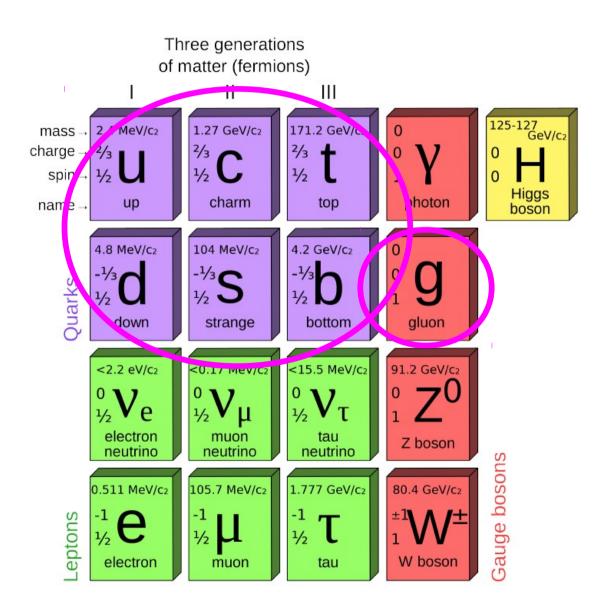
Quarks and Hadrons



Quantum Chromo Dynamics (QCD)

3 color charges (red, green, blue)

Not real colors but e.g. qx, qy, qz that can be +qx for quarks (red) and -qx for anti-quarks (anti-red)

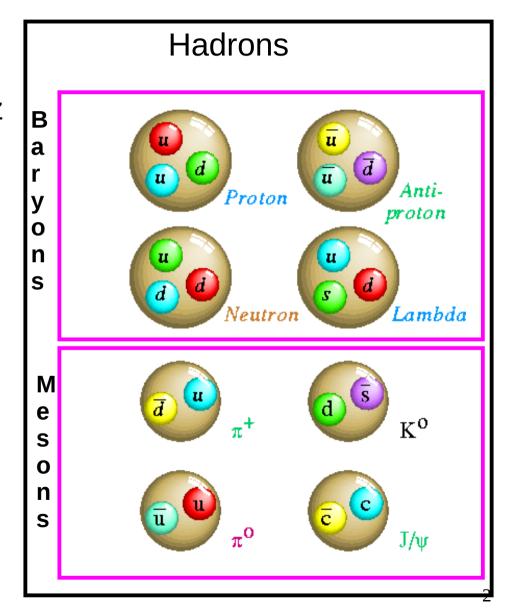
Hadrons have to be colorless

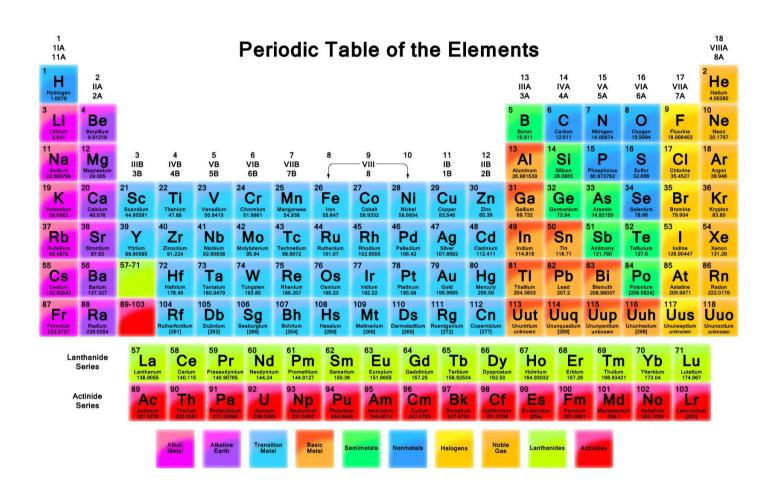
Baryons have all 3 colors

Mesons has a color and an anticolor

A single quark cannot be observed because it has color!

The quarks are confined inside the hadrons!





