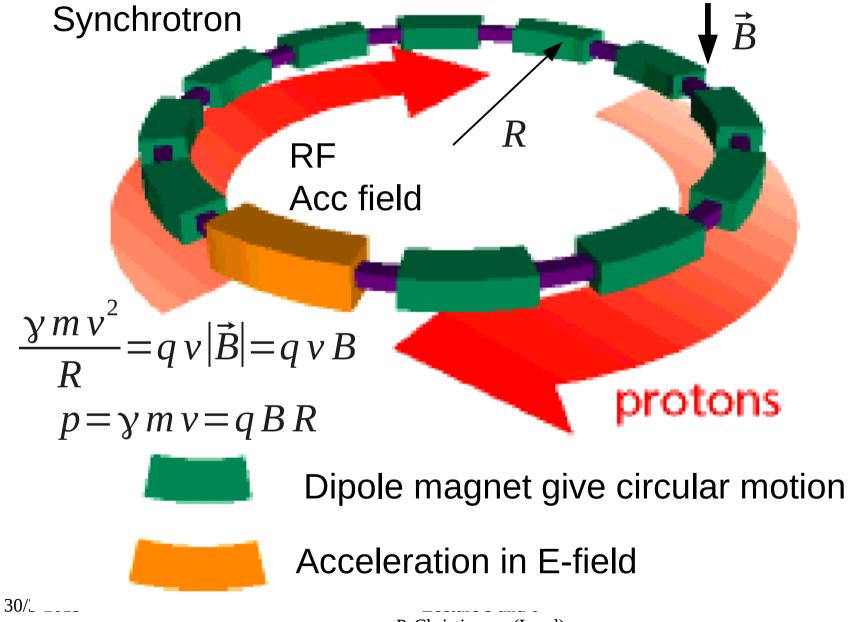
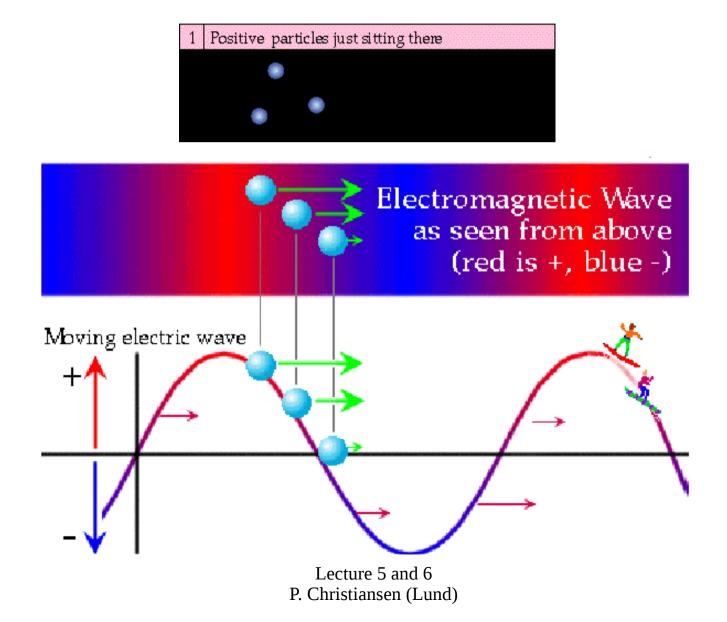
Lectures on accelerator physics

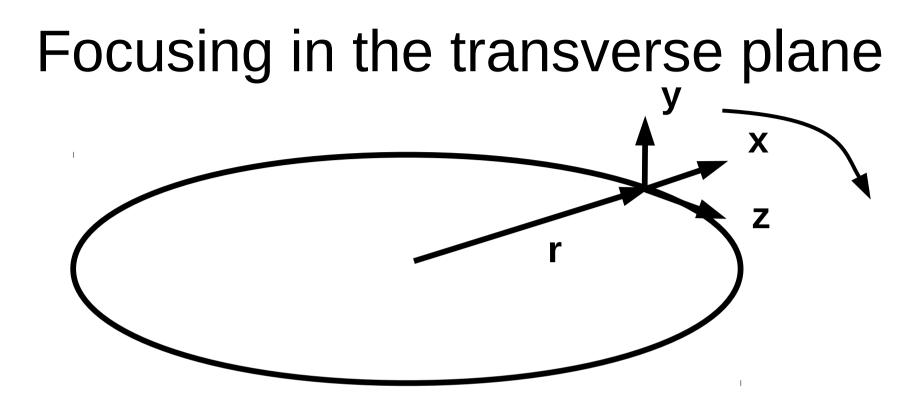
- Lecture 5 and 6: Advanced topics
 - Transverse motion, strong focusing, and LHC
- Material borrowed from
 - Lecture by Anders Oskarsson
 - Lecture by Eric Torrence (University of Oregon)
 - LHC lectures by Danillo Vranic (GSI)
- Weak focusing follows "Principles of Charged Particle Acceleration" by Stanley Humpries Jr. Chapter 7.

Towards the Synchrotron



The particles are "surfing" the acceleration wave

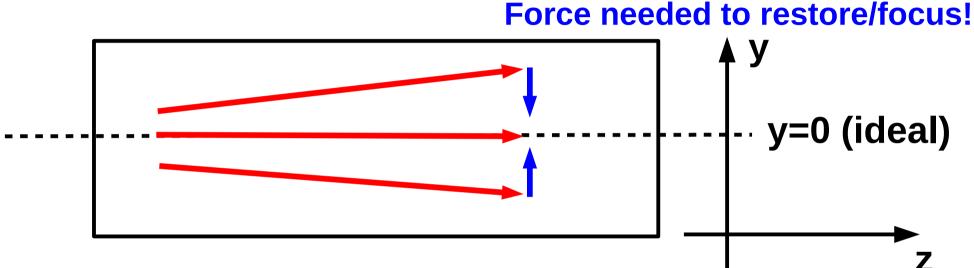




- Assume $(v_x, v_y, v_z) \sim (0, 0, v)$ and $v \sim c = constant!$
 - Very good assumption!
- $z=vt \rightarrow t=z/v (\sim z/c)$
 - d/dt ~ v d/dz (~c d/dz)

