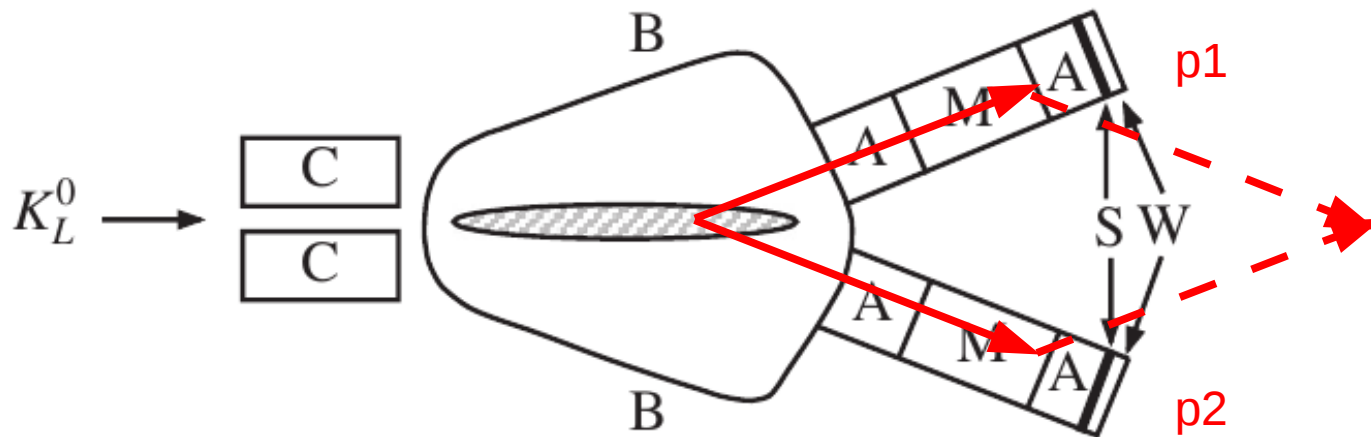


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Clear signature for CP violating 2 pion decay :
 + and – track, $M_{\text{inv}}^2 = (p_1 + p_2)^2 \sim MK^2$,
 sum of 3 momenta parallel to beam