Today: the periodic table of the hadrons:-)

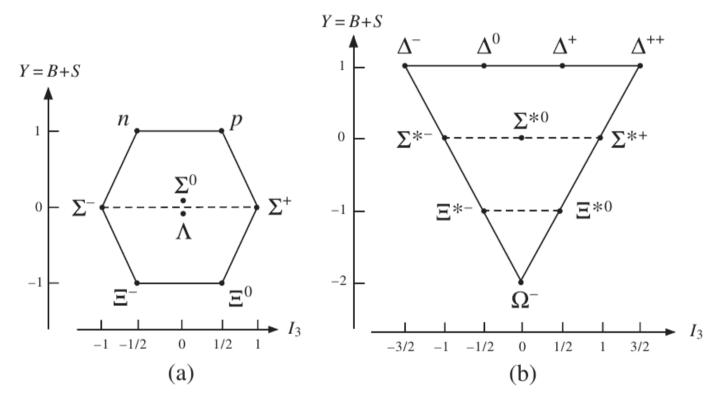


Figure 6.2 Weight diagrams for (a) the $J^P = \frac{1}{2}^+$ octet of light baryons and (b) the $J^P = \frac{3}{2}^+$ baryon decuplet.

Example of mesonic "energy levels": the su system

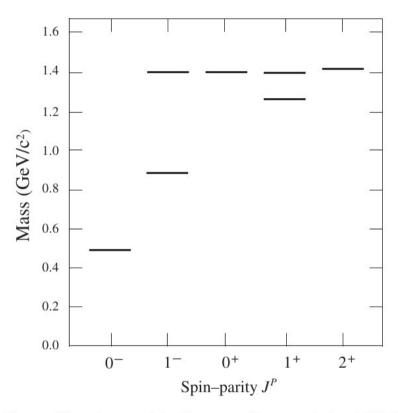


Figure 3.12 Observed bound states of the $s\bar{u}$ system with masses below 1.5 GeV/c², together with values of their spin-parities⁹ J^P . The ground state is the $K^-(494)$ and the others can be interpreted as its excited states.

Example of mesonic "energy levels": the su system

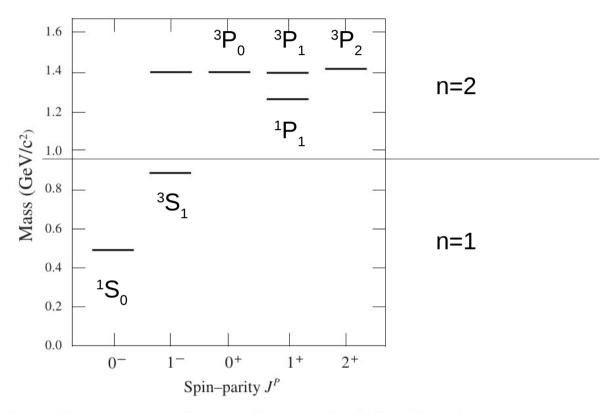


Figure 3.12 Observed bound states of the $s\bar{u}$ system with masses below 1.5 GeV/c², together with values of their spin-parities⁹ J^P . The ground state is the $K^-(494)$ and the others can be interpreted as its excited states.

L=0 and n=1 mesonic uds states

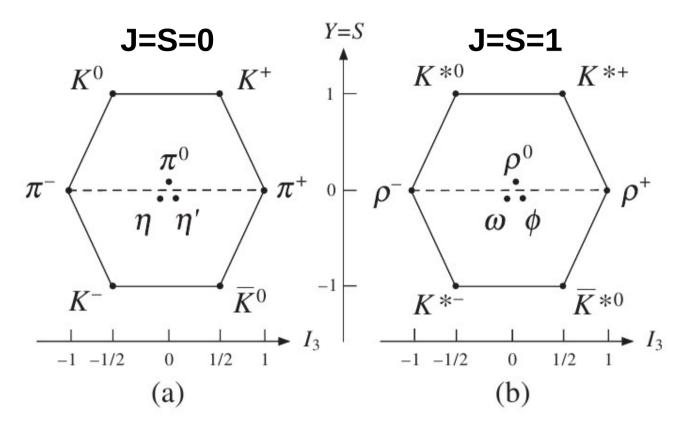


Figure 6.1 Weight diagrams for (a) the 0⁻ meson nonet and (b) the 1⁻ meson nonet.