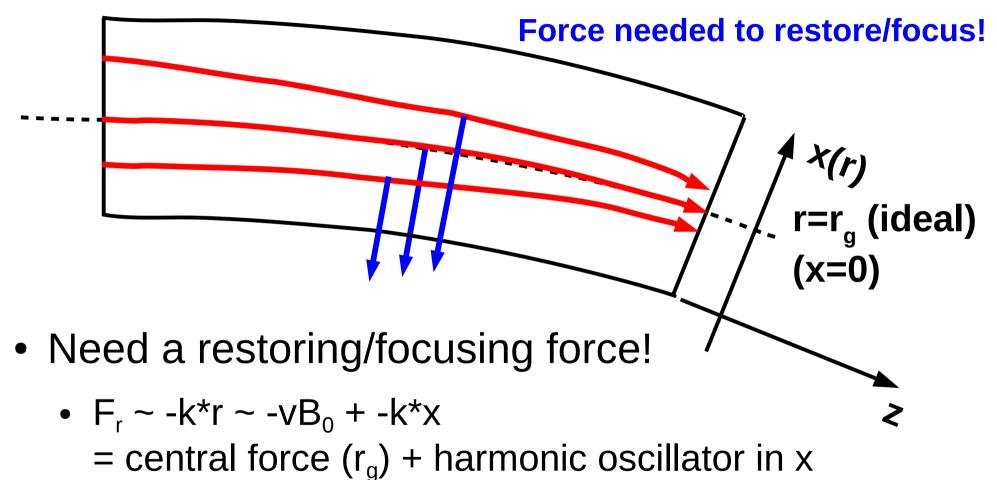
Lecture 5 and 6 P. Christiansen (Lund)

Transverse focusing in y direction



- Need a restoring/focusing force!
 - $F_y \sim -k^*y$
- Harmonic oscillator (like string)

Transverse focusing in x direction



Let us first solve harmonic equation (ignoring magnet realities!)

$$\gamma m \frac{d^2 y}{dt^2} = \gamma m v^2 \frac{d^2 y}{dz^2} = -ky$$

$$y(z) = y_0 \cos\left(\frac{2\pi}{\lambda}z + \varphi\right),$$

where

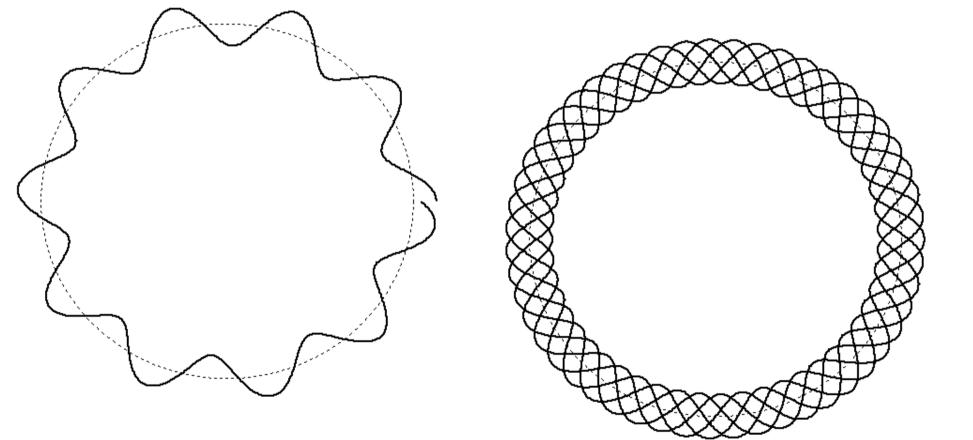
$$\lambda = 2\pi \sqrt{\frac{\gamma m v^2}{k}}.$$

 Note that the wavelength does not depend on the amplitude y₀. There is only one wavelength for all amplitudes!

Tune interlude

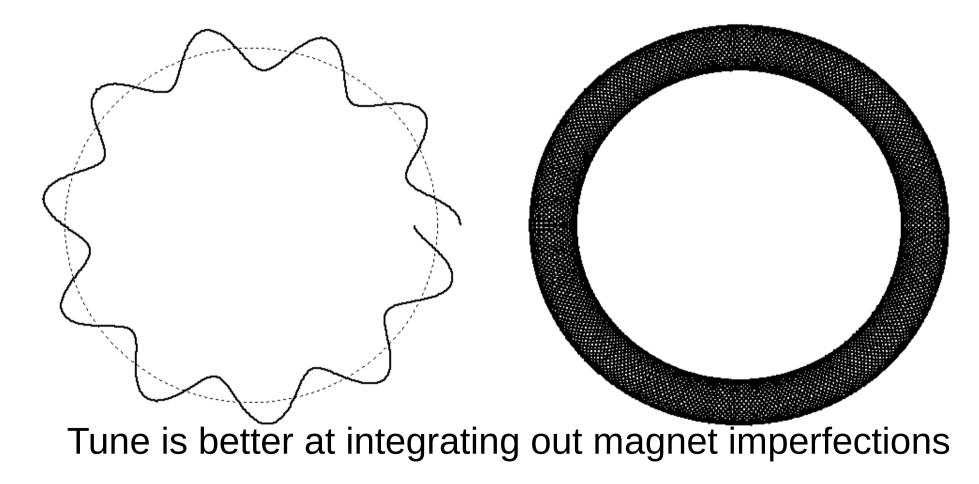
- One defines
 - Q (v[nu]) = C/ λ , where C=2 πr_g is the circumference of the synchrotron ring
- Q is the number of transverse (betatron) oscillations per turn
- It is different for x and y directions
- Very important for beam stability!

Bad harmonic tune (Q=10.2) 1 turn 100 turns

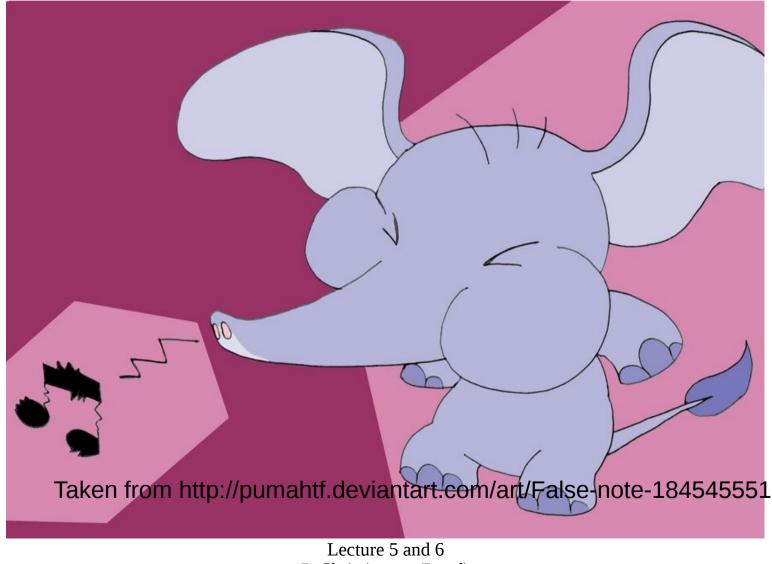


Problem: tune does not integrate out magnet imperfections

Better (less harmonic) tune (Q=10.48) 1 turn 100 turns



Lesson: bad musicians makes great accelerator physicists



LHC TUNES

AT 7TeV HORIZONTAL TUNE: $Q_x = 64.31$ $VERTICAL TUNE: Q_y = 59.32$ $\Delta Q \le 3 \cdot 10^{-3}$ HORIZONTAL TUNE: $Q_z = 64.28$

AT 450GeV HORIZONTAL TUNE: $Q_x = 64.28$ VERTICAL TUNE: $Q_y = 59.31$

Betatron tunes should avoid linear coupling resonances at

$$nQ_x + mQ_y = p$$

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Back to transverse motion and magnet realities!