Result

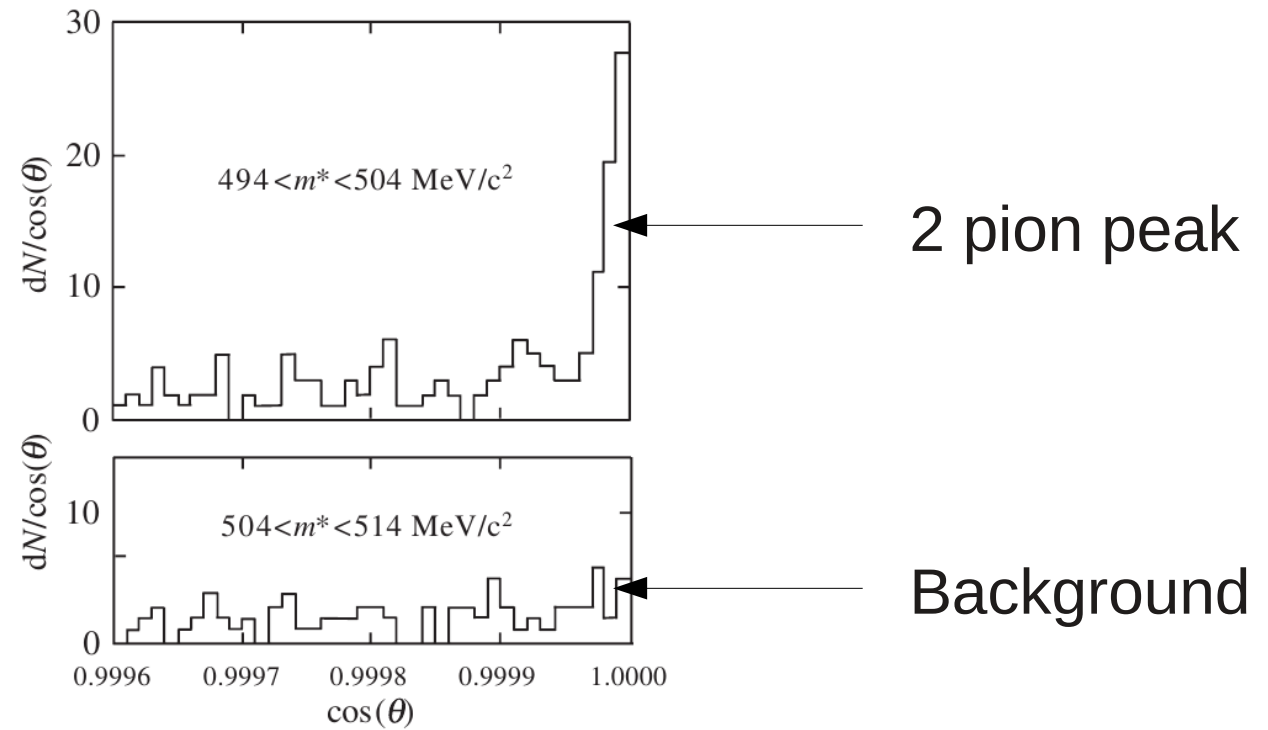


Figure 10.11 Angular distribution of the $\pi^+\pi^-$ pairs detected using the apparatus of Figure 10.10, where θ is the angle between the line-of-flight of the centre-of-mass of the pair and the initial beam direction. Results are shown for an invariant mass range including the K_L^0 mass ($498 \text{ MeV}/c^2$) and for a neighbouring mass range. (Reprinted Figure 1 with permission from J. H. Christenson *et al.*, *Phys. Rev. Lett.*, **13**, 138. Copyright 1964 American Physical Society.)

CP is not good symmetry for all weak decays

$$|K_S^0, \mathbf{p}=\mathbf{0}\rangle = \frac{1}{(1+|\varepsilon|^2)^{1/2}} [|K_1^0, \mathbf{p}=\mathbf{0}\rangle - \varepsilon |K_2^0, \mathbf{p}=\mathbf{0}\rangle] \quad (10.33a)$$

$$|K_L^0, \mathbf{p}=\mathbf{0}\rangle = \frac{1}{(1+|\varepsilon|^2)^{1/2}} [\varepsilon |K_1^0, \mathbf{p}=\mathbf{0}\rangle + |K_2^0, \mathbf{p}=\mathbf{0}\rangle], \quad (10.33b)$$

- So in this case for there is a bit more K_0 than \bar{K}_0 in K_{0L}

$$K_{0L} \sim (1+\varepsilon)K_0 - (1-\varepsilon)\bar{K}_0$$

$$K_{0S} \sim (1-\varepsilon)K_0 + (1+\varepsilon)\bar{K}_0$$

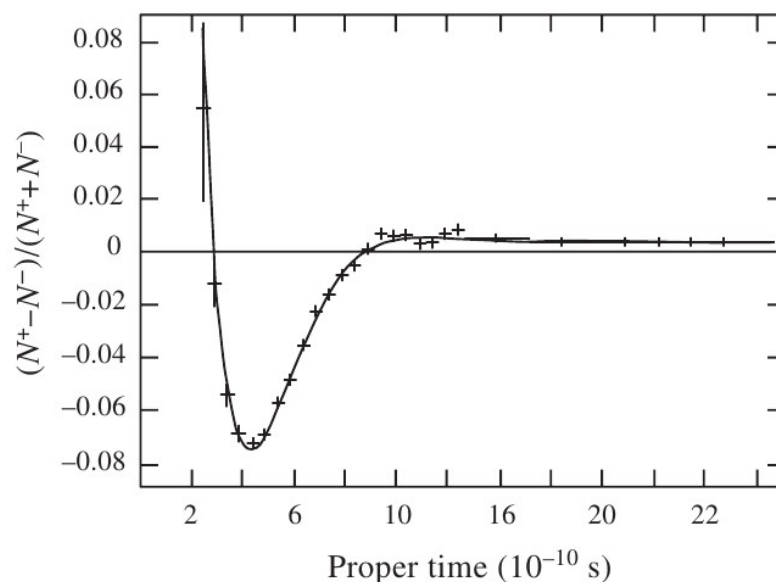
Evidence for mixing

To understand this, we note that the K^0 and \bar{K}^0 can decay by the semileptonic reactions

$$K^0 \rightarrow \pi^- + e^+ + \nu_e \quad \text{and} \quad \bar{K}^0 \rightarrow \pi^+ + e^- + \bar{\nu}_e, \quad (10.35)$$

whereas the corresponding reactions

$$K^0 \rightarrow \pi^+ + e^- + \bar{\nu}_e \quad \text{and} \quad \bar{K}^0 \rightarrow \pi^- + e^+ + \nu_e \quad (10.36)$$



$$|\varepsilon| = (2.232 \pm 0.007) \times 10^{-3}.$$

(10.41)

Figure 10.12 The charge asymmetry observed for $K^0 \rightarrow \pi^- e^+ \nu_e$ and $\bar{K}^0 \rightarrow \pi^+ e^- \bar{\nu}_e$ as a function of proper time, for a beam that is initially predominantly K^0 . (Reprinted from *Physics Letters B*, **52**, 113 Gjesdal, S., *et al.* Copyright 1974, with permission from Elsevier.)

So nature is not completely symmetric

- We can distinguish between left and right because of the P violation
- We can distinguish between matter and anti-matter because of the CP (C) violation

Also observed in other systems

- **LHCb experiment observes new matter-antimatter difference**
- Geneva, 24 April 2013. The LHCb collaboration at CERN today submitted a paper to Physical Review Letters on the first observation of matter-antimatter asymmetry in the decays of the particle known as the B0s. It is only the fourth sub-atomic particle known to exhibit such behaviour.
- Now the LHCb experiment has observed a preference for matter over antimatter known as CP-violation in the decay of neutral B0s particles.....