L=0 and n=1 mesonic uds states

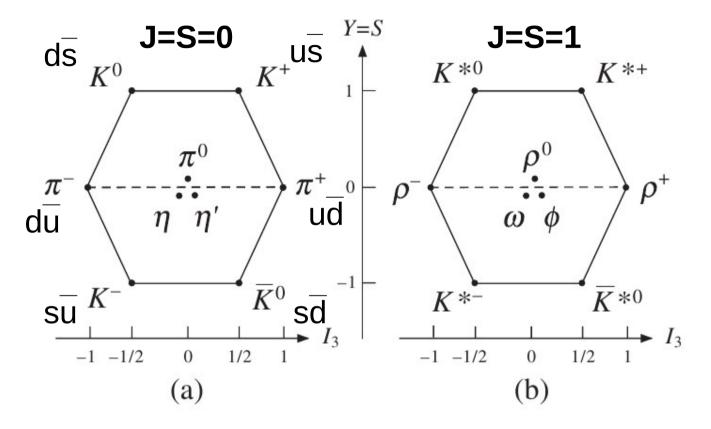


Figure 6.1 Weight diagrams for (a) the 0⁻ meson nonet and (b) the 1⁻ meson nonet.

L=0 and n=1 mesonic uds states

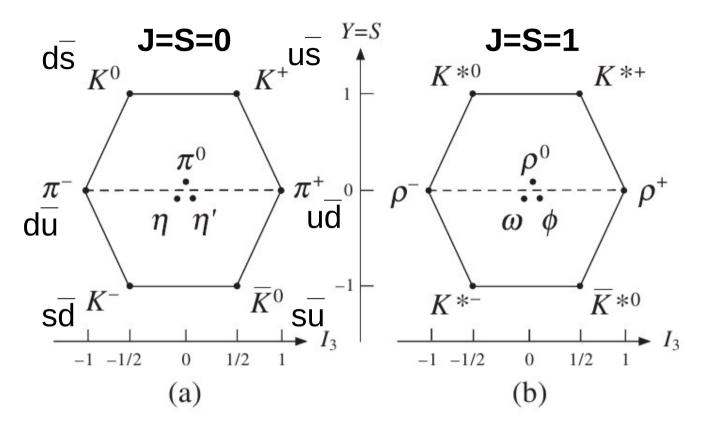


Figure 6.1 Weight diagrams for (a) the 0⁻ meson nonet and (b) the 1⁻ meson nonet.

$$\pi^{0}, \, \rho^{0} \qquad \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \quad (I = 1, I_{3} = 0),$$
(6.23)

Mixture:

$$\eta, \eta', \omega, \Phi$$

$$\begin{cases}
\frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d}) & (I = 0, I_3 = 0) \\
s\bar{s} & (I = 0, I_3 = 0),
\end{cases}$$
(6.24a)

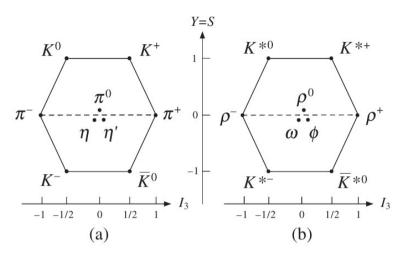


Figure 6.1 Weight diagrams for (a) the 0⁻ meson nonet and (b) the 1⁻ meson nonet.

TABLE 6.4 The states of the light L = 0 meson nonets.