- Taylor expanding the dipole magnetic field AND fulfilling Maxwell equations gives
 - $(B_x, B_y, B_z) \sim (-(n_0 B_0/r_g)y, B_0 (n_0 B_0/r_g)*x, 0)$

- NB! note that - sign is not good!

- Ideally we want n_0 as large as possible to confine the beam!
 - (And make the magnet as small as possible)
- Let us look at solution for x (r)!

The equation of motion for x

$$\gamma m \frac{d^2 r}{dt^2} = \gamma m v^2 \frac{d^2 r}{dz^2} = \gamma m \frac{v^2}{r} - q v B_y$$

$$\frac{d^2r}{dz^2} = \frac{1}{r} - \frac{q}{\gamma mv}B_y$$

Substituting $x = r - r_g + \text{expanding } \frac{1}{r}$:

$$\frac{d^2x}{dz^2} = \frac{1}{r_g} - \frac{1}{r_g^2}x - \frac{q}{\gamma mv}B_y$$

Inserting the Taylor expansion of B_y :

$$\frac{d^2x}{dz^2} = \frac{1}{r_g} - \frac{qB_0}{\gamma mv} - \frac{1}{r_g^2}x + \frac{qn_0B_0}{\gamma mvr_g}x$$

The first two terms gives the solution for the ideal trajectory \rightarrow : $\frac{1}{r_g} = \frac{qB_0}{\gamma mv}$ so that:

$$\frac{d^2x}{dz^2} = -\frac{1}{r_g^2}(1-n_0)x$$

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Weak focusing: $0 < n_0 < 1$

$$\frac{d^2x}{dz^2} = -\frac{1}{r_g^2}(1-n_0)x.$$

- Only harmonic oscillation solution when (1n₀)>0 (and y equation requires n₀>0)
 - Otherwise exponential growth!
- This means that the focusing is limited!
 - That is why this solution is called weak focusing

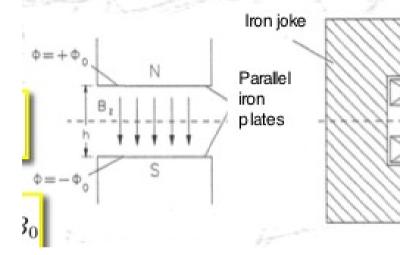
LHC example

- LHC focuses the beam using magnetic fields of 223 T/m
- If one would use weak focusing the magnetic field should be smaller than
 8.33T/2800m = 0.003 T/m!!!!
- And so the ring would have to be enormous!
 - And the luminosity would be very small!

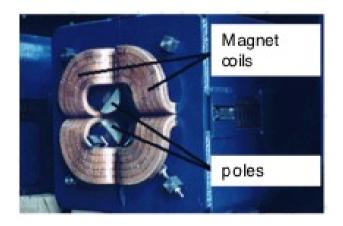
Can we find better focusing? Beamline Elements

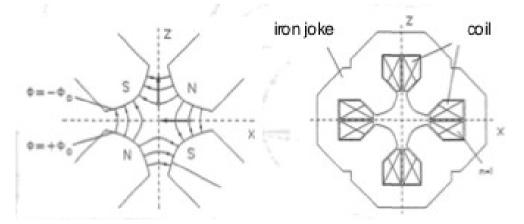
Dipole (bend) magnets





Quadrupole (focusing) magnets



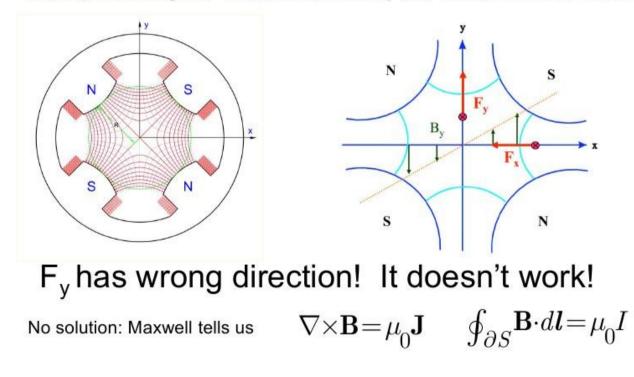


30/3-2015 Also Sextupoles and beyond ture 5 and 6 P. Christiansen (Lund) ωi

Quadropoles has similar problem!

FOCUSING OF THE PROTON BEAM

Quadrupole looks good - field increases linarly with distance from the center.



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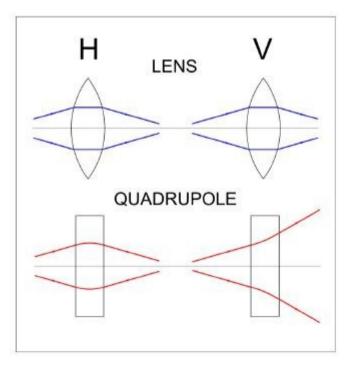
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PROBLEM

Quarupole is convergent lens in horizontal, but divergent in vertical direction!