Quark content	0 ⁻ state	1 ⁻ state	I_3	I	Y = S
$u\bar{s}$ $d\bar{s}$	K ⁺ (494) K ⁰ (498)	<i>K</i> *+(892) <i>K</i> *0(896)	1/2 -1/2	1/2 1/2	1 1
$\frac{u\bar{d}}{(u\bar{u} - d\bar{d})}$	$\pi^{+}(140)$ $\pi^{0}(135)$	$\rho^+(768)^\#$ $\rho^0(768)^\#$	1 0	1 1	0
$\sqrt{2}$ $d\bar{u}$	$\pi^{-}(140)$	$\rho^{-}(768)^{\#}$	-1	1	0
sā sū	$\bar{K}^{0}(498)$ $K^{-}(494)$	$\bar{K}^{*0}(896)$ $K^{*-}(892)$	1/2 -1/2	1/2 1/2	-1 -1
See text See text	η(549) η'(958)	$\omega(782) \\ \phi(1019)$	0	0	0

^{*} The measured mass difference between the neutral and charged ρ mesons is $m(\rho^0)-m(\rho^+)=0.3\pm2.2\,{\rm MeV/c^2}.$

TABLE 6.9 Predicted $c\bar{c}$ and $b\bar{b}$ states with principal quantum numbers n=1 and 2 and radial quantum number $n_r=n-L$, compared with experimentally observed states. Masses are given in MeV/ c^2 .

$^{2S+1}L_J$	n	n_r	$J^{\scriptscriptstyle PC}$	$c\bar{c}$ state	$b\bar{b}$ state
$^{1}S_{0}$	1	1	0-+	$\eta_c(2980)$	$\eta_b(9300)$ #
$^{3}S_{1}$	1	1	1	$J/\psi(3097)$	$\Upsilon(9460)$
${}^{3}P_{0}$	2	1	0_{++}	$\chi_{c0}(3415)$	$\chi_{b0}(9859)$
${}^{3}P_{1}$	2	1	1++	$\chi_{c1}(3511)$	$\chi_{b1}(9893)$
$^{3}P_{2}$	2	1	2++	$\chi_{c2}(3556)$	$\chi_{b2}(9913)$
${}^{1}P_{1}$	2	1	1+-	$h_c(3526)\#$	
${}^{1}S_{0}$	2	2	0_{-+}	$\eta_c(3638)$	
${}^{3}S_{1}$	2	2	1	$\psi(3686)$	$\Upsilon(10023)$

[#] State is not well established and its quantum number assignments are unknown.

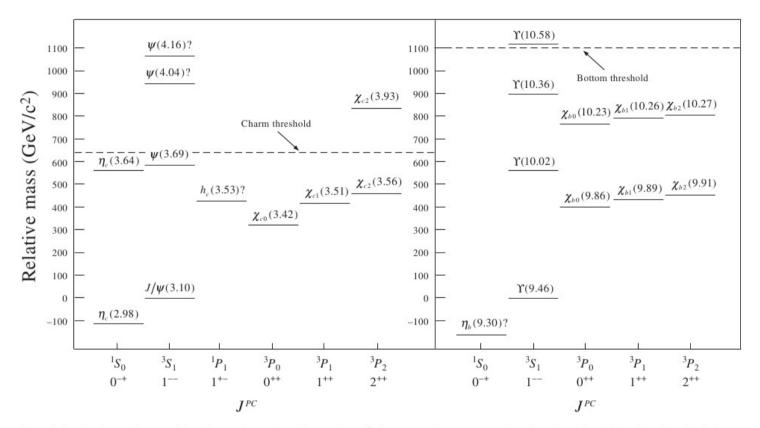


Figure 6.6 The observed states of the charmonium $(c\bar{c})$ and bottomium $(b\bar{b})$ for $L \leq 1$. The masses are given in units of GeV/ c^2 and are plotted relative to that of the 3S_1 ground state.

Why is the J/ ψ (1⁻) famous? (the η_c (0⁻)is lighter!)

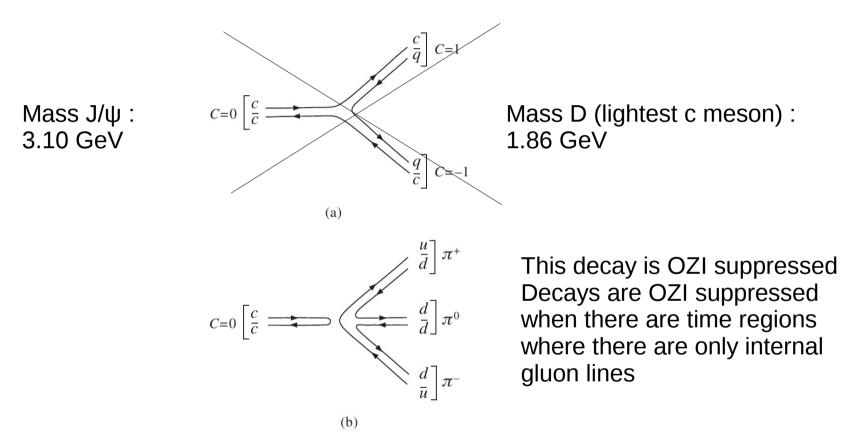
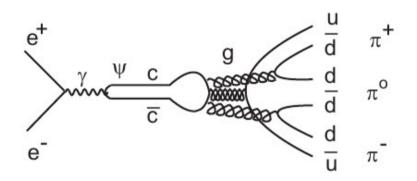
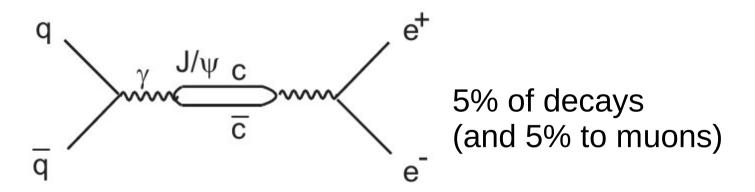


Figure 6.5 Quark diagrams for (a) the decay of a charmonium state to a pair of charmed mesons and (b) an example of a decay to noncharmed mesons.

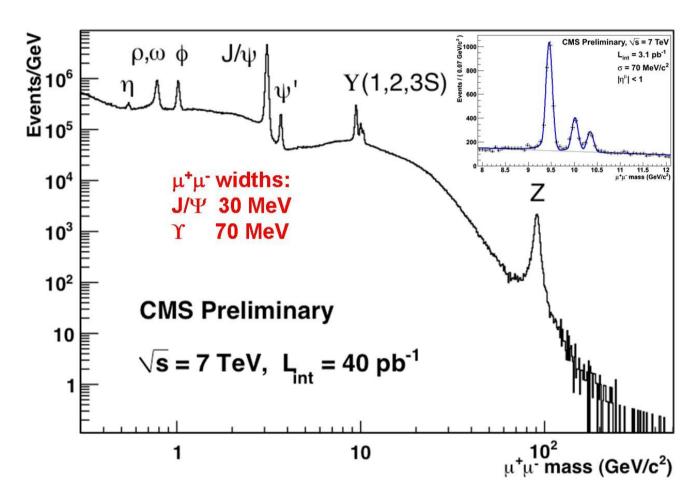
J/ψ (1⁻) has quantum number of virtual photon!



OZI suppressed!



Very easy identification! (if you have a good detector:-)



Note that except for η these are all 1⁻ states!

What can we more do with energy levels: model them!