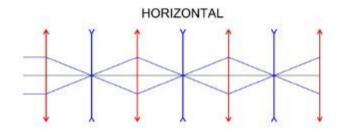


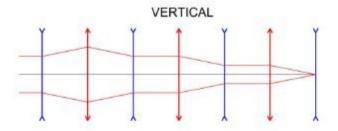
There was no solution until 1952, and it is beautiful and simple:

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SOLUTION: AG OR STRONG FOCUSING





FODO LATTICE

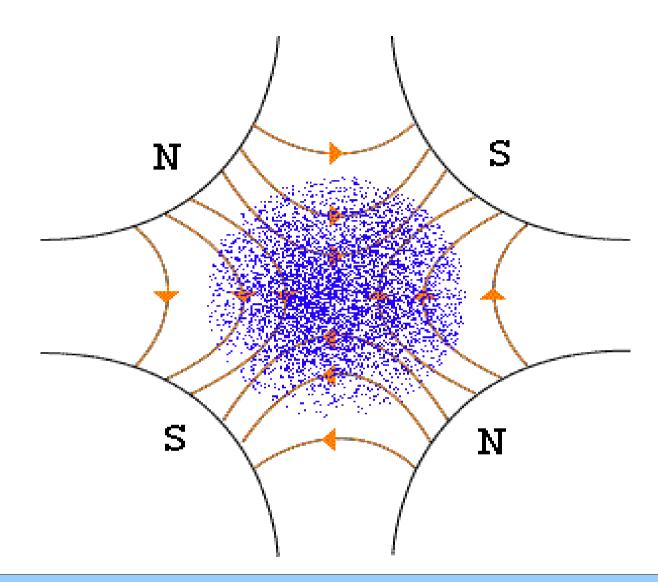
- F focusing
- D defocusing
- O drift space or dipoles

If we have alternating convergent and divergent lenses with right spacing overall effect is focusing!

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Lecture 5 and 6 P. Christiansen (Lund) 23

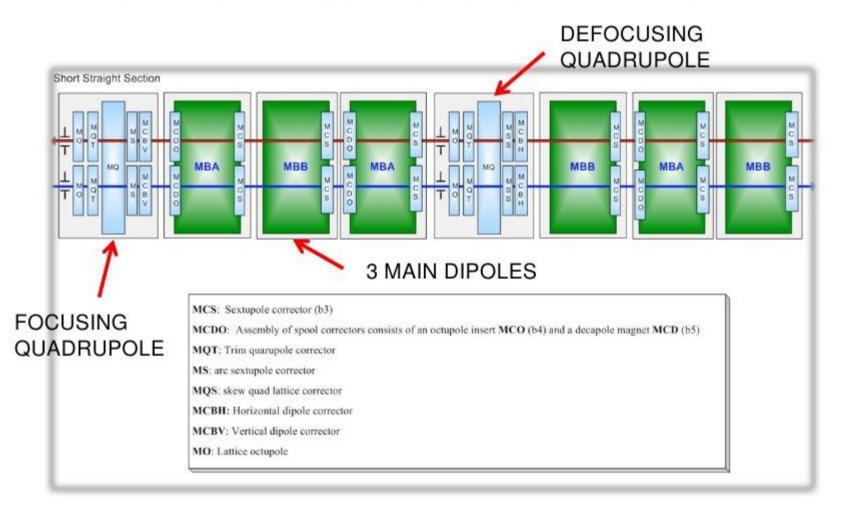


Animation taken from The Physics of Accelerators

Slides by C.R. Prior Rutherford Appleton Laboratory and Trinity College, Oxford

LHC FODO LATTICE CELL (106.9 m)

The pattern of bending and focusing magnets is called lattice.

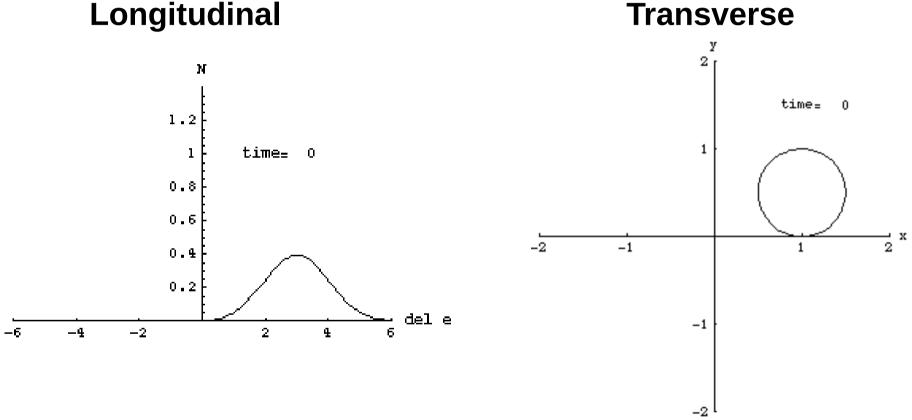


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Energy adjusting by AC (longitudinal) & transverse strong focusing

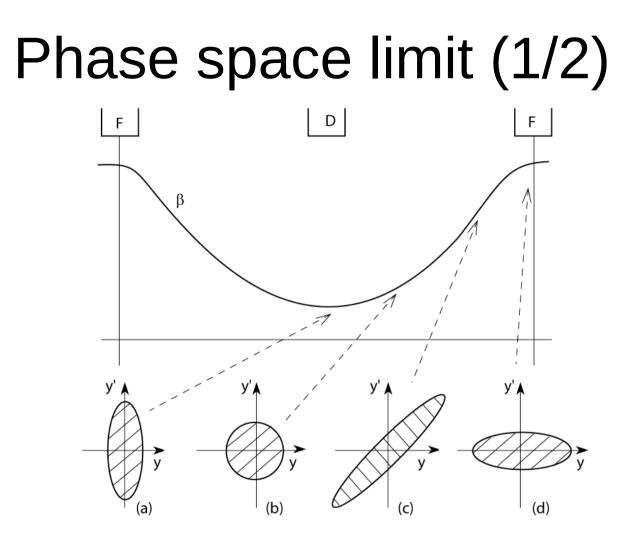


- "Catching the beam" animations taken from
 - http://www.lns.cornell.edu/~dugan/USPAS/

Lecture 5 and 6 P. Christiansen (Lund)

Beam dynamics

- Phase space limits
- The beta function
- Focusing the beam at the interaction point



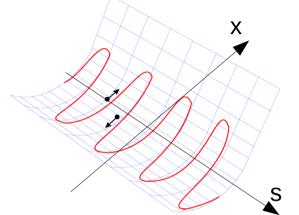
• Liouville's theorem states that for most beams the phase space area cannot change

Phase space limit (2/2)

- In reality the limit is on p_y*y:
 - Constant = $p_y * y \sim \gamma m \beta_y * y$ where $\beta_y = dy/dt$ $\sim \gamma m \beta * y' * y$ where y' = dy/dz
 - So the phasespace limit implies that the area of the phase space ellipse for y'*y (the beam emittance) decreases as 1/p
 - This is called adiabatic damping. The physical size of the beam decreases as it is accelerated. The width decreases as 1/√p [the other 1/√p is the decreasing divergence].

The beta function (1/2)

The "gutter" analogy



Based on: ACCELERATORS FOR PEDESTRIANS

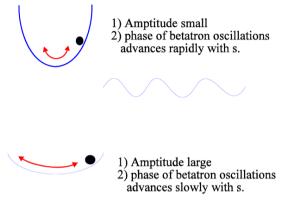
Simon Baird

$$x = \sqrt{\varepsilon \beta(s)} \cos(\Psi(s) + \varphi)$$

- The idea is to separate the transverse motion into two parts:
 - The initial conditions: ϵ (emittance) and ϕ
 - A part depending the focusing and de-focusing by magnets: $\beta(s)$ and $\Psi(s)$
 - The $\beta(s)$ function depends on where in the accelerator we are and <u>not</u> related to the velocity β)

The beta function (2/2)

The "gutter" analogy



Based on: ACCELERATORS FOR PEDESTRIANS

Simon Baird

$$x = \sqrt{\varepsilon \beta(s)} \cos(\Psi(s) + \phi)$$
$$x' = \sqrt{\frac{\varepsilon}{\beta(s)}} \sin(\Psi(s) + \phi)$$

- So the beta function is related to how strong focusing we do and is the one we want to optimize in particular at the interaction points
- Note that x*x'~ε so that the phase space area is still the same

Recall lecture 1 and 2

Intensity or brightness of an accelerator

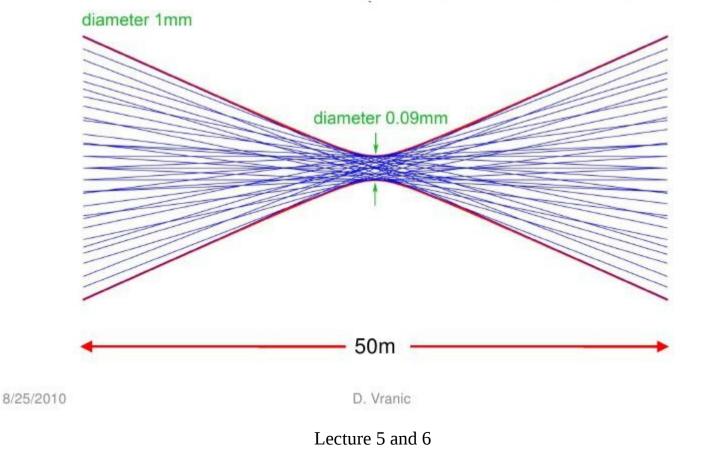
 $N = \mathcal{L} \cdot \sigma$

- Events Seen = Luminosity * cross-section
 Rare processes (fb) need lots of luminosity (fb⁻¹)
- In a storage ring

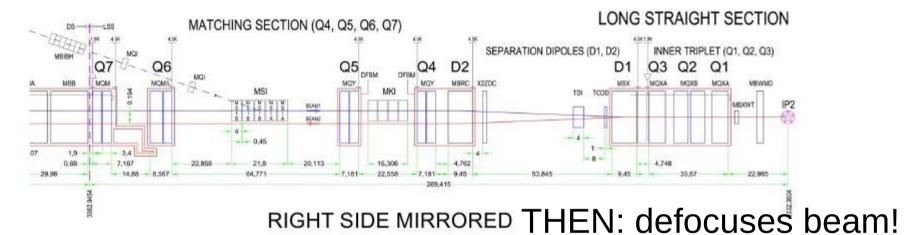
$$\mathcal{L} = \underbrace{\frac{1}{4\pi} \frac{f \cdot N_1 \cdot N_2}{\sigma_x \cdot \sigma_y}}_{\text{(Spot size)}}$$

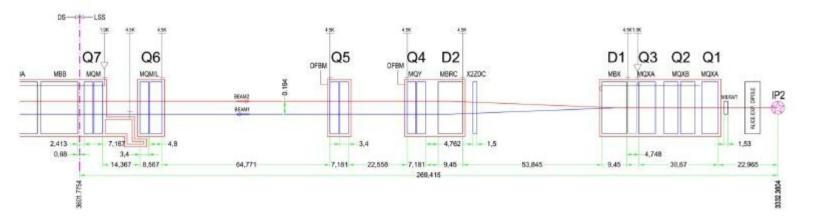
Where *f* is the revelation frequency multiplied by *#* of colliding bunches More particles through a smaller area means more collisions

Example of focusing for collisions at P2 (ALICE)



SYMMETRY! LEFT SIDE EXAMPLE: Focuses beam!