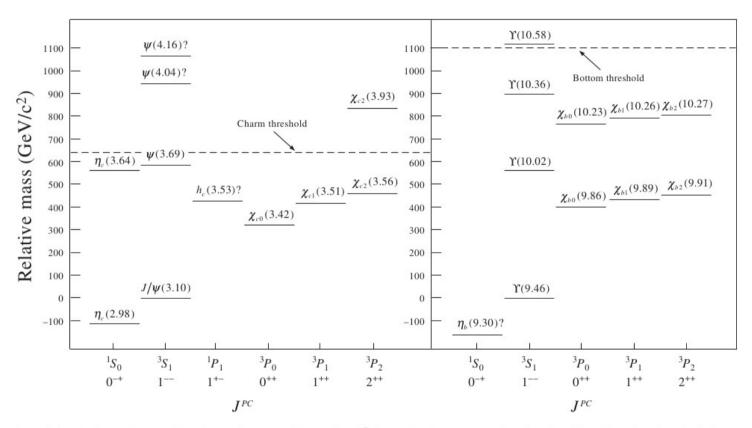


Figure 6.6 The observed states of the charmonium $(c\bar{c})$ and bottomium $(b\bar{b})$ for $L \le 1$. The masses are given in units of GeV/ c^2 and are plotted relative to that of the 3S_1 ground state.

A way to experimentally measure the strong potential

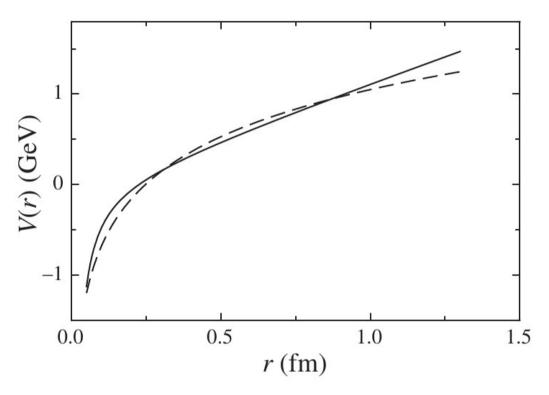
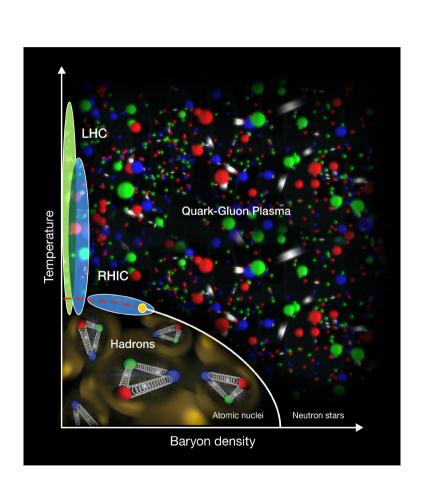


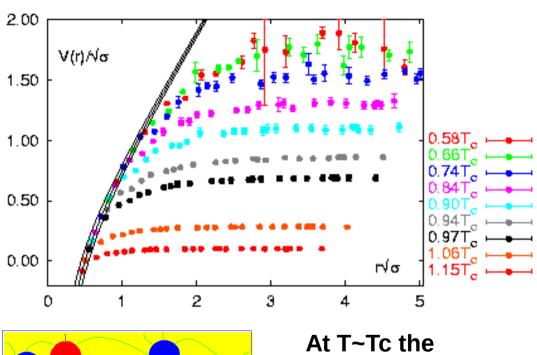
Figure 6.8 Heavy quark—antiquark potentials obtained from fitting the energy levels of charmonium and bottomium. The solid and dashed lines show the results obtained from the forms (6.57) and (6.58), respectively.

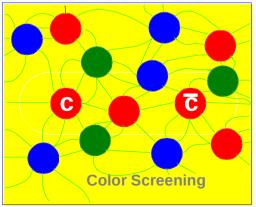
A small excursion

Lattice QCD results (Numerical non-perturbative)