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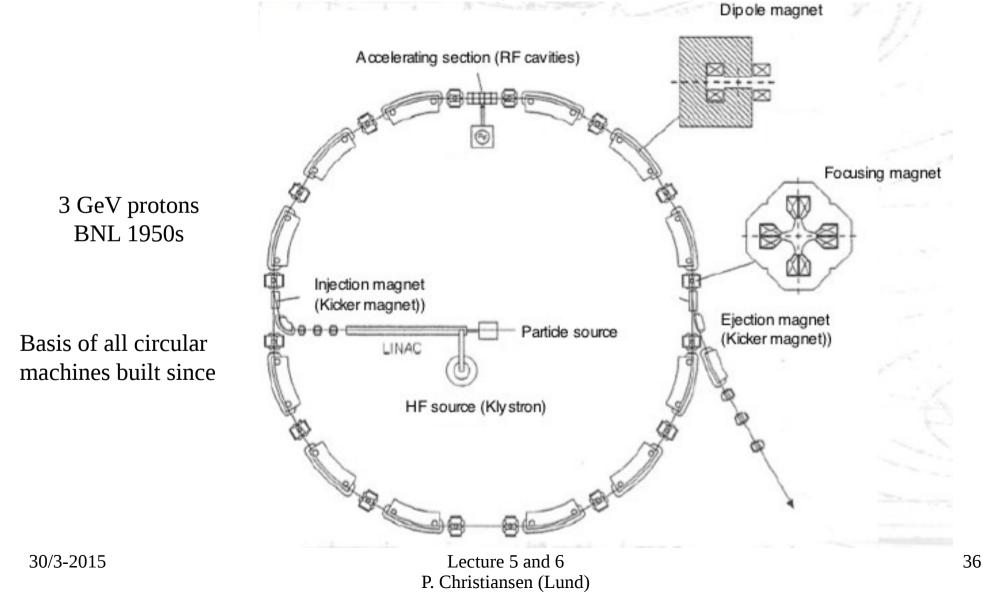
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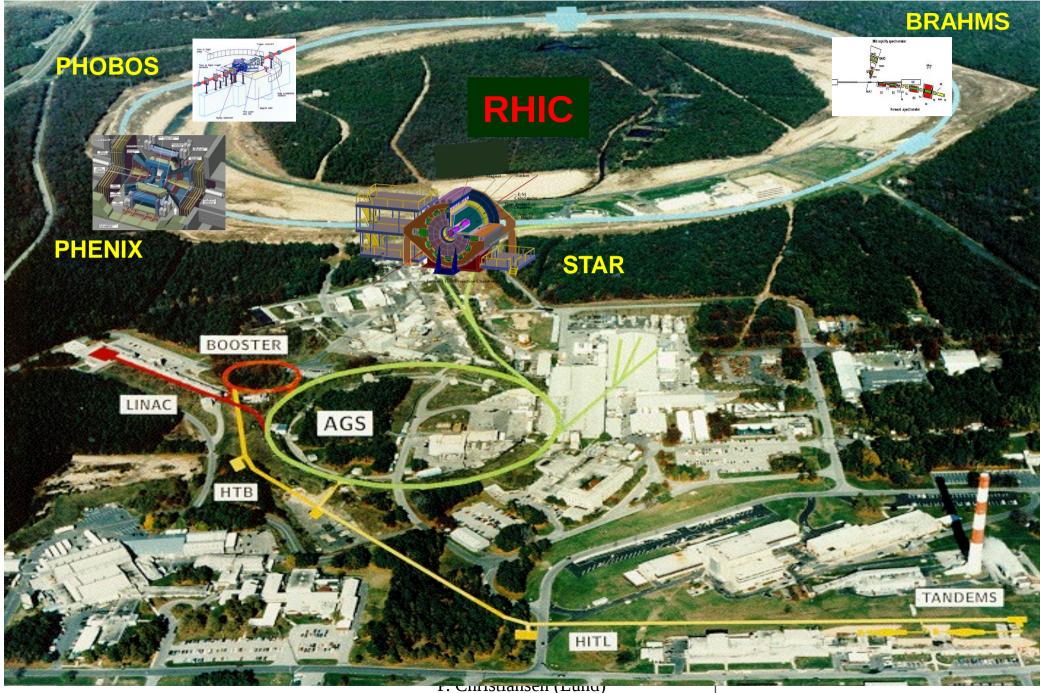
Examples of synchrotrons

Synchrotrons

Use smaller magnets in a ring + accelerating station

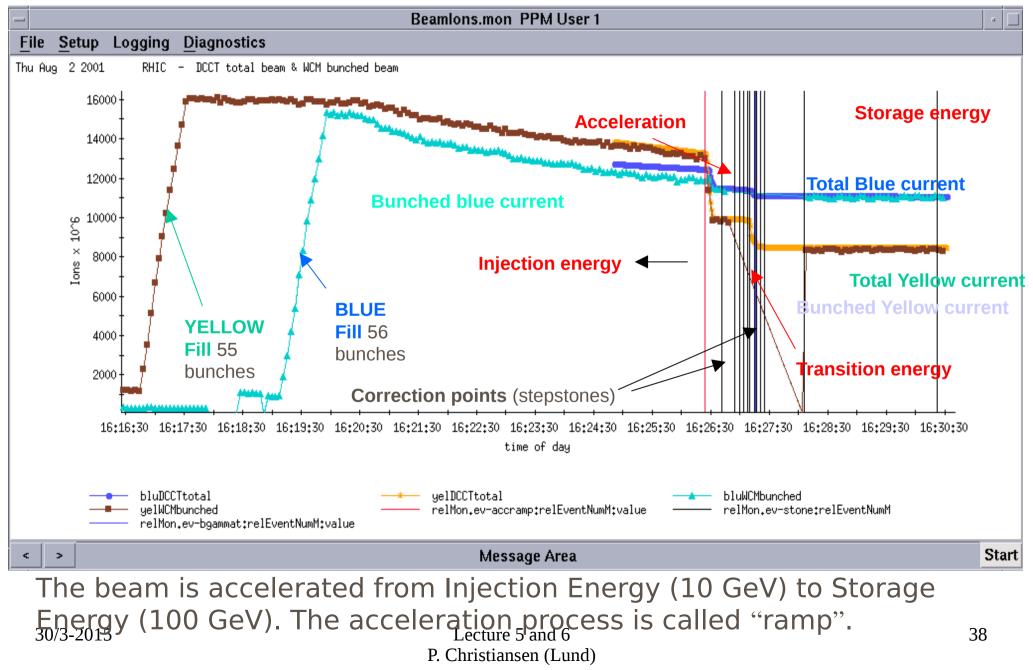


at Brookhaven National

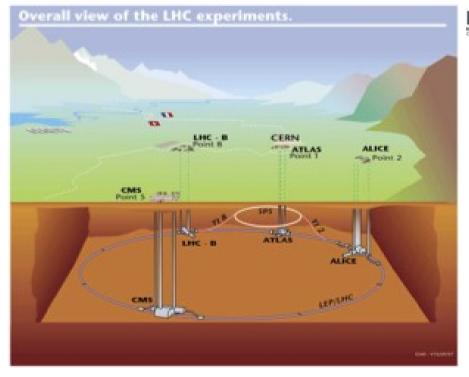


fysN15 Accelerators 4

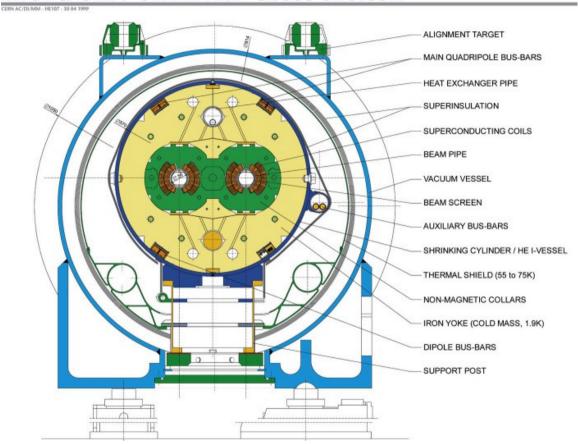
RHIC ramp with 56 bunches

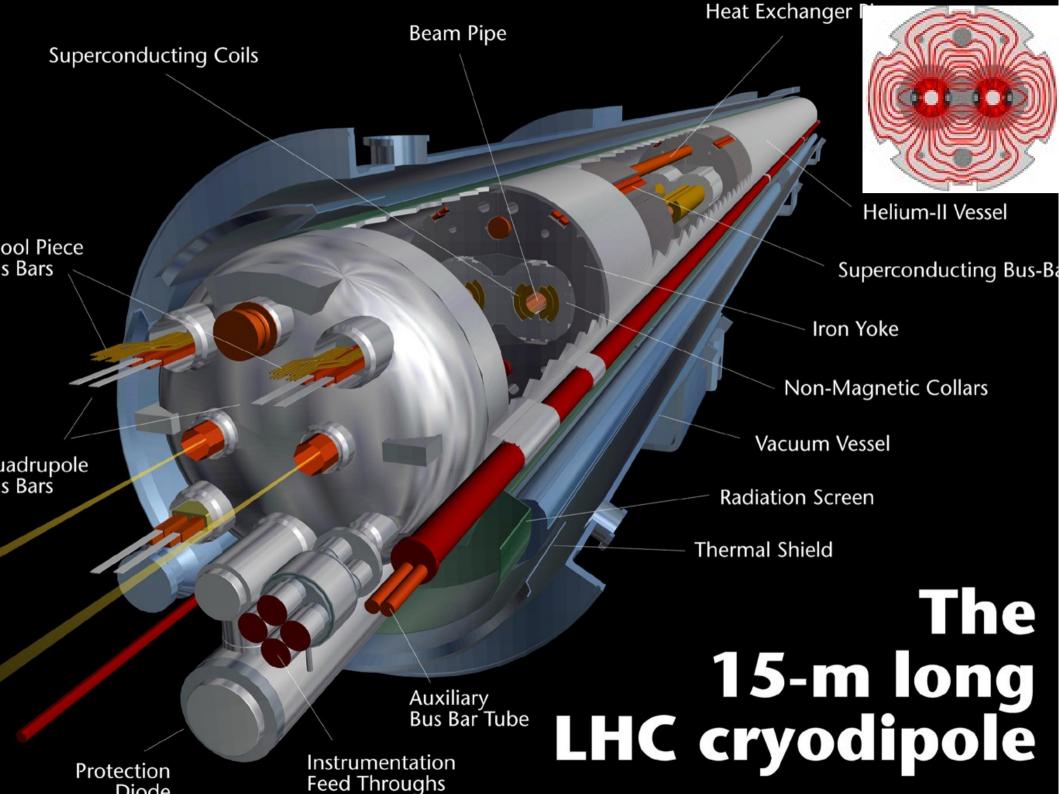


CERN Large Hadron Collider

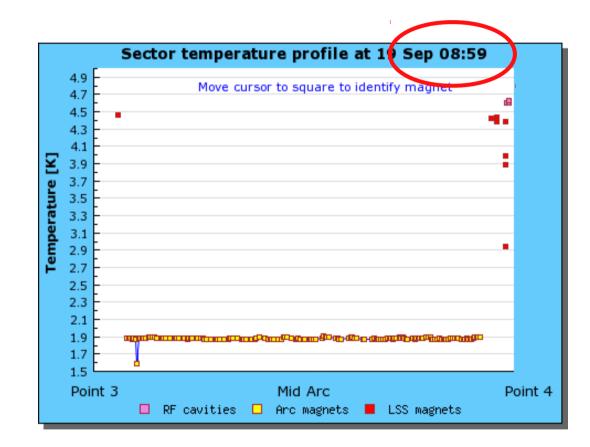


LHC DIPOLE : STANDARD CROSS-SECTION





The 19 September 2008 accident

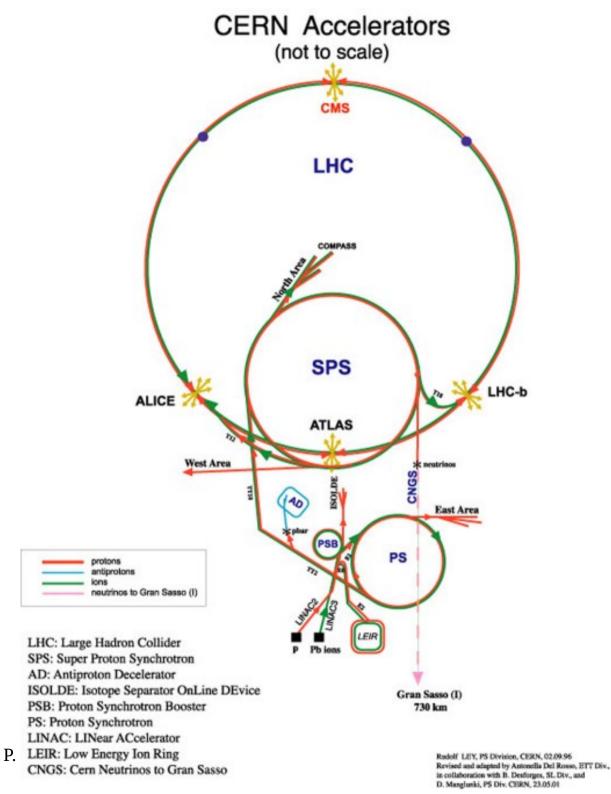


THE 19 SEPTEMBER 2008 INCIDENT



CERN Complex

Old rings still in use Many different programs

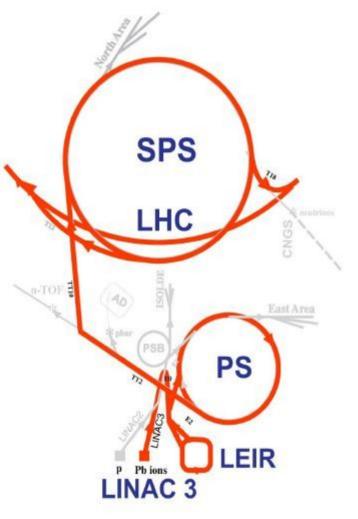


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ION BEAM IN THE LHC

ECR ion source

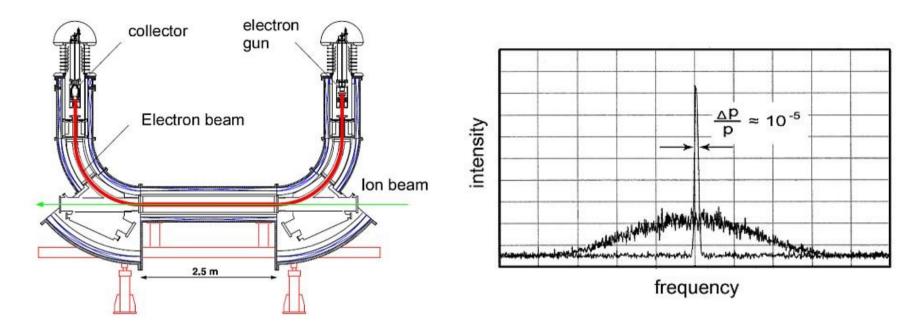
- Provide highest possible intensity of Pb29+
- RFQ + Linac 3
- Adapt to LEIR injection energy
- strip to Pb54+
- LEIR
- Accumulate and cool Linac 3 beam
- Prepare bunch structure for PS
- PS
- Define LHC bunch structure
- Strip to Pb82+
- SPS
- Define filling scheme