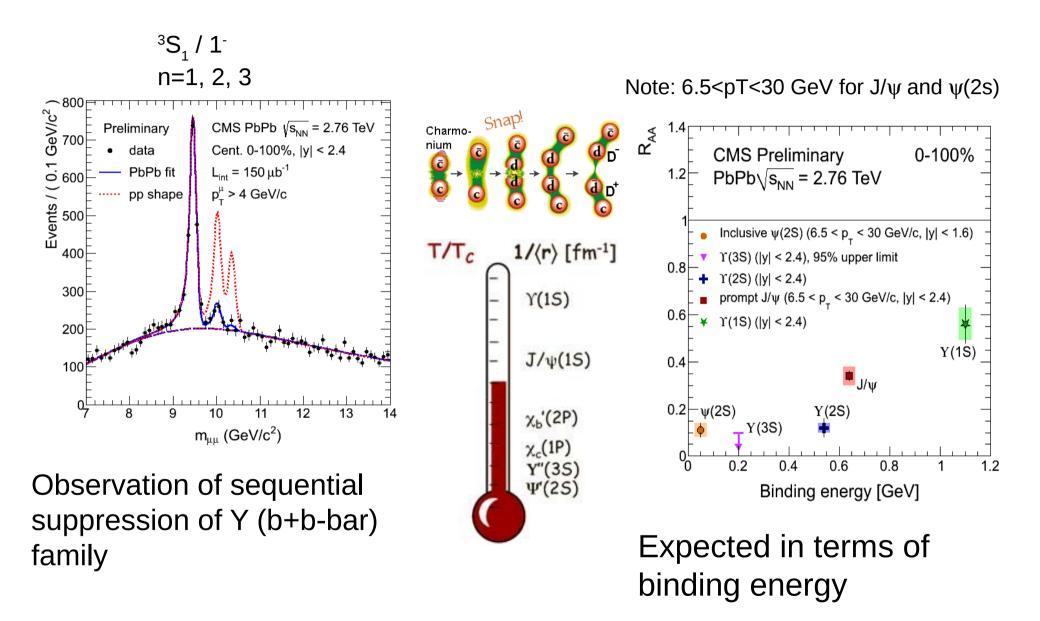
Heavy quark potential







strong potential is screened so e.g. c+c-bar states can disassociate.



Unfortunately heavy quark results are more complex when systematically studied!

Next: let us understand the baryons!

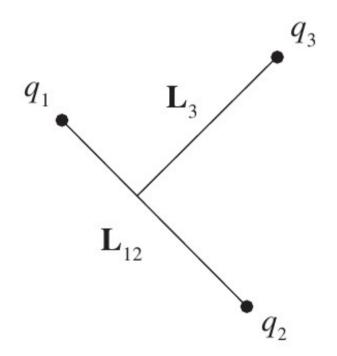


Figure 5.1 Internal orbital angular momenta of a three-quark state.

Only consider L = 0!

The baryonic systems with L=0 and n = 0

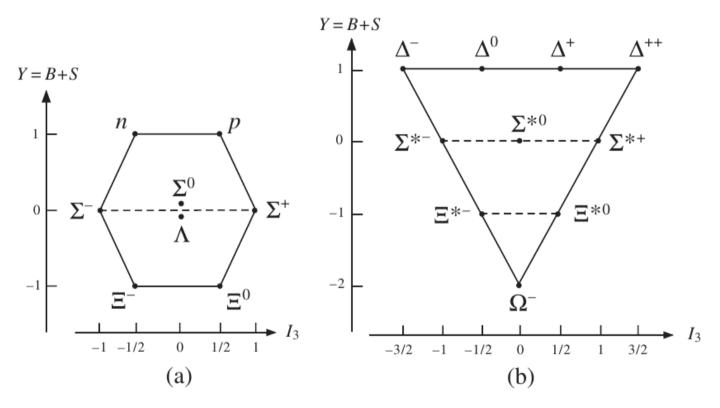


Figure 6.2 Weight diagrams for (a) the $J^P = \frac{1}{2}^+$ octet of light baryons and (b) the $J^P = \frac{3}{2}^+$ baryon decuplet.

Color is needed to make the Δ^{++} wavefunction antisymmetric!

$$\chi_B^C = \frac{1}{\sqrt{6}} (r_1 g_2 b_3 - g_1 r_2 b_3 + b_1 r_2 g_3 - b_1 g_2 r_3 + g_1 b_2 r_3 - r_1 b_2 g_3), \tag{6.36}$$

- This means that the rest of the wavefunction (spin, flavor, L) symmetric!