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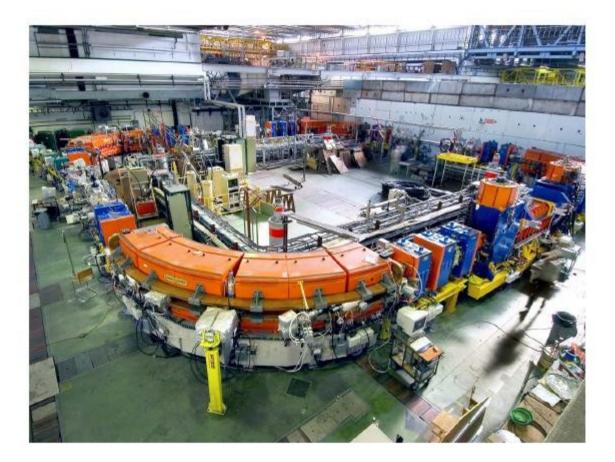
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LEIR: Electron cooling example



- From: http://web-docs.gsi.de/
- Elastic collision e+ion will decrease the relative momentum spread in the beam

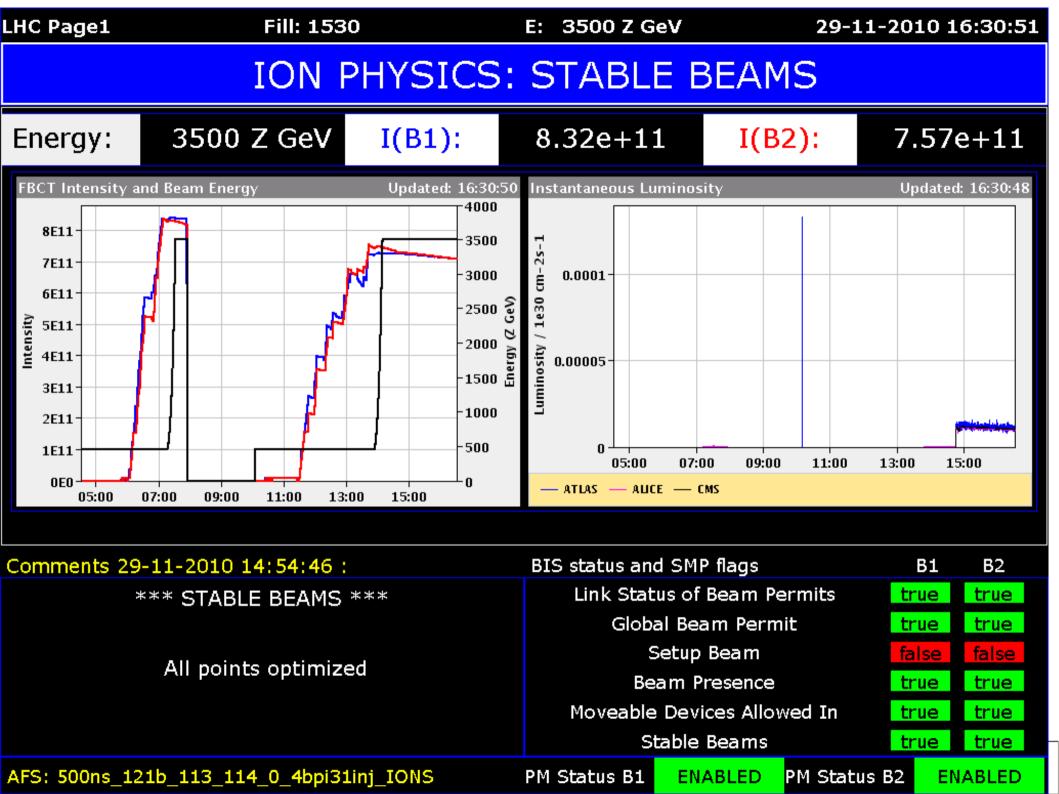
LEIR



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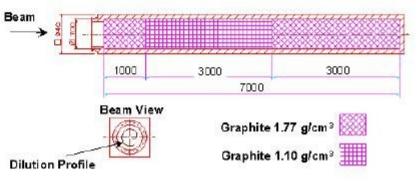
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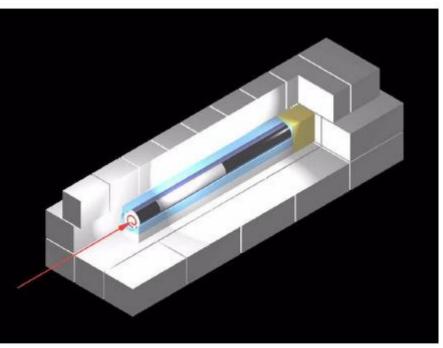
DUMP CORE TDE

7m long C / C-C TDE in steel shrinkcylinder, followed by 1m Al, 2m Fe ~1000 T of concrete shielding



This is the ONLY element in the LHC that can withstand the impact of the full 7 TeV beam ! Nevertheless, the dumped beam must be painted to keep the peak energy densities at a tolerable level !

Why graphite? If the material were heavy, all the beam's energy would concentrate in the first half meter of the block.

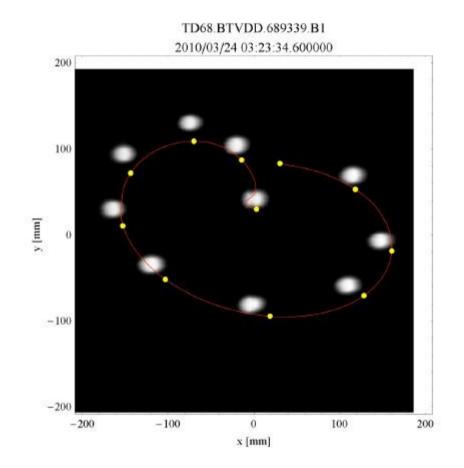


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Lecture 5 and 6 P. Christiansen (Lund)

The beam size has increased to an extent where the sigma is 1.6mm in both planes.







Lecture 5 and 6 P. Christiansen (Lund)